Carrollian c-functions and flat space holographic RG flows

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with Max Riegler, 2309.11539



Outline

Flat space holography à la Carroll

Holographic $\mathit{c}\text{-functions}$ in $\mathsf{AdS}_3/\mathsf{CFT}_2$

Flat space domain walls and BMS_3 *c*-functions

Outlook to Casini-Huerta-like c-function

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Holographic c-functions in AdS₃/CFT₂

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Two main approaches:

- Carrollian (codimension-1 holography)
- Celestial (codimension-2 holography)

Translation between Carrollian and Celestial holography possible

Donnay, Fiorucci, Herfray, Ruzziconi; Bagchi, Banerjee, Basu, Dutta '22

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This talk:

Focus on $\mathsf{BMS}_3/\mathsf{CCFT}_2$ correspondence

We want holography to make useful statements about

1. symmetries

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- 2. spectra

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- 3. microstates

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- 5. correlation functions

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Review five examples

1. analog of Brown–Henneaux bc's leading to BMS₃ Barnich, Compère '06 $ds^{2} = \left(\mathcal{O}(1) \ du^{2} - 2 \ du \ dr + \mathcal{O}(1) \ du \ d\varphi + r^{2} \ d\varphi^{2}\right) \left(1 + \mathcal{O}(1/r)\right)$ preserved by asymptotic Killing vectors

$$M_n = ie^{in\varphi}\partial_u + \dots \qquad L_n = ie^{in\varphi}\left(inu\partial_u - inr\partial_r + (1 + \frac{u}{r}n^2)\partial_\varphi\right) + \dots$$

whose Lie bracket algebra generates superrotations & -translations

$$[L_n, L_m] = (n - m) L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m,0}$$
$$[L_n, M_m] = (n - m) M_{n+m} + \frac{c_M}{12} (n^3 - n) \delta_{n+m,0}$$
$$[M_n, M_m] = 0$$

canonical realization of asymptotic symmetries can lead to central extensions, depending on the bulk theory

algebra above is Carrollian conformal algebra Duval, Gibbons, Horvathy '14 and isomorphic to Galilean conformal algebra in 2d Bagchi '10

- 1. analog of Brown–Henneaux bc's leading to BMS_3 Barnich, Compère '06
- 2. match of linearized spectra & gap in spectrum Bagchi, Detournay, DG '12

e.g. in flat space chiral gravity

$$\operatorname{CCFT}_2: \psi^{(n)} = L_{-n} |0\rangle$$
 $\operatorname{BMS}_3: \psi^{(n)}_{uu} = -2ne^{-in\varphi} \dots$

Virsoro descendants of vacuum mapped to "boundary gravitons"

gap: ground state (global Minkowski vacuum) has Virasoro charge $Q_{L_0} = -k$ while non-perturbative states (flat space cosmologies) have Virasoro charge $Q_{L_0} = k \alpha^2$ global Minkowski:

$$-\,\mathrm{d} u^2 - 2\,\mathrm{d} u\,\mathrm{d} r + r^2\,\mathrm{d} \varphi^2$$

flat space cosmologies (with horizon radius r_0):

$$\alpha^2 \left(1 - \frac{r_0^2}{r^2}\right) \,\mathrm{d}u^2 - 2 \,\mathrm{d}u \,\mathrm{d}r + r^2 \left(\,\mathrm{d}\varphi - \frac{\alpha r_0}{r^2} \,\mathrm{d}u\right)^2$$

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- 3. microstates à la Cardy Barnich; Bagchi, Detournay, Fareghbal, Simón '12

thermal entropy of flat space cosmologies:

$$S_{
m macro} = S_{
m FSC} = S_{
m BH} = rac{2\pi r_0}{4G}$$

thermal entropy in Carrollian CFT_2 using Cardyology:

$$S_{\rm micro} = S_{\rm CCFT_2} = 2\pi h_L \sqrt{\frac{c_M}{24h_M}}$$

fineprint: assumed here vanishing Virasoro central charge and non-vanishing BMS central charge, like in Einstein gravity

explicit computation shows
$$h_L = \frac{\alpha r_0}{4G}$$
, $h_M = \frac{\alpha^2}{8G}$, $c_M = \frac{3}{G}$ and thus $S_{\rm macro} = S_{\rm micro}$

