

Flat space higher spin gravity

Daniel Grumiller

Institute for Theoretical Physics
TU Wien

Dutch String meeting, Groningen, February 2015



based on work w. Afshar, Bagchi, Basu, Detournay, Fareghbal,
Gary, Riegler, Rosseel, Salzer, Sarkar, Schöller, Simon, ...

Outline

Motivations

Holography basics

Flat space gravity

Flat space higher spin gravity

Outline

Motivations

Holography basics

Flat space gravity

Flat space higher spin gravity

General motivations

- ▶ Quantum gravity
 - ▶ Address conceptual issues of quantum gravity



General motivations

- ▶ Quantum gravity
 - ▶ Address conceptual issues of quantum gravity
 - ▶ Black holes (thermodynamics, evaporation, information loss, microstate counting, entanglement entropy, firewalls, ...)



General motivations

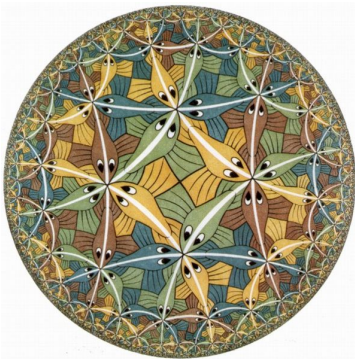
▶ Quantum gravity

- ▶ Address conceptual issues of quantum gravity
- ▶ Black holes (thermodynamics, evaporation, information loss, microstate counting, entanglement entropy, firewalls, ...)
- ▶ String theory (is it the right theory? can there be any alternative? ...)



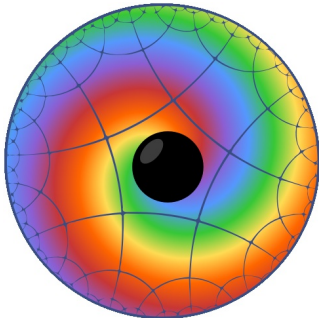
General motivations

- ▶ Quantum gravity
 - ▶ Address conceptual issues of quantum gravity
 - ▶ Black holes (thermodynamics, evaporation, information loss, microstate counting, entanglement entropy, firewalls, ...)
 - ▶ String theory (is it the right theory? can there be any alternative? ...)
- ▶ Holography
 - ▶ Holographic principle realized in Nature? (yes/no)



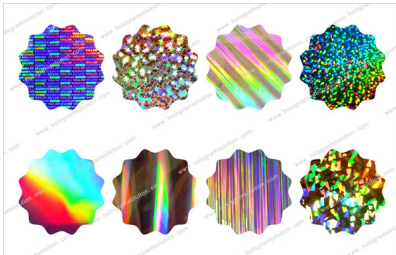
General motivations

- ▶ Quantum gravity
 - ▶ Address conceptual issues of quantum gravity
 - ▶ Black holes (thermodynamics, evaporation, information loss, microstate counting, entanglement entropy, firewalls, ...)
 - ▶ String theory (is it the right theory? can there be any alternative? ...)
- ▶ Holography
 - ▶ Holographic principle realized in Nature? (yes/no)
 - ▶ Quantum gravity via AdS/CFT? (define quantum gravity in AdS by constructing/postulating dual CFT)



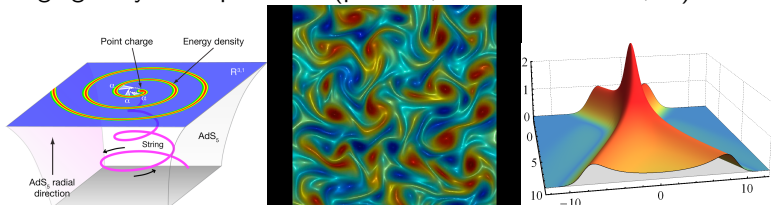
General motivations

- ▶ Quantum gravity
 - ▶ Address conceptual issues of quantum gravity
 - ▶ Black holes (thermodynamics, evaporation, information loss, microstate counting, entanglement entropy, firewalls, ...)
 - ▶ String theory (is it the right theory? can there be any alternative? ...)
- ▶ Holography
 - ▶ Holographic principle realized in Nature? (yes/no)
 - ▶ Quantum gravity via AdS/CFT? (define quantum gravity in AdS by constructing/postulating dual CFT)
 - ▶ How general is holography? (non-unitary holography, **higher spin holography**, **flat space holography**, non-AdS holography, ...)



General motivations

- ▶ Quantum gravity
 - ▶ Address conceptual issues of quantum gravity
 - ▶ Black holes (thermodynamics, evaporation, information loss, microstate counting, entanglement entropy, firewalls, ...)
 - ▶ String theory (is it the right theory? can there be any alternative? ...)
- ▶ Holography
 - ▶ Holographic principle realized in Nature? (yes/no)
 - ▶ Quantum gravity via AdS/CFT? (define quantum gravity in AdS by constructing/postulating dual CFT)
 - ▶ How general is holography? (non-unitary holography, **higher spin holography**, **flat space holography**, non-AdS holography, ...)
- ▶ Applications
 - ▶ Gauge gravity correspondence (plasmas, condensed matter, ...)



Specific motivation for 3D

Gravity in 3D is simpler than in higher dimensions

Specific motivation for 3D

Gravity in 3D is simpler than in higher dimensions

A quote by Albert Einstein: "Simplicity is the ultimate sophistication". The text is centered on a light gray background with a subtle gradient and a faint diagonal line.

Specific motivation for flat space higher spin gravity

Massless higher spin theories constrained by no-gos!

- ▶ Coleman–Mandula '67
- ▶ Aragone–Deser '79
- ▶ Weinberg–Witten '80
- ▶ recent summary: [Bekaert, Boulanger, Sundell '12](#)

Conclusion: there are no consistent interacting massless higher spin theories in 4- (or higher-) dimensional flat space

Specific motivation for flat space higher spin gravity

Massless higher spin theories constrained by no-gos!

- ▶ Coleman–Mandula '67
- ▶ Aragone–Deser '79
- ▶ Weinberg–Witten '80
- ▶ recent summary: Bekaert, Boulanger, Sundell '12

Conclusion: there are no consistent interacting massless higher spin theories in 4- (or higher-) dimensional flat space

Circumventing no-gos:

- ▶ Vasiliev '87-'90: higher spin theories in (A)dS

Specific motivation for flat space higher spin gravity

Massless higher spin theories constrained by no-gos!

- ▶ Coleman–Mandula '67
- ▶ Aragone–Deser '79
- ▶ Weinberg–Witten '80
- ▶ recent summary: [Bekaert, Boulanger, Sundell '12](#)

Conclusion: there are no consistent interacting massless higher spin theories in 4- (or higher-) dimensional flat space

Circumventing no-gos:

- ▶ [Vasiliev '87-'90](#): higher spin theories in (A)dS
- ▶ [Afshar, Bagchi, Fareghbal, DG, Rosseel '13](#); [Gonzalez, Matulich, Pino, Troncoso '13](#): flat space higher spin theories in 3d

Goals of this talk

1. Review general aspects of holography in 3D

Goals of this talk

1. Review general aspects of holography in 3D
2. Discuss flat space holography

Goals of this talk

1. Review general aspects of holography in 3D
2. Discuss flat space holography
3. Generalize to higher spin holography

Goals of this talk

1. Review general aspects of holography in 3D
2. Discuss flat space holography
3. Generalize to higher spin holography
4. List selected open issues

Goals of this talk

1. Review general aspects of holography in 3D
2. Discuss flat space holography
3. Generalize to higher spin holography
4. List selected open issues

Address these issues in 3D!



