

# Near Horizon Soft Hairs as Black Hole Microstates

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Seminario-almuerzo  
Valdivia, November 2016



1603.04824, 1607.05360, 1607.00009

## Two simple punchlines

### 1. Heisenberg algebra

$$[X_n, P_m] = i \delta_{n,m}$$

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### 2. Black hole microstates identified as specific “soft hair” descendants

based on work with

- ▶ **Hamid Afshar** [IPM Teheran]
- ▶ Stephane Detournay [ULB]
- ▶ Wout Merbis [TU Wien]
- ▶ Blagoje Oblak [ULB]
- ▶ Alfredo Perez [CECS Valdivia]
- ▶ Stefan Prohazka [TU Wien]
- ▶ **Shahin Sheikh-Jabbari** [IPM Teheran]
- ▶ David Tempo [CECS Valdivia]
- ▶ Ricardo Troncoso [CECS Valdivia]

# Outline

Motivation

Near horizon boundary conditions

Explicit construction of BTZ microstates

Discussion

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- ▶ Generalizations in 2+1 gravity/gravity-like theories (Galilean CFT, warped CFT, ...)

warped CFT: Detournay, Hartman, Hofman '12

Galilean CFT: Bagchi, Detournay, Fareghbal, Simon '13; Barnich '13

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- ▶ Generalizations in 2+1 gravity/gravity-like theories (Galilean CFT, warped CFT, ...)
- ▶ Main idea: consider near horizon symmetries for non-extremal horizons
- ▶ Near horizon line-element with **Rindler acceleration  $a$** :

$$ds^2 = -2a\rho dv^2 + 2dv d\rho + \gamma^2 d\varphi^2 + \dots$$

### Meaning of coordinates:

- ▶  $\rho$ : radial direction ( $\rho = 0$  is horizon)
- ▶  $\varphi \sim \varphi + 2\pi$ : angular direction
- ▶  $v$ : (advanced) time

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of **Rindler** metric

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We make this choice in this talk!

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- ▶ Work in 3d Einstein gravity in Chern–Simons formulation

$$I_{CS} = \pm \sum_{\pm} \frac{k}{4\pi} \int \langle A^{\pm} \wedge dA^{\pm} + \frac{2}{3} A^{\pm} \wedge A^{\pm} \wedge A^{\pm} \rangle$$

with  $sl(2)$  connections  $A^{\pm}$  and  $k = \ell/(4G_N)$  with AdS radius  $\ell = 1$

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## Diagonal gauge

Standard trick: partially fix gauge

$$A^\pm = b_\pm^{-1}(\rho) (d + \mathbf{a}_\pm(x^0, x^1)) b_\pm(\rho)$$

with some group element  $b \in SL(2)$  depending on radius  $\rho$  with  $\delta b = 0$

Drop  $\pm$  decorations in most of talk

Manifold topologically a cylinder or torus, with radial coordinate  $\rho$  and boundary coordinates  $(x^0, x^1) \sim (v, \varphi)$

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- ▶ Standard AdS<sub>3</sub> approach: highest weight gauge

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- ▶ Precise boundary conditions ( $\zeta$ : chemical potential):

$$\mathfrak{a} = (\mathcal{J} d\varphi + \zeta dv) L_0 \quad \delta \mathfrak{a} = \delta \mathcal{J} d\varphi L_0$$

and  $b = \exp(\frac{1}{\zeta} L_+) \cdot \exp(\frac{\rho}{2} L_-)$ . (assume constant  $\zeta$  for simplicity)

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$$ds^2 = -2a\rho f dv^2 + 2dv d\rho - 2\omega a^{-1} d\varphi d\rho \\ + 4\omega\rho f dv d\varphi + \left[ \gamma^2 + \frac{2\rho}{a} f (\gamma^2 - \omega^2) \right] d\varphi^2$$

state-dependent functions  $\mathcal{J}^{\pm} = \gamma \pm \omega$ , chemical potentials  $\zeta^{\pm} = -a \pm \Omega$

For simplicity set  $\Omega = 0$  and  $a = \text{const.}$  in metric above

EOM imply  $\partial_v \mathcal{J}^{\pm} = \pm \partial_{\varphi} \zeta^{\pm}$ ; in this case  $\partial_v \mathcal{J}^{\pm} = 0$

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Neglecting rotation terms ( $\omega = 0$ ) yields **Rindler** plus higher order terms:

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Comments:

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- ▶  $\gamma = \gamma(\varphi)$ : “black flower”

## Canonical boundary charges

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- ▶ Zero mode charges: mass and angular momentum

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Background independent result for Chern–Simons yields

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Meaningful near horizon boundary conditions and non-trivial theory!