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- 4. (holographic) entanglement entropy Bagchi, Basu, DG, Riegler '14

Carrollian CFT $_2$ calculation yields EE for vacuum on the plane

$$S_{\rm EE} = \frac{c_L}{6} \ln \frac{\Delta x}{\varepsilon_x} + \frac{c_M}{6} \left(\frac{\Delta u}{\Delta x} - \frac{\varepsilon_u}{\varepsilon_x} \right)$$

get also EE for other states related to vacuum by uniformization DG, Parekh, Riegler '19 another example: EE for vacuum on the cylinder

$$S_{\rm EE} = \frac{c_L}{6} \ln \frac{2\sin\frac{\Delta\varphi}{2}}{\varepsilon_{\varphi}} + \frac{c_M}{6} \left(\frac{\Delta u}{2} \cot\frac{\Delta\varphi}{2} - \frac{\varepsilon_u}{\varepsilon_{\varphi}}\right)$$

results reproduced on gravity side using Wilson lines Basu, Riegler '15 or swing construction with geodesics Jiang, Song, Wen '17; Apolo, Jiang, Song, Zhong '20

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- 5. (holographic) stress tensor correlation functions Bagchi, DG, Merbis '15

CCFT₂ conservation equations (M: supertranslations, L: superrotations)

$$\partial_u \langle M \mathcal{O} \rangle = 0 \qquad \qquad \partial_u \langle L \mathcal{O} \rangle = \partial_\varphi \langle M \mathcal{O} \rangle$$

yield BPZ-like recursion relations $(s_{ij} = 2\sin[(\varphi_1 - \varphi_2)/2], c_{ij} = \cot[(\varphi_1 - \varphi_2)/2])$

$$\langle M_1 L_2 \dots L_n \rangle = \sum_{i=2}^n \left(\frac{2}{s_{1i}^2} + \frac{c_{1i}}{2} \,\partial_{\varphi_i} \right) \langle M_2 L_3 \dots L_n \rangle$$
$$\langle L_1 L_2 \dots L_n \rangle = \frac{c_L}{c_M} \left\langle M_1 L_2 \dots L_n \right\rangle + \sum_{i=1}^n u_i \partial_{\varphi_i} \left\langle M_1 L_2 \dots L_n \right\rangle$$

reproduced on gravity side (in Chern-Simons formulation)

Wout Merbis

Institute for Theoretical Physics, Vienna University of Tachnology Wedner Hauptstrasse 8-10/138 A-1040 Wien, AUSTRIA wei microsoften ph. Univ. John Statusch



 $(T(z_1)T(z_2)T(z_3)T(z_4)T(z_5)T(z_6)T(z_7)) =$

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Bessel or work presented in "Stress lensor considers in three-dimensional gravity" x - 21 at 1677, 201823 with A. Begchi and D. Gramilian. Special thanks to Friedrich Schölter for Mathematica visandry

Asymptotic symmetries in asymptotically flat space for 3d Einstein gravity: Ashtekar, Bicak, Schmidt '96; Barnich, Compère '06

$$[L_n, L_m] = (n - m) L_{n+m}$$
$$[L_n, M_m] = (n - m) M_{n+m} + \frac{c_M}{12} n(n^2 - 1) \delta_{n+m,0}$$
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numerous holographic checks & results based on BMS₃ symmetries