Outline

Motivations

Holography basics

Flat space gravity

Flat space higher spin gravity

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons
- ▶ Weyl = 0, thus Riemann = Ricci

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons
- ▶ Weyl = 0, thus Riemann = Ricci
- ▶ Einstein gravity: no on-shell gravitons

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons
- ▶ Weyl = 0, thus Riemann = Ricci
- ▶ Einstein gravity: no on-shell gravitons
- ▶ Formulation as topological gauge theory (Chern–Simons)

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons
- ▶ Weyl = 0, thus Riemann = Ricci
- ▶ Einstein gravity: no on-shell gravitons
- ▶ Formulation as topological gauge theory (Chern–Simons)
- ▶ Dual field theory (if it exists): 2D

Interesting generic constraints from CFT₂!

e.g. [Hellerman '09](#), [Hartman, Keller, Stoica '14](#)

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons
- ▶ Weyl = 0, thus Riemann = Ricci
- ▶ Einstein gravity: no on-shell gravitons
- ▶ Formulation as topological gauge theory (Chern–Simons)
- ▶ Dual field theory (if it exists): 2D
- ▶ Infinite dimensional asymptotic symmetries (Brown–Henneaux)

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons
- ▶ Weyl = 0, thus Riemann = Ricci
- ▶ Einstein gravity: no on-shell gravitons
- ▶ Formulation as topological gauge theory (Chern–Simons)
- ▶ Dual field theory (if it exists): 2D
- ▶ Infinite dimensional asymptotic symmetries (Brown–Henneaux)
- ▶ Black holes as orbifolds of AdS₃ (BTZ)

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons
- ▶ Weyl = 0, thus Riemann = Ricci
- ▶ Einstein gravity: no on-shell gravitons
- ▶ Formulation as topological gauge theory (Chern–Simons)
- ▶ Dual field theory (if it exists): 2D
- ▶ Infinite dimensional asymptotic symmetries (Brown–Henneaux)
- ▶ Black holes as orbifolds of AdS₃ (BTZ)
- ▶ Simple microstate counting from AdS₃/CFT₂

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons
- ▶ Weyl = 0, thus Riemann = Ricci
- ▶ Einstein gravity: no on-shell gravitons
- ▶ Formulation as topological gauge theory (Chern–Simons)
- ▶ Dual field theory (if it exists): 2D
- ▶ Infinite dimensional asymptotic symmetries (Brown–Henneaux)
- ▶ Black holes as orbifolds of AdS₃ (BTZ)
- ▶ Simple microstate counting from AdS₃/CFT₂
- ▶ Hawking–Page phase transition hot AdS ↔ BTZ

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons
- ▶ Weyl = 0, thus Riemann = Ricci
- ▶ Einstein gravity: no on-shell gravitons
- ▶ Formulation as topological gauge theory (Chern–Simons)
- ▶ Dual field theory (if it exists): 2D
- ▶ Infinite dimensional asymptotic symmetries (Brown–Henneaux)
- ▶ Black holes as orbifolds of AdS₃ (BTZ)
- ▶ Simple microstate counting from AdS₃/CFT₂
- ▶ Hawking–Page phase transition hot AdS ↔ BTZ
- ▶ Simple checks of Ryu–Takayanagi proposal

Gravity in 3D

AdS₃ gravity

- ▶ Lowest dimension with black holes and (off-shell) gravitons
- ▶ Weyl = 0, thus Riemann = Ricci
- ▶ Einstein gravity: no on-shell gravitons
- ▶ Formulation as topological gauge theory (Chern–Simons)
- ▶ Dual field theory (if it exists): 2D
- ▶ Infinite dimensional asymptotic symmetries (Brown–Henneaux)
- ▶ Black holes as orbifolds of AdS₃ (BTZ)
- ▶ Simple microstate counting from AdS₃/CFT₂
- ▶ Hawking–Page phase transition hot AdS ↔ BTZ
- ▶ Simple checks of Ryu–Takayanagi proposal

Caveat: while there are many string compactifications with AdS₃ factor, applying holography just to AdS₃ factor does not capture everything!

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle

Example: Einstein gravity with Dirichlet boundary conditions

$$I = -\frac{1}{16\pi G_N} \int d^3x \sqrt{|g|} \left(R + \frac{2}{\ell^2} \right)$$

with $\delta g = \text{fixed}$ at the boundary

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions

Example: asymptotically AdS

$$ds^2 = d\rho^2 + \left(e^{2\rho/\ell} \gamma_{ij}^{(0)} + \gamma_{ij}^{(2)} + \dots \right) dx^i dx^j$$

with $\delta\gamma^{(0)} = 0$ for $\rho \rightarrow \infty$

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions
3. Perform canonical analysis and check consistency of bc's
 - ▶ Find and classify all constraints
 - ▶ Construct canonical gauge generators
 - ▶ Add boundary terms and get (variation of) canonical charges
 - ▶ Check integrability of canonical charges
 - ▶ Check finiteness of canonical charges
 - ▶ Check conservation (in time) of canonical charges
 - ▶ Calculate Dirac bracket algebra of canonical charges

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions
3. Perform canonical analysis and check consistency of bc's
 - ▶ Find and classify all constraints
 - ▶ Construct canonical gauge generators
 - ▶ Add boundary terms and get (variation of) canonical charges
 - ▶ Check integrability of canonical charges
 - ▶ Check finiteness of canonical charges
 - ▶ Check conservation (in time) of canonical charges
 - ▶ Calculate Dirac bracket algebra of canonical charges

Example: Brown–Henneaux analysis for 3D Einstein gravity

$$\{Q[\varepsilon], Q[\eta]\} = \delta_\varepsilon Q[\eta]$$

$$Q[\varepsilon] \sim \oint d\varphi \mathcal{L}(\varphi) \varepsilon(\varphi)$$

$$\delta_\varepsilon \mathcal{L} = \mathcal{L} \varepsilon + 2\mathcal{L} \varepsilon' + \frac{\ell}{16\pi G_N} \varepsilon'''$$

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions
3. Perform canonical analysis and check consistency of bc's
4. Derive (classical) asymptotic symmetry algebra and central charges

Example: Two copies of Virasoro algebra

$$[\mathcal{L}_n, \mathcal{L}_m] = (n - m) \mathcal{L}_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0}$$

with Brown–Henneaux central charge

$$c = \frac{3\ell}{2G_N}$$

Reminder: ASA = quotient algebra of asymptotic symmetries by 'trivial' asymptotic symmetries with zero canonical charges

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions
3. Perform canonical analysis and check consistency of bc's
4. Derive (classical) asymptotic symmetry algebra and central charges
5. Improve to quantum ASA

Example: semi-classical ASA in spin-3 gravity (Henneaux, Rey '10; Campoleoni, Pfenninger, Fredenhagen, Theisen '10)

$$[W_n, W_m] = \frac{16}{5c} \sum_p L_p L_{n+m-p} + \dots$$

quantum ASA

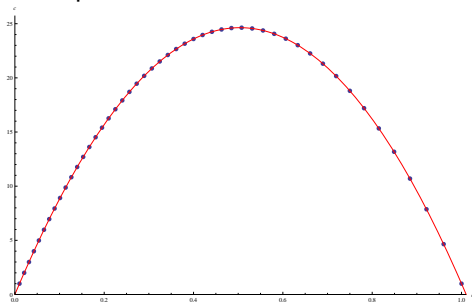
$$[W_n, W_m] = \frac{16}{5c + 22} \sum_p : L_p L_{n+m-p} : + \dots$$

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions
3. Perform canonical analysis and check consistency of bc's
4. Derive (classical) asymptotic symmetry algebra and central charges
5. Improve to quantum ASA
6. Study unitary representations of quantum ASA

Example:



Afshar et al '12

Discrete set of Newton
constant values compatible
with unitarity
(3D spin-N gravity in
next-to-principal embedding)

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions
3. Perform canonical analysis and check consistency of bc's
4. Derive (classical) asymptotic symmetry algebra and central charges
5. Improve to quantum ASA
6. Study unitary representations of quantum ASA
7. Identify/constrain dual field theory

Example: Monster CFT in (flat space) chiral gravity

Witten '07

Li, Song & Strominger '08

Bagchi, Detournay & DG '12

$$Z(q) = J(q) = \frac{1}{q} + (1 + 196883)q + \mathcal{O}(q^2)$$

Note: $\ln 196883 \approx 12.2 = 4\pi + \text{quantum corrections}$

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions
3. Perform canonical analysis and check consistency of bc's
4. Derive (classical) asymptotic symmetry algebra and central charges
5. Improve to quantum ASA
6. Study unitary representations of quantum ASA
7. Identify/constrain dual field theory
8. If unhappy with result go back to previous items and modify

Examples: too many!



Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions
3. Perform canonical analysis and check consistency of bc's
4. Derive (classical) asymptotic symmetry algebra and central charges
5. Improve to quantum ASA
6. Study unitary representations of quantum ASA
7. Identify/constrain dual field theory
8. If unhappy with result go back to previous items and modify

Goal of this talk:

Apply algorithm above to flat space holography in 3D higher spin theories

Outline

Motivations

Holography basics

Flat space gravity

Flat space higher spin gravity

Flat space holography (Barnich et al, Bagchi et al, Strominger et al, ...)

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

Flat space holography (Barnich et al, Bagchi et al, Strominger et al, ...)

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

- ▶ Works straightforwardly sometimes, otherwise not

Flat space holography (Barnich et al, Bagchi et al, Strominger et al, ...)

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

- ▶ Works straightforwardly sometimes, otherwise not
- ▶ Example where it works nicely: asymptotic symmetry algebra

Flat space holography (Barnich et al, Bagchi et al, Strominger et al, ...)

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

- ▶ Works straightforwardly sometimes, otherwise not
- ▶ Example where it works nicely: asymptotic symmetry algebra
- ▶ Take linear combinations of Virasoro generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$

$$L_n = \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n = \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

- ▶ Works straightforwardly sometimes, otherwise not
- ▶ Example where it works nicely: asymptotic symmetry algebra
- ▶ Take linear combinations of Virasoro generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$

$$L_n = \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n = \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

- ▶ Make Inönü–Wigner contraction $\ell \rightarrow \infty$ on ASA

$$[L_n, L_m] = (n - m) L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[L_n, M_m] = (n - m) M_{n+m} + \frac{c_M}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[M_n, M_m] = 0$$

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

- ▶ Works straightforwardly sometimes, otherwise not
- ▶ Example where it works nicely: asymptotic symmetry algebra
- ▶ Take linear combinations of Virasoro generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$

$$L_n = \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n = \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

- ▶ Make Inönü–Wigner contraction $\ell \rightarrow \infty$ on ASA

$$[L_n, L_m] = (n - m) L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[L_n, M_m] = (n - m) M_{n+m} + \frac{c_M}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[M_n, M_m] = 0$$

- ▶ This is nothing but the BMS_3 algebra (or $\text{GCA}_2, \text{URCA}_2, \text{CCA}_2$)!

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

- ▶ Works straightforwardly sometimes, otherwise not
- ▶ Example where it works nicely: asymptotic symmetry algebra
- ▶ Take linear combinations of Virasoro generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$

$$L_n = \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n = \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

- ▶ Make Inönü–Wigner contraction $\ell \rightarrow \infty$ on ASA

$$[L_n, L_m] = (n - m) L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[L_n, M_m] = (n - m) M_{n+m} + \frac{c_M}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[M_n, M_m] = 0$$

- ▶ This is nothing but the BMS_3 algebra (or $GCA_2, URCA_2, CCA_2$)!
- ▶ Example where it does not work easily: boundary conditions!

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

- ▶ Works straightforwardly sometimes, otherwise not
- ▶ Example where it works nicely: asymptotic symmetry algebra
- ▶ Take linear combinations of Virasoro generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$

$$L_n = \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n = \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

- ▶ Make Inönü–Wigner contraction $\ell \rightarrow \infty$ on ASA

$$[L_n, L_m] = (n - m) L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[L_n, M_m] = (n - m) M_{n+m} + \frac{c_M}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[M_n, M_m] = 0$$

- ▶ This is nothing but the BMS_3 algebra (or $\text{GCA}_2, \text{URCA}_2, \text{CCA}_2$)!
- ▶ Example where it does not work easily: boundary conditions!
- ▶ Example where it does not work at all: highest weight conditions!

Flat space holography (Barnich et al, Bagchi et al, Strominger et al, ...)

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

Not in general! Must (also) work intrinsically in flat space!

Interesting example:

- ▶ unitarity of flat space quantum gravity

Flat space holography (Barnich et al, Bagchi et al, Strominger et al, ...)

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

Not in general! Must (also) work intrinsically in flat space!

Interesting example:

- ▶ unitarity of flat space quantum gravity
- ▶ extrapolate from AdS: should be unitary (?)

Flat space holography (Barnich et al, Bagchi et al, Strominger et al, ...)

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

Not in general! Must (also) work intrinsically in flat space!

Interesting example:

- ▶ unitarity of flat space quantum gravity
- ▶ extrapolate from AdS: should be unitary (?)
- ▶ extrapolate from dS: should be non-unitary (?)

Flat space holography (Barnich et al, Bagchi et al, Strominger et al, ...)

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

Not in general! Must (also) work intrinsically in flat space!

Interesting example:

- ▶ unitarity of flat space quantum gravity
- ▶ extrapolate from AdS: should be unitary (?)
- ▶ extrapolate from dS: should be non-unitary (?)
- ▶ directly in flat space: both options realized, depending on details of model

Flat space holography (Barnich et al, Bagchi et al, Strominger et al, ...)

if holography is true \Rightarrow must work in flat space

Just take large AdS radius limit of 10^4 AdS/CFT papers?

Not in general! Must (also) work intrinsically in flat space!

Interesting example:

- ▶ unitarity of flat space quantum gravity
- ▶ extrapolate from AdS: should be unitary (?)
- ▶ extrapolate from dS: should be non-unitary (?)
- ▶ directly in flat space: both options realized, depending on details of model

Many open issues in flat space holography!

(Higher spin) gravity as Chern–Simons gauge theory...

...with weird boundary conditions (Achúcarro & Townsend '86; Witten '88; Bañados '96)

CS action (for AdS: $\mathfrak{sl}(2) \oplus \mathfrak{sl}(2)$):

$$S_{\text{CS}} = \frac{k}{4\pi} \int \text{CS}(A) - \frac{k}{4\pi} \int \text{CS}(\bar{A})$$

Variational principle:

$$\delta S_{\text{CS}}|_{\text{EOM}} = \frac{k}{4\pi} \int \text{Tr} (A \wedge \delta A - \bar{A} \wedge \delta \bar{A})$$

Well-defined for boundary conditions (similarly for \bar{A})

$$A_+ = 0 \quad \text{or} \quad A_- = 0 \quad \text{boundary coordinates } x^\pm$$

Example: asymptotically AdS_3 (Cartan-version of Brown–Henneaux)

(Higher spin) gravity as Chern–Simons gauge theory...

...with weird boundary conditions (Achúcarro & Townsend '86; Witten '88; Bañados '96)

CS action (for AdS: $\mathfrak{sl}(2) \oplus \mathfrak{sl}(2)$):

$$S_{\text{CS}} = \frac{k}{4\pi} \int \text{CS}(A) - \frac{k}{4\pi} \int \text{CS}(\bar{A})$$

Variational principle:

$$\delta S_{\text{CS}}|_{\text{EOM}} = \frac{k}{4\pi} \int \text{Tr} (A \wedge \delta A - \bar{A} \wedge \delta \bar{A})$$

Well-defined for boundary conditions (similarly for \bar{A})

$$A_+ = 0 \quad \text{or} \quad A_- = 0 \quad \text{boundary coordinates } x^\pm$$

Example: asymptotically AdS_3 (Cartan-version of Brown–Henneaux)

$$\begin{aligned} A_\rho &= L_\rho & \bar{A}_\rho &= -L_\rho \\ A_+ &= e^\rho L_1 + e^{-\rho} L(x^+) L_{-1} & \bar{A}_+ &= 0 \\ A_- &= 0 & \bar{A}_- &= -e^\rho L_{-1} - e^{-\rho} \bar{L}(x^-) L_1 \end{aligned}$$

Dreibein: $e/\ell \sim A - \bar{A}$, spin-connection: $\omega \sim A + \bar{A}$

İnönü–Wigner contraction of Virasoro (Barnich & Compère '06)

BMS₃ and GCA₂ (or rather, URCA₂)

- ▶ Take two copies of Virasoro, generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$, central charges c, \bar{c}

İnönü–Wigner contraction of Virasoro (Barnich & Compère '06)

BMS₃ and GCA₂ (or rather, URCA₂)

- ▶ Take two copies of Virasoro, generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$, central charges c, \bar{c}
- ▶ Define superrotations L_n and supertranslations M_n

$$L_n := \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n := \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