## Near horizon symmetry algebra

- ▶ **Near horizon symmetry algebra** = all near horizon boundary conditions preserving trafo, modulo trivial gauge trafo

Most general trafo

$$\delta_\epsilon \mathbf{a} = d\epsilon + [\mathbf{a}, \epsilon] = \mathcal{O}(\delta \mathbf{a})$$

that preserves our boundary conditions for constant  $\zeta$  given by

$$\epsilon = \epsilon^+ L_+ + \eta L_0 + \epsilon^- L_-$$

with

$$\partial_v \eta = 0$$

implying

$$\delta_\epsilon \mathcal{J} = \partial_\varphi \eta$$

## Near horizon symmetry algebra

- ▶ Near horizon symmetry algebra = all near horizon boundary conditions preserving trafos, modulo trivial gauge trafos
- ▶ Expand charges in Fourier modes

$$J_n^\pm = \frac{k}{4\pi} \oint d\varphi e^{in\varphi} \mathcal{J}^\pm(\varphi)$$

What should we expect?

- ▶ Virasoro? (spacetime is locally  $\text{AdS}_3$ )
- ▶  $\text{BMS}_3$ ? (Rindler boundary similar to scri)
- ▶ warped conformal algebra? (this is what we found for Rindleresque holography and what Donnay, Giribet, Gonzalez, Pino found in their near horizon analysis)

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$$[J_n^\pm, J_m^\pm] = \pm \frac{1}{2} kn \delta_{n+m,0} \quad [J_n^+, J_m^-] = 0$$

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- ▶ Map

$$P_0 = J_0^+ + J_0^- \quad P_n = \frac{i}{kn} (J_{-n}^+ + J_{-n}^-) \text{ if } n \neq 0 \quad X_n = J_n^+ - J_n^-$$

yields Heisenberg algebra (with Casimirs  $X_0, P_0$ )

$$[X_n, X_m] = [P_n, P_m] = [X_0, P_n] = [P_0, X_n] = 0$$

$$[X_n, P_m] = i\delta_{n,m} \quad \text{if } n \neq 0$$

## Map to asymptotic variables

- ▶ Usual asymptotic AdS<sub>3</sub> connection with chemical potential  $\mu$ :

$$\hat{A} = \hat{b}^{-1} (d + \hat{\mathbf{a}}) \hat{b} \quad \hat{\mathbf{a}}_\varphi = L_+ - \frac{1}{2} \mathcal{L} L_-$$

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- ▶ Get Virasoro with non-zero central charge  $\delta \mathcal{L} = 2\mathcal{L}\varepsilon' + \mathcal{L}'\varepsilon - \varepsilon'''$

## Remarks on asymptotic and near horizon variables

- ▶ Asymptotic spin-2 currents fulfill Virasoro algebra, but charges obey still **Heisenberg algebra**

$$\delta Q = -\frac{k}{4\pi} \oint d\varphi \varepsilon \delta \mathcal{L} = -\frac{k}{4\pi} \oint d\varphi \eta \delta \mathcal{J}$$

Reason: asymptotic “chemical potentials”  $\mu$  depend on near horizon charges  $\mathcal{J}$  and chemical potentials  $\zeta$

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Solved automatically from map to asymptotic observables; reminder:

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Near horizon boundary conditions natural for near horizon observer

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- ▶ Near horizon algebra (conveniently rescaled)

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- ▶ Call this “near horizon symmetry algebra” (note: independent from  $\ell$ )

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- ▶ Will exploit this property to provide cut-off on soft hair spectrum!

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## Proposed map between near horizon and asymptotic generators

- ▶ Suggestive proposal (see Bañados 9811162)

$$cL_0^\pm = \mathcal{L}_0^\pm - \frac{1}{24}$$

Note: in 9811162 whole Virasoro algebras are mapped, so even for  $n \neq 0$

$$cL_n^\pm = \mathcal{L}_{nc}^\pm$$

we use much weaker map above between zero modes as ansatz

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- ▶ Microstates = all states in near horizon Hilbert space obeying equations above

## Horizon fluffs as microstates

We are now ready to identify all BTZ microstates

- ▶ Vector space  $\mathcal{V}_{\mathcal{B}}$  of BTZ microstates defined by

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- ▶ Agrees with Bekenstein–Hawking and Cardy formula

# Outline

Motivation

Near horizon boundary conditions

Explicit construction of BTZ microstates

Discussion

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- ▶ Highly non-trivial indication for existence of soft Heisenberg hair in 4d

## Microstates of non-extremal Kerr?



Main challenge: how to provide (controlled) cut-off on soft hair spectrum in four dimensions?

Thanks for your attention!



-  H. Afshar, D. Grumiller and M.M. Sheikh-Jabbari “Near Horizon Soft Hairs as Microstates of Three Dimensional Black Holes,” 1607.00009.
-  H. Afshar, S. Detournay, D. Grumiller, W. Merbis, A. Perez, D. Tempo and R. Troncoso “Soft Heisenberg hair on black holes in three dimensions,” Phys.Rev. **D93** (2016) 101503(R); 1603.04824.

Thanks to Bob McNees for providing the  $\LaTeX$  beamerclass!

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- ▶ Spectral flow and discrete conic spaces generated by  $\mathcal{J}_r^\pm$  ( $r = 1, 2, \dots, c-1$ ), the “horizon fluffs”

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- ▶ Mismatch in coefficients; not sure yet if bug or feature