- flat space chiral gravity Bagchi, Detournay, DG '12
- cardyology Bagchi, Detournay, Fareghbal, Simón; Barnich '12
- phase transitions Bagchi, Detournay, DG, Simón '13
- entanglement entropy Bagchi, Basu, DG, Riegler '14; Jiang, Song, Wen '17
- holographic dictionary & 1-point fct.'s Detournay, DG, Schöller, Simón '14
- 1-loop partition fct. & BMS characters Barnich, González, Maloney, Oblak '15
- all stress tensor correlators Bagchi, DG, Merbis '15
- BMS bootstrap Bagchi, Gary, Zodinmawia '16
- BMS blocks Hijano '18

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- L_n : superrotations (diff (S^1))
- ► M_n: supertranslations
- $c_M = 3/G$: BMS₃-central charge

Note: c_M dimensionful \Rightarrow dimensionless ratios still meaningful

$$rac{c_M}{h_M}\,, \qquad \quad rac{c_M^{
m UV}}{c_M^{
m R}}\,, \qquad \quad c_M imes {
m length}$$

where

$$M_0|\Psi\rangle = h_M|\Psi\rangle$$

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- $c_M = 3/G$: BMS₃-central charge
- c_L: non-zero for TMG, but zero in Einstein gravity
- ► BMS₃ emerges as UR limit of CFT₂ symmetries ⇒ Carrollian CFT₂

$$L_n := \mathcal{L}_n^+ - \mathcal{L}_{-n}^- \qquad \qquad M_n := \frac{1}{\ell} \left(\mathcal{L}_n^+ + \mathcal{L}_{-n}^- \right)$$

with Virasoros $[\mathcal{L}_n^{\pm}, \mathcal{L}_m^{\pm}] = (n-m) \mathcal{L}_{n+m}^{\pm} + \frac{c^{\pm}}{12} (n^3 - n) \delta_{n+m,0}$ yields CCFT₂ \simeq BMS₃ with $c_L = c^+ - c^-$ and $c_M = \frac{1}{\ell} (c^+ + c^-)$

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- ► *M_n*: supertranslations
- $c_M = 3/G$: BMS₃-central charge
- ▶ *c*_{*L*}: non-zero for TMG, but zero in Einstein gravity
- ▶ BMS₃ emerges as UR limit of CFT₂ symmetries \Rightarrow Carrollian CFT₂

This talk: bulk theory = Einstein gravity + scalar field

$$I_{\text{bulk}}[g_{\mu\nu}, \phi] = \frac{1}{16\pi G} \int d^3x \sqrt{-g} \left(R - \frac{1}{2} \left(\partial \phi \right)^2 - V(\phi) \right)$$

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Flat space holography à la Carroll

Holographic $\mathit{c}\text{-functions}$ in $\mathsf{AdS}_3/\mathsf{CFT}_2$

Flat space domain walls and BMS_3 c-functions

Outlook to Casini–Huerta-like *c*-function

Consider RG-flow from UV to IR in QFT_2

fineprint: the Euclidean QFT2 is assumed to be renormalizable, reflection-positive, translation- and rotation-invariant

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 \blacktriangleright coupling constants g^i have β -functions

$$\dot{g}^i = \beta^i(g)$$

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$$\dot{g}^i = \beta^i(g)$$

there exists a function c(g) with the properties 1. Monotonicity

$$\dot{c}(g) := \beta^i(g) \, \frac{\partial c(g)}{\partial g^i} \le 0$$

Equivalently: c(g) non-increasing under dilatations!