İnönü–Wigner contraction of Virasoro (Barnich & Compère '06)

BMS₃ and GCA₂ (or rather, URCA₂)

- ▶ Take two copies of Virasoro, generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$, central charges c, \bar{c}
- ▶ Define superrotations L_n and supertranslations M_n

$$L_n := \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n := \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

- ▶ Make ultrarelativistic boost, $\ell \rightarrow \infty$

$$[L_n, L_m] = (n - m) L_{n+m} + c_L \frac{1}{12} \delta_{n+m,0}$$

$$[L_n, M_m] = (n - m) M_{n+m} + c_M \frac{1}{12} \delta_{n+m,0}$$

$$[M_n, M_m] = 0$$

Global part: contracted to $\text{isl}(2)$ (generators: $L_{\pm 1}, L_0, M_{\pm 1}, M_0$)

İnönü–Wigner contraction of Virasoro (Barnich & Compère '06)

BMS₃ and GCA₂ (or rather, URCA₂)

- ▶ Take two copies of Virasoro, generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$, central charges c, \bar{c}
- ▶ Define superrotations L_n and supertranslations M_n

$$L_n := \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n := \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

- ▶ Make ultrarelativistic boost, $\ell \rightarrow \infty$

$$[L_n, L_m] = (n - m) L_{n+m} + c_L \frac{1}{12} \delta_{n+m,0}$$

$$[L_n, M_m] = (n - m) M_{n+m} + c_M \frac{1}{12} \delta_{n+m,0}$$

$$[M_n, M_m] = 0$$

- ▶ Is precisely the (centrally extended) BMS₃ algebra!
- ▶ Central charges:

$$c_L = c - \bar{c} \quad c_M = (c + \bar{c})/\ell$$

Inönü–Wigner contraction of Virasoro (Barnich & Compère '06)

BMS₃ and GCA₂ (or rather, URCA₂)

- ▶ Take two copies of Virasoro, generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$, central charges c, \bar{c}
- ▶ Define superrotations L_n and supertranslations M_n

$$L_n := \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n := \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

- ▶ Make ultrarelativistic boost, $\ell \rightarrow \infty$

$$[L_n, L_m] = (n - m) L_{n+m} + c_L \frac{1}{12} \delta_{n+m,0}$$

$$[L_n, M_m] = (n - m) M_{n+m} + c_M \frac{1}{12} \delta_{n+m,0}$$

$$[M_n, M_m] = 0$$

- ▶ Is precisely the (centrally extended) BMS₃ algebra!
- ▶ Central charges:

$$c_L = c - \bar{c} \quad c_M = (c + \bar{c})/\ell$$

Example TMG (with gravitational CS coupling μ and Newton constant G):

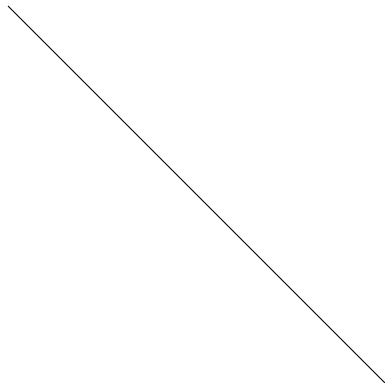
$$c_L = \frac{3}{\mu G} \quad c_M = \frac{3}{G}$$

Consequence of ultrarelativistic boost for AdS boundary

AdS-boundary:



Flat space boundary:



Limit $l \rightarrow \infty$

Null infinity holography!

Contraction on gravity side

AdS metric ($\varphi \sim \varphi + 2\pi$):

$$ds_{\text{AdS}}^2 = d(\ell\rho)^2 - \cosh^2\left(\frac{\ell\rho}{\ell}\right) dt^2 + \ell^2 \sinh^2\left(\frac{\ell\rho}{\ell}\right) d\varphi^2$$

Contraction on gravity side

AdS metric ($\varphi \sim \varphi + 2\pi$):

$$ds_{\text{AdS}}^2 = d(\ell\rho)^2 - \cosh^2\left(\frac{\ell\rho}{\ell}\right) dt^2 + \ell^2 \sinh^2\left(\frac{\ell\rho}{\ell}\right) d\varphi^2$$

Limit $\ell \rightarrow \infty$ ($r = \ell\rho$):

$$ds_{\text{Flat}}^2 = dr^2 - dt^2 + r^2 d\varphi^2 = -du^2 - 2du dr + r^2 d\varphi^2$$

Contraction on gravity side

AdS metric ($\varphi \sim \varphi + 2\pi$):

$$ds_{\text{AdS}}^2 = d(\ell\rho)^2 - \cosh^2\left(\frac{\ell\rho}{\ell}\right) dt^2 + \ell^2 \sinh^2\left(\frac{\ell\rho}{\ell}\right) d\varphi^2$$

Limit $\ell \rightarrow \infty$ ($r = \ell\rho$):

$$ds_{\text{Flat}}^2 = dr^2 - dt^2 + r^2 d\varphi^2 = -du^2 - 2du dr + r^2 d\varphi^2$$

BTZ metric:

$$ds_{\text{BTZ}}^2 = -\frac{\left(\frac{r^2}{\ell^2} - \frac{r_+^2}{\ell^2}\right)(r^2 - r_-^2)}{r^2} dt^2 + \frac{r^2 dr^2}{\left(\frac{r^2}{\ell^2} - \frac{r_+^2}{\ell^2}\right)(r^2 - r_-^2)} + r^2 \left(d\varphi - \frac{r_+}{r^2} r_- dt\right)^2$$

Contraction on gravity side

AdS metric ($\varphi \sim \varphi + 2\pi$):

$$ds_{\text{AdS}}^2 = d(\ell\rho)^2 - \cosh^2\left(\frac{\ell\rho}{\ell}\right) dt^2 + \ell^2 \sinh^2\left(\frac{\ell\rho}{\ell}\right) d\varphi^2$$

Limit $\ell \rightarrow \infty$ ($r = \ell\rho$):

$$ds_{\text{Flat}}^2 = dr^2 - dt^2 + r^2 d\varphi^2 = -du^2 - 2du dr + r^2 d\varphi^2$$

BTZ metric:

$$ds_{\text{BTZ}}^2 = -\frac{\left(\frac{r_+^2}{\ell^2} - \frac{r_-^2}{\ell^2}\right)(r^2 - r_-^2)}{r^2} dt^2 + \frac{r^2 dr^2}{\left(\frac{r_+^2}{\ell^2} - \frac{r_-^2}{\ell^2}\right)(r^2 - r_-^2)} + r^2 \left(d\varphi - \frac{r_+}{r^2} r_- dt\right)^2$$

Limit $\ell \rightarrow \infty$ ($\hat{r}_+ = \frac{r_+}{\ell} = \text{finite}$):

$$ds_{\text{FSC}}^2 = \hat{r}_+^2 \left(1 - \frac{r_-^2}{r^2}\right) dt^2 - \frac{1}{1 - \frac{r_-^2}{r^2}} \frac{dr^2}{\hat{r}_+^2} + r^2 \left(d\varphi - \frac{\hat{r}_+ r_-}{r^2} dt\right)^2$$

Contraction on gravity side

AdS metric ($\varphi \sim \varphi + 2\pi$):

$$ds_{\text{AdS}}^2 = d(\ell\rho)^2 - \cosh^2\left(\frac{\ell\rho}{\ell}\right) dt^2 + \ell^2 \sinh^2\left(\frac{\ell\rho}{\ell}\right) d\varphi^2$$

Limit $\ell \rightarrow \infty$ ($r = \ell\rho$):

$$ds_{\text{Flat}}^2 = dr^2 - dt^2 + r^2 d\varphi^2 = -du^2 - 2du dr + r^2 d\varphi^2$$

BTZ metric:

$$ds_{\text{BTZ}}^2 = -\frac{\left(\frac{r_+^2}{\ell^2} - \frac{r_-^2}{\ell^2}\right)(r^2 - r_-^2)}{r^2} dt^2 + \frac{r^2 dr^2}{\left(\frac{r_+^2}{\ell^2} - \frac{r_-^2}{\ell^2}\right)(r^2 - r_-^2)} + r^2 \left(d\varphi - \frac{r_+}{r^2} r_- dt\right)^2$$

