Black Holes I (136.028)

Daniel Grumiller

Institute for Theoretical Physics TU Wien

Outlook on black hole research, January 2020



grumil@hep.itp.tuwien.ac.at

Simple idea:

Harmonic oscillator: take a physical system and shake it

Amazingly successful:

Simple idea:

Harmonic oscillator: take a physical system and shake it

Amazingly successful:

QFT corrections to Hydrogen atom

Feynman diagrams contributing to Lamb shift

Simple idea:

Harmonic oscillator: take a physical system and shake it

Amazingly successful:

- QFT corrections to Hydrogen atom
- weakly coupled phonons and electrons in cond-mat





Simple idea:

Harmonic oscillator: take a physical system and shake it

Amazingly successful:

- QFT corrections to Hydrogen atom
- weakly coupled phonons and electrons
- Standard Model of particle physics





Simple idea:

Harmonic oscillator: take a physical system and shake it

Amazingly successful:

- QFT corrections to Hydrogen atom
- weakly coupled phonons and electrons in cond-mat
- Standard Model of particle physics
- see also the TU Wien curriculum

Lectures at TU Wien containing harmonic oscillator (PP = particle physics)

- Intro to PP
- Atomic, nuclear & PP
- Intro to QFT
- Theor. methods in PP

- Astro-PP
- Black holes
- Intro to QED
- Field theory and pheno

- Path integrals
- Thermal field theory
- Cosmology
- various basic lectures

Appetizer, Part II Physics of the $21^{\rm st}$ century: black holes?

Application of harmonic oscillator limited to perturbative phenomena

Appetizer, Part II Physics of the 21^{st} century: black holes?

Application of harmonic oscillator limited to perturbative phenomena

Many physical systems require non-perturbative physics:

- QCD at low energies
- High T_c superconductors
- Graphene
- Cold atoms
- Gravity at high curvature

Generally speaking:

Strongly coupled systems require new techniques

Application of harmonic oscillator limited to perturbative phenomena

Many physical systems require non-perturbative physics:

- QCD at low energies
- High T_c superconductors
- Graphene
- Cold atoms
- Gravity at high curvature

Generally speaking:

Strongly coupled systems require new techniques

Punch-line of this outlook:

Black hole holography can provide such a technique

Appetizer, Part III

Black holes have apparently paradoxical properties

Black holes: The simplest macroscopic objects in the Universe



Properties determined by:

- ► Mass M
- Angular momentum J
- Charge(s) Q

Black hole \sim elementary particle!

Appetizer, Part III

Black holes have apparently paradoxical properties

Black holes: The simplest macroscopic objects in the Universe



Properties determined by:

- ► Mass M
- Angular momentum J
- Charge(s) Q

Black hole \sim elementary particle!

Black holes: The most complicated objects conceivable



Quantum mechanics:

- Black holes radiate
- Black holes have entropy
- Black holes are holographic

Bekenstein–Hawking entropy: $S_{\rm BH} \sim A/4$

Outline

Black hole experiments

Black hole theory

Black hole holography

Outline

Black hole experiments

Black hole theory

Black hole holography

Do black holes exist in our Universe? (image source: random webpage) Ham's scho ans g'sehn? (Ernst Mach to Ludwig Boltzmann conerning the existence of atoms)

Artistic black hole binary impression



Do black holes exist in our Universe? (image source: 1502.03808) Ham's scho ans g'sehn? (Ernst Mach to Ludwig Boltzmann conerning the existence of atoms)

Numerical black hole simulation (interstellar)



Do black holes exist in our Universe? (image source: 1906.11238) Ham's scho ans g'sehn? (Ernst Mach to Ludwig Boltzmann conerning the existence of atoms)

Photo of black hole M87 (EHT collaboration)



"Seeing is believing" is infantile; why did people believe in black holes before the photo?

"Seeing is believing" is infantile; why did people believe in black holes before the photo?

Experimental evidence for/against various black hole candidates:

stellar black holes: gravitational collapse Chandrasekhar 1930

after fusion in star stops:

Fermi pressure of electrons/neutrons prevents gravitational collapse critical mass from Fermi pressure = gravitational pressure rough estimate yields

$$M_{\rm critical} \sim \frac{1}{m_N^2} \sim 10^{38} \sim 10^{30} kg \sim M_{\odot}$$

more refined calculation yields

$$M_{\rm critical} \approx 3 M_{\odot}$$

stellar objects with mass $> 3 M_{\odot}$ thus are black holes

"Seeing is believing" is infantile; why did people believe in black holes before the photo?