Consider RG-flow from UV to IR in QFT_2

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2. Saturation at fixed points

$$\dot{c}(g_*) = 0 \qquad \leftrightarrow \qquad \beta^i(g_*) = 0$$

Note: at fixed points enhancement to CFT₂ Virasoro symmetries

$$[L_n, L_m] = (n-m) L_{n+m} + \frac{c_{\text{Vir}}(g_*)}{12} n(n^2 - 1) \delta_{n+m,0}$$

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3. Equality to Virasoro central charge

$$c(g_*) = c_{\rm Vir}(g_*)$$
Zamolodchikov *c*-function

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Consequence:

$$c_{\rm Vir}^{\rm UV} > c_{\rm Vir}^{\rm IR} \qquad \Rightarrow \qquad {\rm more \ dof \ in \ UV}$$

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Explicit construction of a c-function using stress tensor and its 2-point correlators Zamolodchikov '86

domain walls in AdS_3 (set AdS radius to unity)

$$\mathrm{d}s^2 = \mathrm{d}\rho^2 + e^{2A(\rho)} \left(-\mathrm{d}t^2 + \mathrm{d}x^2 \right) \qquad \lim_{\rho \to \infty} A(\rho) = \rho + \dots$$

holographic model for RG flow (UV: $\rho \rightarrow \infty$)

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 $\blacktriangleright \rho = \text{const. slices: Poincaré invariant}$

 $KVs: \quad \partial_t \qquad \partial_x \qquad x\partial_t + t\partial_x$ conformal KVs: infinitely many (CFT₂ symmetries)

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holographic model for RG flow (UV: $\rho \rightarrow \infty)$

• $\rho = \text{const. slices: Poincaré invariant}$

KVs: ∂_t ∂_x $x\partial_t + t\partial_x$

conformal KVs: infinitely many (CFT₂ symmetries)

holographic domain wall c-function

$$c_{
m dw}(
ho) = rac{c^{
m UV}}{A'(
ho)}$$

monotonicity implied by reality of bulk scalar field [see later!]





- SSA of EE
 - $S(C) + S(B) \ge S(A) + S(D)$

Lieb, Ruskai '73; Kiefer '59



- SSA of EE
 - $S(C) + S(B) \ge S(A) + S(D)$
- Use Minkowski diagram
 - $S(B) S(A) \geq S(\lambda B) S(\lambda A)$



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Differential instead difference

 $c_{\rm CH}(L) = \# L S'(L)$

non-increasing under dilatations!



- SSA of EE
 - $S(\mathbb{C}) + S(\mathbb{B}) \ge S(A) + S(D)$
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Differential instead difference

 $c_{\rm CH}(L) = \# L S'(L)$

non-increasing under dilatations!Fix normalization

$$c_{\rm CH}(L) = 3 \, L \, S'(L)$$

Properties of CH c-function

Monotonicity

$$c'_{\rm CH}(L) = 3L \, S''(L) + 3S'(L) \le 0$$

Properties of CH c-function

Monotonicity

$$c_{\rm CH}'(L) = 3L\,S''(L) + 3S'(L) \le 0$$

Fixed point values

$$\lim_{L \to 0} c_{\rm CH}(L) = c^{\rm UV} \qquad \qquad \lim_{L \to \infty} c_{\rm CH}(L) = c^{\rm IR}$$

Note: at fixed points we have

$$\begin{split} S(L \to 0) &= \frac{c^{\text{UV}}}{3} \ln L \quad \Rightarrow \quad c_{\text{CH}}(L \to 0) = \lim_{L \to 0} \left(3L \, S'(L) \right) \to c^{\text{UV}} \\ S(L \to \infty) &= \frac{c^{\text{IR}}}{3} \ln L \quad \Rightarrow \quad c_{\text{CH}}(L \to \infty) = \lim_{L \to \infty} \left(3L \, S'(L) \right) \to c^{\text{IR}} \\ \text{Holzbey Larsen Wilczek '94: Cardy Calabrese '06} \end{split}$$

 \Rightarrow is indeed a *c*-function different from previous ones!