Limit $\ell \rightarrow \infty$ ($\hat{r}_+ = \frac{r_+}{\ell} = \text{finite}$):

$$ds_{\text{FSC}}^2 = \hat{r}_+^2 \left(1 - \frac{r_-^2}{r^2}\right) dt^2 - \frac{1}{1 - \frac{r_-^2}{r^2}} \frac{dr^2}{\hat{r}_+^2} + r^2 \left(d\varphi - \frac{\hat{r}_+ r_-}{r^2} dt\right)^2$$

Shifted-boost orbifold studied by Cornalba & Costa more than decade ago

Describes expanding (contracting) Universe in flat space

Cosmological horizon at $r = r_-$, screening CTCs at $r < 0$

Outline

Motivations

Holography basics

Flat space gravity

Flat space higher spin gravity

Flat space higher spin gravity

Afshar, Bagchi, Fareghbal, DG, Rosseel '13, Gonzalez, Matulich, Pino, Troncoso '13

- ▶ AdS gravity in CS formulation: $\text{spin } 2 \rightarrow \text{spin } 3 \sim \text{sl}(2) \rightarrow \text{sl}(3)$

Flat space higher spin gravity

Afshar, Bagchi, Fareghbal, DG, Rosseel '13, Gonzalez, Matulich, Pino, Troncoso '13

- ▶ AdS gravity in CS formulation: spin 2 \rightarrow spin 3 \sim $\mathfrak{sl}(2) \rightarrow \mathfrak{sl}(3)$
- ▶ Flat space: similar!

$$S_{\text{CS}}^{\text{flat}} = \frac{k}{4\pi} \int \text{CS}(\mathcal{A})$$

with $\mathfrak{isl}(3)$ connection ($e^a =$ “zuvielbein”)

$$\mathcal{A} = e^a T_a + \omega^a J_a \quad T_a = (M_n, V_m) \quad J_a = (L_n, U_m)$$

$\mathfrak{isl}(3)$ algebra (spin 3 extension of global part of BMS/GCA algebra)

$$[L_n, L_m] = (n - m)L_{n+m}$$

$$[L_n, M_m] = (n - m)M_{n+m}$$

$$[L_n, U_m] = (2n - m)U_{n+m}$$

$$[M_n, U_m] = [L_n, V_m] = (2n - m)V_{n+m}$$

$$[U_n, U_m] = (n - m)(2n^2 + 2m^2 - nm - 8)L_{n+m}$$

$$[U_n, V_m] = (n - m)(2n^2 + 2m^2 - nm - 8)M_{n+m}$$

Flat space higher spin gravity

Afshar, Bagchi, Fareghbal, DG, Rosseel '13, Gonzalez, Matulich, Pino, Troncoso '13

- ▶ AdS gravity in CS formulation: $\text{spin } 2 \rightarrow \text{spin } 3 \sim \text{sl}(2) \rightarrow \text{sl}(3)$
- ▶ Flat space: similar!

$$S_{\text{CS}}^{\text{flat}} = \frac{k}{4\pi} \int \text{CS}(\mathcal{A})$$

with $\text{isl}(3)$ connection ($e^a =$ “zuvielbein”)

$$\mathcal{A} = e^a T_a + \omega^a J_a \quad T_a = (M_n, V_m) \quad J_a = (L_n, U_m)$$

- ▶ Same type of boundary conditions as for spin 2:

$$\mathcal{A}(r, t, \varphi) = b^{-1}(r) (d + a(t, \varphi) + o(1)) b(r)$$

Flat space higher spin gravity

Afshar, Bagchi, Fareghbal, DG, Rosseel '13, Gonzalez, Matulich, Pino, Troncoso '13

- ▶ AdS gravity in CS formulation: spin 2 \rightarrow spin 3 \sim $sl(2) \rightarrow sl(3)$
- ▶ Flat space: similar!

$$S_{\text{CS}}^{\text{flat}} = \frac{k}{4\pi} \int \text{CS}(\mathcal{A})$$

with $isl(3)$ connection ($e^a =$ “zuvielbein”)

$$\mathcal{A} = e^a T_a + \omega^a J_a \quad T_a = (M_n, V_m) \quad J_a = (L_n, U_m)$$

- ▶ Same type of boundary conditions as for spin 2:

$$\mathcal{A}(r, t, \varphi) = b^{-1}(r) (d + a(t, \varphi) + o(1)) b(r)$$

- ▶ Flat space boundary conditions: $b(r) = \exp(\frac{1}{2} r M_{-1})$ and

$$a(t, \varphi) = (M_1 - M(\varphi)M_{-1} - V(\varphi)V_{-2}) dt \\ + (L_1 - M(\varphi)L_{-1} - V(\varphi)U_{-2} - L(\varphi)M_{-1} - Z(\varphi)V_{-2}) d\varphi$$

Flat space higher spin gravity

Afshar, Bagchi, Fareghbal, DG, Rosseel '13, Gonzalez, Matulich, Pino, Troncoso '13

- ▶ AdS gravity in CS formulation: spin 2 \rightarrow spin 3 \sim $\mathfrak{sl}(2) \rightarrow \mathfrak{sl}(3)$
- ▶ Flat space: similar!

$$S_{\text{CS}}^{\text{flat}} = \frac{k}{4\pi} \int \text{CS}(\mathcal{A})$$

with $\mathfrak{isl}(3)$ connection ($e^a =$ “zuvielbein”)

$$\mathcal{A} = e^a T_a + \omega^a J_a \quad T_a = (M_n, V_m) \quad J_a = (L_n, U_m)$$

- ▶ Same type of boundary conditions as for spin 2:

$$\mathcal{A}(r, t, \varphi) = b^{-1}(r) (d + a(t, \varphi) + o(1)) b(r)$$

- ▶ Flat space boundary conditions: $b(r) = \exp(\frac{1}{2} r M_{-1})$ and

$$a(t, \varphi) = (M_1 - M(\varphi)M_{-1} - V(\varphi)V_{-2}) dt \\ + (L_1 - M(\varphi)L_{-1} - V(\varphi)U_{-2} - L(\varphi)M_{-1} - Z(\varphi)V_{-2}) d\varphi$$

- ▶ Spin 3 charges:

$$Q[\varepsilon_M, \varepsilon_L, \varepsilon_V, \varepsilon_U] \sim \oint (\varepsilon_M(\varphi)M(\varphi) + \varepsilon_L(\varphi)L(\varphi) + \varepsilon_V(\varphi)V(\varphi) + \varepsilon_U(\varphi)U(\varphi))$$

Flat space higher spin gravity

Asymptotic symmetry algebra at finite level k Afshar, Bagchi, Fareghbal, DG, Rosseel '13

- ▶ Do either Brown–Henneaux type of analysis or İnönü–Wigner contraction of two copies of quantum W_3 -algebra

- ▶ Do either Brown–Henneaux type of analysis or İnönü–Wigner contraction of two copies of quantum W_3 -algebra
- ▶ Obtain new type of W -algebra as extension of BMS (“BMW”)

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[L_n, M_m] = (n - m)M_{n+m} + \frac{c_M}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[U_n, U_m] = (n - m)(2n^2 + 2m^2 - nm - 8)L_{n+m} + \frac{192}{c_M} (n - m)\Lambda_{n+m} \\ - \frac{96(c_L + \frac{44}{5})}{c_M^2} (n - m)\Theta_{n+m} + \frac{c_L}{12} n(n^2 - 1)(n^2 - 4) \delta_{n+m, 0}$$

$$[U_n, V_m] = (n - m)(2n^2 + 2m^2 - nm - 8)M_{n+m} + \frac{96}{c_M} (n - m)\Theta_{n+m} \\ + \frac{c_M}{12} n(n^2 - 1)(n^2 - 4) \delta_{n+m, 0}$$

$$\Lambda_n = \sum_p : L_p M_{n-p} : - \frac{3}{10} (n + 2)(n + 3)M_n \quad \Theta_n = \sum_p M_p M_{n-p}$$

other commutators as in $\text{isl}(3)$ with $n \in \mathbb{Z}$

- ▶ Do either Brown–Henneaux type of analysis or İnönü–Wigner contraction of two copies of quantum W_3 -algebra
- ▶ Obtain new type of W -algebra as extension of BMS (“BMW”)

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[L_n, M_m] = (n - m)M_{n+m} + \frac{c_M}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[U_n, U_m] = (n - m)(2n^2 + 2m^2 - nm - 8)L_{n+m} + \frac{192}{c_M} (n - m)\Lambda_{n+m} \\ - \frac{96(c_L + \frac{44}{5})}{c_M^2} (n - m)\Theta_{n+m} + \frac{c_L}{12} n(n^2 - 1)(n^2 - 4) \delta_{n+m, 0}$$

$$[U_n, V_m] = (n - m)(2n^2 + 2m^2 - nm - 8)M_{n+m} + \frac{96}{c_M} (n - m)\Theta_{n+m} \\ + \frac{c_M}{12} n(n^2 - 1)(n^2 - 4) \delta_{n+m, 0}$$

- ▶ Note quantum shift and poles in central terms!