Properties of CH c-function

Monotonicity

$$c'_{\rm CH}(L) = 3L \, S''(L) + 3S'(L) \le 0$$

Fixed point values

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Monotonicity of CH c-function

$$\frac{c_{\rm CH}'(L)}{3L} = S''(L) - \frac{S'(L)}{L} + \frac{6}{c_{\rm CH}(L)} (S'(L))^2 \le 0$$

implies ground state QNEC [see next slide]

$$S''(L) - \frac{S'(L)}{L} + \frac{6}{c_{\rm UV}} (S'(L))^2 \le 0$$

Ecker, DG, Soltanpanahi, Stanzer '20

$\mathsf{QNEC} \text{ in } \mathsf{2d}$

Quantum Null Energy Condition (QNEC) in 2d:

$$2\pi \left\langle T_{\pm\pm} \right\rangle \ge \frac{\mathrm{d}^2 S}{\mathrm{d}\lambda^2} \bigg|_{\lambda=0} + \frac{6}{c^{\mathrm{UV}}} \left(\frac{\mathrm{d}S}{\mathrm{d}\lambda}\right)^2 \bigg|_{\lambda=0}$$



Bousso, Fisher, Leichenauer, Wall '15

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Bousso, Fisher, Leichenauer, Wall '15

both sides of QNEC transform with Schwarzian derivative Wall '11

under bulk diffeos (or boundary conformal trafos):

$$\delta_{\xi}S = \underbrace{\xi S' - \frac{c}{12}\,\xi'}_{}$$

anomalous scalar

implies

$$\delta_{\xi}Q = \underbrace{\xi Q' + 2\xi' Q - \frac{c}{12}\xi'''}_{\text{infinitesimal Schwarzian}}$$

with
$$Q:=S''+\frac{6}{c}\,(S')^2$$

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- QNEC saturates for states dual to vacuum Einstein solutions
 Khandker, Kundu, Li '18; Ecker, DG, Sheikh-Jabbari, Stanzer, van der Schee '20

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- both sides of QNEC transform with Schwarzian derivative Wall '11
- QNEC saturates for states dual to vacuum Einstein solutions Khandker, Kundu, Li '18; Ecker, DG, Sheikh-Jabbari, Stanzer, van der Schee '20
- for boost invariant ground states:

$$0 \ge \frac{\mathrm{d}^2 S}{\mathrm{d}\lambda^2} \bigg|_{\lambda=0} + \frac{6}{c^{\mathrm{UV}}} \left(\frac{\mathrm{d}S}{\mathrm{d}\lambda}\right)^2 \bigg|_{\lambda=0} = S''(L) - \frac{S'(L)}{L} + \frac{6}{c^{\mathrm{UV}}} \left(S'\right)^2$$

since $S(\lambda, L + \lambda) = S(0, \sqrt{(L + \lambda)^2 - \lambda^2})$ Ecker, DG, Sheikh-Jabbari, Stanzer '20

Quantum Null Energy Condition (QNEC) in 2d:

$$2\pi \left\langle T_{\pm\pm} \right\rangle \geq \frac{\mathrm{d}^2 S}{\mathrm{d}\lambda^2} \bigg|_{\lambda=0} + \frac{6}{c^{\mathrm{UV}}} \left(\frac{\mathrm{d}S}{\mathrm{d}\lambda}\right)^2 \bigg|_{\lambda=0}$$

Bousso, Fisher, Leichenauer, Wall '15

- both sides of QNEC transform with Schwarzian derivative Wall '11
- QNEC saturates for states dual to vacuum Einstein solutions
 Khandker, Kundu, Li '18; Ecker, DG, Sheikh-Jabbari, Stanzer, van der Schee '20
- for boost invariant ground states:

$$0 \ge \frac{\mathrm{d}^2 S}{\mathrm{d}\lambda^2} \bigg|_{\lambda=0} + \frac{6}{c^{\mathrm{UV}}} \left(\frac{\mathrm{d}S}{\mathrm{d}\lambda}\right)^2 \bigg|_{\lambda=0} = S''(L) - \frac{S'(L)}{L} + \frac{6}{c^{\mathrm{UV}}} \left(S'\right)^2$$