- ▶ Do either Brown–Henneaux type of analysis or İnönü–Wigner contraction of two copies of quantum W_3 -algebra
- ▶ Obtain new type of W -algebra as extension of BMS (“BMW”)

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[L_n, M_m] = (n - m)M_{n+m} + \frac{c_M}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[U_n, U_m] = (n - m)(2n^2 + 2m^2 - nm - 8)L_{n+m} + \frac{192}{c_M} (n - m)\Lambda_{n+m} \\ - \frac{96(c_L + \frac{44}{5})}{c_M^2} (n - m)\Theta_{n+m} + \frac{c_L}{12} n(n^2 - 1)(n^2 - 4) \delta_{n+m, 0}$$

$$[U_n, V_m] = (n - m)(2n^2 + 2m^2 - nm - 8)M_{n+m} + \frac{96}{c_M} (n - m)\Theta_{n+m} \\ + \frac{c_M}{12} n(n^2 - 1)(n^2 - 4) \delta_{n+m, 0}$$

- ▶ Note **quantum shift** and **poles** in central terms!
- ▶ Analysis generalizes to flat space contractions of other W -algebras

Unitarity in flat space

Unitarity leads to further contraction DG, Riegler, Rosseel '14

Facts:

- ▶ Unitarity in GCA requires $c_M = 0$ (see paper for caveats!)

Unitarity in flat space

Unitarity leads to further contraction DG, Riegler, Rosseel '14

Facts:

- ▶ Unitarity in GCA requires $c_M = 0$ (see paper for caveats!)
- ▶ Non-triviality requires then $c_L \neq 0$

Unitarity in flat space

Unitarity leads to further contraction DG, Riegler, Rosseel '14

Facts:

- ▶ Unitarity in GCA requires $c_M = 0$ (see paper for caveats!)
- ▶ Non-triviality requires then $c_L \neq 0$
- ▶ Generalization to contracted higher spin algebras straightforward

Unitarity in flat space

Unitarity leads to further contraction DG, Riegler, Rosseel '14

Facts:

- ▶ Unitarity in GCA requires $c_M = 0$ (see paper for caveats!)
- ▶ Non-triviality requires then $c_L \neq 0$
- ▶ Generalization to contracted higher spin algebras straightforward
- ▶ All of them contain GCA as subalgebra

Unitarity in flat space

Unitarity leads to further contraction DG, Riegler, Rosseel '14

Facts:

- ▶ Unitarity in GCA requires $c_M = 0$ (see paper for caveats!)
- ▶ Non-triviality requires then $c_L \neq 0$
- ▶ Generalization to contracted higher spin algebras straightforward
- ▶ All of them contain GCA as subalgebra
- ▶ $c_M = 0$ is necessary for unitarity

Unitarity in flat space

Unitarity leads to further contraction DG, Riegler, Rosseel '14

Facts:

- ▶ Unitarity in GCA requires $c_M = 0$ (see paper for caveats!)
- ▶ Non-triviality requires then $c_L \neq 0$
- ▶ Generalization to contracted higher spin algebras straightforward
- ▶ All of them contain GCA as subalgebra
- ▶ $c_M = 0$ is necessary for unitarity

Limit $c_M \rightarrow 0$ requires further contraction: $U_n \rightarrow c_M U_n$

Unitarity in flat space

Unitarity leads to further contraction DG, Riegler, Rosseel '14

Facts:

- ▶ Unitarity in GCA requires $c_M = 0$ (see paper for caveats!)
- ▶ Non-triviality requires then $c_L \neq 0$
- ▶ Generalization to contracted higher spin algebras straightforward
- ▶ All of them contain GCA as subalgebra
- ▶ $c_M = 0$ is necessary for unitarity

Limit $c_M \rightarrow 0$ requires further contraction: $U_n \rightarrow c_M U_n$

Doubly contracted algebra has unitary representations:

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[L_n, M_m] = (n - m)M_{n+m}$$

$$[L_n, U_m] = (2n - m)U_{n+m}$$

$$[M_n, U_m] = [L_n, V_m] = (2n - m)V_{n+m}$$

$$[U_n, U_m] \propto [U_n, V_m] = 96(n - m) \sum_p M_p M_{n-p}$$

Unitarity in flat space

Unitarity leads to further contraction DG, Riegler, Rosseel '14

Facts:

- ▶ Unitarity in GCA requires $c_M = 0$ (see paper for caveats!)
- ▶ Non-triviality requires then $c_L \neq 0$
- ▶ Generalization to contracted higher spin algebras straightforward
- ▶ All of them contain GCA as subalgebra
- ▶ $c_M = 0$ is necessary for unitarity

Limit $c_M \rightarrow 0$ requires further contraction: $U_n \rightarrow c_M U_n$

Doubly contracted algebra has unitary representations:

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

$$[L_n, M_m] = (n - m)M_{n+m}$$

$$[L_n, U_m] = (2n - m)U_{n+m}$$

$$[M_n, U_m] = [L_n, V_m] = (2n - m)V_{n+m}$$

$$[U_n, U_m] \propto [U_n, V_m] = 96(n - m) \sum_p M_p M_{n-p}$$

Higher spin states decouple and become null states!

Unitarity in flat space

Generic flat space W -algebras DG, Riegler, Rosseel '14

1. NO-GO:

Generically (see paper) you can have only two out of three:

- ▶ Unitarity
- ▶ Flat space
- ▶ Non-trivial higher spin states

Unitarity in flat space

Generic flat space W -algebras DG, Riegler, Rosseel '14

1. NO-GO:

Generically (see paper) you can have only two out of three:

- ▶ Unitarity
- ▶ Flat space
- ▶ Non-trivial higher spin states

Example:

Flat space chiral gravity

Bagchi, Detournay, DG, 1208.1658

Unitarity in flat space

Generic flat space W -algebras DG, Riegler, Rosseel '14

1. NO-GO:

Generically (see paper) you can have only two out of three:

- ▶ Unitarity
- ▶ Flat space
- ▶ Non-trivial higher spin states

Example:

Minimal model holography

Gaberdiel, Gopakumar, 1011.2986, 1207.6697

Unitarity in flat space

Generic flat space W -algebras DG, Riegler, Rosseel '14

1. NO-GO:

Generically (see paper) you can have only two out of three:

- ▶ Unitarity
- ▶ Flat space
- ▶ Non-trivial higher spin states

Example:

Flat space higher spin gravity (Galilean W_3 algebra)

Afshar, Bagchi, Fareghbal, DG and Rosseel, 1307.4768

Gonzalez, Matulich, Pino and Troncoso, 1307.5651

Unitarity in flat space

Generic flat space W -algebras DG, Riegler, Rosseel '14

1. NO-GO:

Generically (see paper) you can have only two out of three:

- ▶ Unitarity
- ▶ Flat space
- ▶ Non-trivial higher spin states

Compatible with “spirit” of various no-go results in higher dimensions!

Unitarity in flat space

Generic flat space W -algebras DG, Riegler, Rosseel '14

1. NO-GO:

Generically (see paper) you can have only two out of three:

- ▶ Unitarity
- ▶ Flat space
- ▶ Non-trivial higher spin states

Compatible with “spirit” of various no-go results in higher dimensions!

2. YES-GO:

There is (at least) one counter-example, namely a Vasiliev-type of theory, where you can have all three properties!

Unitary, non-trivial flat space higher spin algebra exists!
Vasiliev-type flat space chiral higher spin gravity?