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ground state QNEC necessary for monotonicity of CH c-function

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Flat space holography à la Carroll

Holographic c-functions in AdS_3/CFT_2

Flat space domain walls and BMS_3 c-functions

Outlook to Casini–Huerta-like *c*-function

$$ds^{2} = -e^{A(r)} 2 du dr + e^{2A(r)} dx^{2}$$

as solutions of Einstein-Klein-Gordon bulk theory

$$I_{\text{bulk}}[g_{\mu\nu}, \phi] = \frac{1}{16\pi G} \int d^3x \sqrt{-g} \left(R - \frac{1}{2} (\partial \phi)^2 - V(\phi) \right)$$

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 r = r₀ = const. slices have degenerate metric 0 du² + 0 du dx + e^{2A(r₀)} dx² whose conformal KVs form BMS₃ conformal KVs: (ξ_M(x) + u ξ'_L(x)) ∂_u + ξ_L(x) ∂_x ξ_L(x): generates diff(S¹) ξ_M(x): generates supertranslations

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can translate all AdS₃ domain walls into flat space domain walls!

Is holographic domain wall function

$$c_{
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 $\mathit{c}\text{-function}$ for Carrollian QFTs with BMS_3 invariant UV & IR fixed points?

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all desired properties of flat space holographic c-function

Take flat space domain wall related to AdS domain wall with super-potential $W(\phi) = -2 - \frac{1}{4} \phi^2 - \frac{\alpha}{8} \phi^4$ (mass $m^2 = -\frac{3}{4}$ above BF)

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yields profile function

$$A(r) = r + A_0 - \frac{j^2}{16} \left(\frac{1}{e^r - \alpha j^2} + \frac{r - \ln(e^r - \alpha j^2)}{\alpha j^2} \right)$$

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can read off BMS₃ central charges at fixed points

$$c_M^{\rm IR} = \frac{c_M^{\rm UV}}{1 - \frac{1}{16\alpha}} < c_M^{\rm UV}$$

more info in plot of flat space domain wall c-function
Domain wall example

Take flat space domain wall related to AdS domain wall with super-potential $W(\phi) = -2 - \frac{1}{4}\phi^2 - \frac{\alpha}{8}\phi^4$ (mass $m^2 = -\frac{3}{4}$ above BF) $C_{dw}(r_0)$ 1.0 0.8 0.6 $C_{dw}(r_0)$ 0.4 0.2 r_0 2 -8 -6 -2 0 4 _4

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▶ split EE into *L*- and *M*-parts:

$$S_{\text{EE}} = S_L + S_M$$
 $S_L = \frac{c_L}{6} \ln \frac{\Delta x}{\epsilon_x}$ $S_M = \frac{c_M}{6} \left(\frac{\Delta u}{\Delta x} - \frac{\epsilon_u}{\epsilon_x}\right)$

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QECs at our disposal — Casini–Huerta-like *c*-function conceivable!

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Monotonicity of conjectured c-functions yields ground state QEC

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All results in 3 bulk dimensions. Do not know how/if all results generalize to other dimensions.

Ladies and gentlemen, esteemed colleagues, and fellow enthusiasts of the Carrollian cosmos. As we prepare to wrap up our journey through the whimsical wonderland of "Carrollian c-functions and flat space holographic RG flows,'' I can't help but feel a bit like Alice herself, exploring the curious corners of spacetime where BMS symmetries, Carrollian CFTs, and c-functions come together in a harmonious dance. Our adventure today has taken us down the rabbit hole of flat space holography. But as we prepare to return from this captivating journey, let us remember that in science, as in Wonderland, curiosity is our guiding star. As Carrollian characters once said, "Begin at the beginning, and go on till you come to the end: then stop.''

My conclusion

Carrollian holography is a lot of fun - feel free to join!



Al prompt: Carroll and Zamolodchikov meet 't Hooft and Susskind at the rabbit hole to wonderland