Unitarity in flat space

Flat space W_∞ -algebra compatible with unitarity DG, Riegler, Rosseel '14

- ▶ We do not know if flat space **chiral** higher spin gravity exists...

Unitarity in flat space

Flat space W_∞ -algebra compatible with unitarity DG, Riegler, Rosseel '14

- ▶ We do not know if flat space **chiral** higher spin gravity exists...
- ▶ ...but its existence is at least not ruled out by the no-go result!

Unitarity in flat space

Flat space W_∞ -algebra compatible with unitarity DG, Riegler, Rosseel '14

- ▶ We do not know if flat space **chiral** higher spin gravity exists...
- ▶ ...but its existence is at least not ruled out by the no-go result!
- ▶ If it exists, this must be its asymptotic symmetry algebra:

$$[\mathcal{V}_m^i, \mathcal{V}_n^j] = \sum_{r=0}^{\lfloor \frac{i+j}{2} \rfloor} g_{2r}^{ij}(m, n) \mathcal{V}_{m+n}^{i+j-2r} + c_{\mathcal{V}}^i(m) \delta^{ij} \delta_{m+n,0}$$

$$[\mathcal{V}_m^i, \mathcal{W}_n^j] = \sum_{r=0}^{\lfloor \frac{i+j}{2} \rfloor} g_{2r}^{ij}(m, n) \mathcal{W}_{m+n}^{i+j-2r} \quad [\mathcal{W}_m^i, \mathcal{W}_n^j] = 0$$

where

$$c_{\mathcal{V}}^i(m) = \#(i, m) \times c \quad \text{and} \quad c = -\bar{c}$$

Unitarity in flat space

Flat space W_∞ -algebra compatible with unitarity DG, Riegler, Rosseel '14

- ▶ We do not know if flat space **chiral** higher spin gravity exists...
- ▶ ...but its existence is at least not ruled out by the no-go result!
- ▶ If it exists, this must be its asymptotic symmetry algebra:

$$[\mathcal{V}_m^i, \mathcal{V}_n^j] = \sum_{r=0}^{\lfloor \frac{i+j}{2} \rfloor} g_{2r}^{ij}(m, n) \mathcal{V}_{m+n}^{i+j-2r} + c_{\mathcal{V}}^i(m) \delta^{ij} \delta_{m+n,0}$$

$$[\mathcal{W}_m^i, \mathcal{W}_n^j] = \sum_{r=0}^{\lfloor \frac{i+j}{2} \rfloor} g_{2r}^{ij}(m, n) \mathcal{W}_{m+n}^{i+j-2r} \quad [\mathcal{W}_m^i, \mathcal{W}_n^j] = 0$$

where

$$c_{\mathcal{V}}^i(m) = \#(i, m) \times c \quad \text{and} \quad c = -\bar{c}$$

- ▶ Vacuum descendants $\mathcal{W}_m^i|0\rangle$ are null states for all i and m !

Unitarity in flat space

Flat space W_∞ -algebra compatible with unitarity DG, Riegler, Rosseel '14

- ▶ We do not know if flat space **chiral** higher spin gravity exists...
- ▶ ...but its existence is at least not ruled out by the no-go result!
- ▶ If it exists, this must be its asymptotic symmetry algebra:

$$[\mathcal{V}_m^i, \mathcal{V}_n^j] = \sum_{r=0}^{\lfloor \frac{i+j}{2} \rfloor} g_{2r}^{ij}(m, n) \mathcal{V}_{m+n}^{i+j-2r} + c_{\mathcal{V}}^i(m) \delta^{ij} \delta_{m+n,0}$$

$$[\mathcal{W}_m^i, \mathcal{W}_n^j] = \sum_{r=0}^{\lfloor \frac{i+j}{2} \rfloor} g_{2r}^{ij}(m, n) \mathcal{W}_{m+n}^{i+j-2r} \quad [\mathcal{W}_m^i, \mathcal{W}_n^j] = 0$$

where

$$c_{\mathcal{V}}^i(m) = \#(i, m) \times c \quad \text{and} \quad c = -\bar{c}$$

- ▶ Vacuum descendants $\mathcal{W}_m^i |0\rangle$ are null states for all i and m !
- ▶ AdS parent theory: no trace anomaly, but **gravitational anomaly** (Like in conformal Chern–Simons gravity \rightarrow Vasiliev type analogue?)

Long story short:

Long story short:

$$A_u \rightarrow A_u + \mu$$

Works nicely in Chern–Simons formulation!

Long story short:

$$A_u \rightarrow A_u + \mu$$

Works nicely in Chern–Simons formulation!

Line-element with spin-2 and spin-3 chemical potentials:

$$g_{\mu\nu} dx^\mu dx^\nu = \left(r^2 (\mu_L^2 - 4\mu_U'' \mu_U + 3\mu_U'^2 + 4\mathcal{M}\mu_U^2) + r g_{uu}^{(r)} + g_{uu}^{(0)} + g_{uu}^{(0')} \right) du^2 + \left(r^2 \mu_L - r\mu_M' + \mathcal{N}(1 + \mu_M) + 8\mathcal{Z}\mu_V \right) 2 du d\varphi - (1 + \mu_M) 2 dr du + r^2 d\varphi^2$$

$$g_{uu}^{(0)} = \mathcal{M}(1 + \mu_M)^2 + 2(1 + \mu_M)(\mathcal{N}\mu_L + 12\mathcal{V}\mu_V + 16\mathcal{Z}\mu_U) + 16\mathcal{Z}\mu_L\mu_V + \frac{4}{3}(\mathcal{M}^2\mu_V^2 + 4\mathcal{M}\mathcal{N}\mu_U\mu_V + \mathcal{N}^2\mu_U^2)$$

Spin-3 field with same chemical potentials:

$$\Phi_{\mu\nu\lambda} dx^\mu dx^\nu dx^\lambda = \Phi_{uuu} du^3 + \Phi_{ruu} dr du^2 + \Phi_{uu\varphi} du^2 d\varphi - (2\mu_U r^2 - r\mu_V' + 2\mathcal{N}\mu_V) dr du d\varphi + \mu_V dr^2 du - (\mu_U' r^3 - \frac{1}{3}r^2(\mu_V'' - \mathcal{M}\mu_V + 4\mathcal{N}\mu_U) + r\mathcal{N}\mu_V' - \mathcal{N}^2\mu_V) du d\varphi^2$$

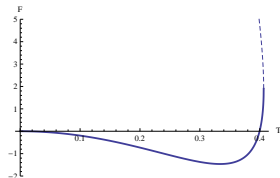
$$\begin{aligned} \Phi_{uuu} = & r^2 [2(1 + \mu_M)\mu_U(\mathcal{M}\mu_L - 4\mathcal{V}\mu_U) - \frac{1}{3}\mu_L^2(\mathcal{M}\mu_V - 4\mathcal{N}\mu_U) + 16\mu_L\mu_U(\mathcal{V}\mu_V + \mathcal{Z}\mu_U) - \frac{4}{3}\mathcal{M}\mu_U^2(\mathcal{M}\mu_V \\ & + 2\mathcal{N}\mu_U)] + 2\mathcal{V}(1 + \mu_M)^3 + \frac{2}{3}(1 + \mu_M)^2(6\mathcal{Z}\mu_L + \mathcal{M}^2\mu_V + 2\mathcal{M}\mathcal{N}\mu_U) + 16\mu_L\mu_V^2(\mathcal{N}\mathcal{V} - \frac{1}{3}\mathcal{M}\mathcal{Z}) \\ & + \frac{2}{3}(1 + \mu_M)((\mathcal{N}\mu_L + 16\mathcal{Z}\mu_U)(2\mathcal{M}\mu_V + \mathcal{N}\mu_U) + 12\mathcal{M}\mathcal{V}\mu_V^2) + \frac{64}{3}\mathcal{Z}\mu_U\mu_V(\mathcal{N}\mu_L + 12\mathcal{V}\mu_V + 12\mathcal{Z}\mu_U) \\ & + \mathcal{N}^2\mu_L^2\mu_V + 64\mathcal{V}^2\mu_V^3 - \frac{8}{27}(\mathcal{M}^3\mu_V^3 - \mathcal{N}^3\mu_U^3) - \frac{4}{9}\mathcal{M}\mathcal{N}\mu_U\mu_V(4\mathcal{M}\mu_V + 5\mathcal{N}\mu_U) + \sum_{n=0}^3 r^n \Phi_{uuu}^{(r^n)} \end{aligned}$$

Long story short:

$$A_u \rightarrow A_u + \mu$$

Works nicely in Chern–Simons formulation!

Interesting novel phase transitions of zeroth/first order:



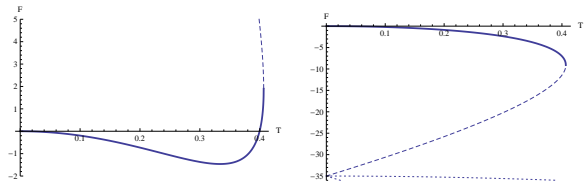
Free energy of four branches of regular solutions as function of temperature for different values of higher spin chemical potential ratio (in AdS: see [David, Ferlino, Kumar '12](#))

Long story short:

$$A_u \rightarrow A_u + \mu$$

Works nicely in Chern–Simons formulation!

Interesting novel phase transitions of zeroth/first order:



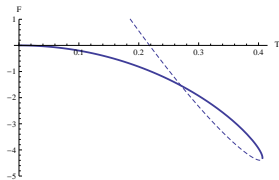
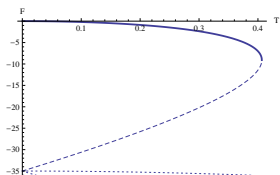
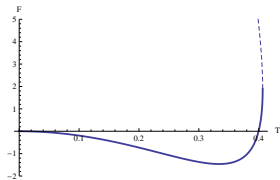
Free energy of four branches of regular solutions as function of temperature for different values of higher spin chemical potential ratio (in AdS: see [David, Ferlino, Kumar '12](#))

Long story short:

$$A_u \rightarrow A_u + \mu$$

Works nicely in Chern–Simons formulation!

Interesting novel phase transitions of zeroth/first order:



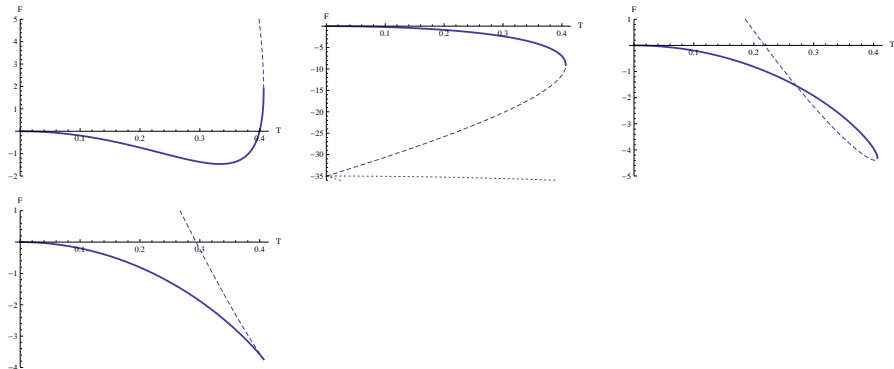
Free energy of four branches of regular solutions as function of temperature for different values of higher spin chemical potential ratio (in AdS: see David, Ferlino, Kumar '12)

Long story short:

$$A_u \rightarrow A_u + \mu$$

Works nicely in Chern–Simons formulation!

Interesting novel phase transitions of zeroth/first order:



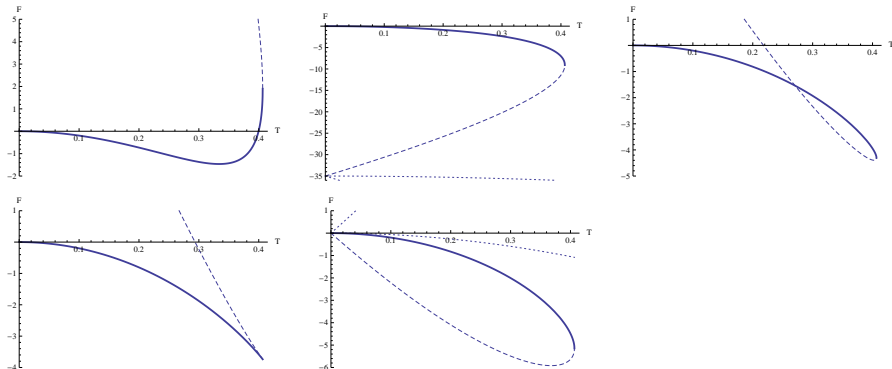
Free energy of four branches of regular solutions as function of temperature for different values of higher spin chemical potential ratio (in AdS: see [David, Ferlino, Kumar '12](#))

Long story short:

$$A_u \rightarrow A_u + \mu$$

Works nicely in Chern–Simons formulation!

Interesting novel phase transitions of zeroth/first order:



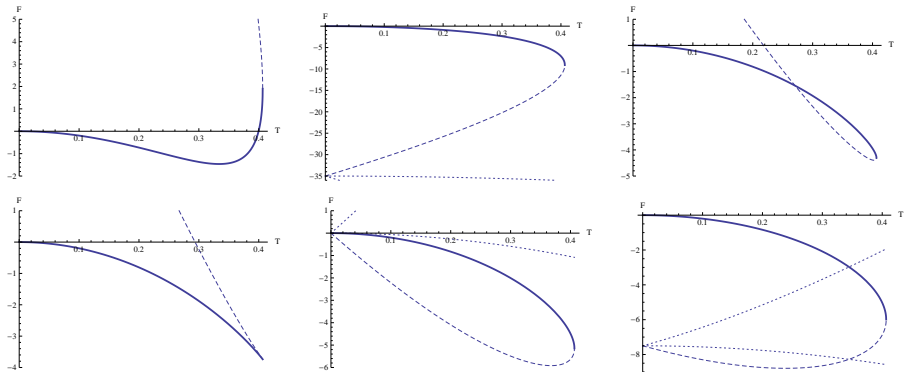
Free energy of four branches of regular solutions as function of temperature for different values of higher spin chemical potential ratio (in AdS: see David, Ferlino, Kumar '12)

Long story short:

$$A_u \rightarrow A_u + \mu$$

Works nicely in Chern–Simons formulation!

Interesting novel phase transitions of zeroth/first order:



Free energy of four branches of regular solutions as function of temperature for different values of higher spin chemical potential ratio (in AdS: see David, Ferlino, Kumar '12)

Flat space higher spin holography is a meaningful notion in 3D

Selected open issues

Flat space higher spin holography is a meaningful notion in 3D

Here are selected open issues:

- ▶ landscape of all possible phase transitions?

Flat space higher spin holography is a meaningful notion in 3D

Here are selected open issues:

- ▶ landscape of all possible phase transitions?
- ▶ existence of flat space chiral higher spin gravity?

Flat space higher spin holography is a meaningful notion in 3D

Here are selected open issues:

- ▶ landscape of all possible phase transitions?
- ▶ existence of flat space chiral higher spin gravity?
- ▶ other unitary examples?

Flat space higher spin holography is a meaningful notion in 3D

Here are selected open issues:

- ▶ landscape of all possible phase transitions?
- ▶ existence of flat space chiral higher spin gravity?
- ▶ other unitary examples?
- ▶ (holographic) entanglement entropy? (Bagchi, Basu, DG, Riegler '14)

Flat space higher spin holography is a meaningful notion in 3D

Here are selected open issues:

- ▶ landscape of all possible phase transitions?
- ▶ existence of flat space chiral higher spin gravity?
- ▶ other unitary examples?
- ▶ (holographic) entanglement entropy? ([Bagchi, Basu, DG, Riegler '14](#))
- ▶ flat space local quantum quench?

Flat space higher spin holography is a meaningful notion in 3D

Here are selected open issues:

- ▶ landscape of all possible phase transitions?
- ▶ existence of flat space chiral higher spin gravity?
- ▶ other unitary examples?
- ▶ (holographic) entanglement entropy? ([Bagchi, Basu, DG, Riegler '14](#))
- ▶ flat space local quantum quench?

Flat space higher spin holography provides a new playground
Contributes to long-term goal: find how general is holography

Thanks for your attention!