- stellar black holes: gravitational collapse Chandrasekhar 1930
- ▶ stellar black holes: accretion disk physics Bolton, Webster, Murdin 1972 Objects whose mass is clearly beyond critical $M > 3M_{\odot}$ (from '06):

System	$P_{\rm orb}$	f(M)	Donor	Classification	M_x^{\dagger}
	[days]	$[M_{\odot}]$	Spect. Type		$[M_{\odot}]$
GRS 1915+105 ^a	33.5	9.5 ± 3.0	K/M III	LMXB/Transient	14 ± 4
V404 Cyg	6.471	6.09 ± 0.04	K0 IV	,,	12 ± 2
Cyg X-1	5.600	0.244 ± 0.005	09.7 lab	HMXB/Persistent	10 ± 3
LMC X-1	4.229	0.14 ± 0.05	07 III	,,	> 4
XTE J1819-254	2.816	3.13 ± 0.13	B9 III	IMXB/Transient	7.1 ± 0.3
GRO J1655-40	2.620	2.73 ± 0.09	F3/5 IV	,,	6.3 ± 0.3
BW Cir ^b	2.545	5.74 ± 0.29	G5 IV	LMXB/Transient	> 7.8
GX 339-4	1.754	5.8 ± 0.5	-	,,	
LMC X-3	1.704	2.3 ± 0.3	B3 V	HMXB/Persistent	7.6 ± 1.3
XTE J1550-564	1.542	6.86 ± 0.71	G8/K8 IV	LMXB/Transient	9.6 ± 1.2
4U 1543-475	1.125	0.25 ± 0.01	Á2 V	IMXB/Transient	9.4 ± 1.0
H1705-250	0.520	4.86 ± 0.13	K3/7 V	LMXB/Transient	6 ± 2
GS 1124-684	0.433	3.01 ± 0.15	K3/5 V	.,,	7.0 ± 0.6
XTE J1859+226 ^c	0.382	7.4 ± 1.1	_	,,	
GS2000+250	0.345	5.01 ± 0.12	K3/7 V	,,	7.5 ± 0.3
A0620-003	0.325	2.72 ± 0.06	K4 V	,,	11 ± 2
XTE J1650-500	0.321	2.73 ± 0.56	K4 V	,,	
GRS 1009-45	0.283	3.17 ± 0.12	K7/M0 V	,,	5.2 ± 0.6
GRO J0422+32	0.212	1.19 ± 0.02	M2 V	,,	4 ± 1
XTE J1118+480	0.171	6.3 ± 0.2	K5/M0 V	,,	6.8 ± 0.4

"Seeing is believing" is infantile; why did people believe in black holes before the photo?

- stellar black holes: gravitational collapse Chandrasekhar 1930
- stellar black holes: accretion disk physics Bolton, Webster, Murdin 1972
- stellar black hole mergers: gravitational wave production LIGO '16 Hanford, Washington (H1) Livingston, Louisiana (L1)



Pre-photographic evidence for black holes "Seeing is believing" is infantile; why did people believe in black holes before the photo?

- stellar black holes: gravitational collapse Chandrasekhar 1930
- stellar black holes: accretion disk physics Bolton, Webster, Murdin 1972
- stellar black hole mergers: gravitational wave production LIGO '16
- supermassive black holes: Kepler orbits Gillessen, Eisenhauer, Trippe et al '09



Evidence for black holes "Seeing is believing"

Experimental evidence for/against various black hole candidates:

- stellar black holes: gravitational collapse Chandrasekhar 1930
- stellar black holes: accretion disk physics Bolton, Webster, Murdin 1972
- stellar black hole mergers: gravitational wave production LIGO '16
- supermassive black holes: Kepler orbits Gillessen, Eisenhauer, Trippe et al '09

supermassive black holes: shadow EHT '19



Evidence for black holes "Seeing is believing"

- stellar black holes: gravitational collapse Chandrasekhar 1930
- ► stellar black holes: accretion disk physics Bolton, Webster, Murdin 1972
- stellar black hole mergers: gravitational wave production LIGO '16
- supermassive black holes: Kepler orbits Gillessen, Eisenhauer, Trippe et al '09
- supermassive black holes: shadow EHT '19
- ▶ intermediate black holes: some evidence $(100 10^6 M_{\odot})$

Evidence for black holes "Seeing is believing"

- stellar black holes: gravitational collapse Chandrasekhar 1930
- ► stellar black holes: accretion disk physics Bolton, Webster, Murdin 1972
- stellar black hole mergers: gravitational wave production LIGO '16
- supermassive black holes: Kepler orbits Gillessen, Eisenhauer, Trippe et al '09
- supermassive black holes: shadow EHT '19
- ▶ intermediate black holes: some evidence $(100 10^6 M_{\odot})$
- primordial black holes: no signatures from cosmology

Evidence for black holes "Seeing is believing"

- stellar black holes: gravitational collapse Chandrasekhar 1930
- ► stellar black holes: accretion disk physics Bolton, Webster, Murdin 1972
- stellar black hole mergers: gravitational wave production LIGO '16
- supermassive black holes: Kepler orbits Gillessen, Eisenhauer, Trippe et al '09
- supermassive black holes: shadow EHT '19
- ▶ intermediate black holes: some evidence $(100 10^6 M_{\odot})$
- primordial black holes: no signatures from cosmology
- particle-produced black holes: no signatures from LHC

Evidence for black holes "Seeing is believing"

- stellar black holes: gravitational collapse Chandrasekhar 1930
- ► stellar black holes: accretion disk physics Bolton, Webster, Murdin 1972
- stellar black hole mergers: gravitational wave production LIGO '16
- supermassive black holes: Kepler orbits Gillessen, Eisenhauer, Trippe et al '09
- supermassive black holes: shadow EHT '19
- ▶ intermediate black holes: some evidence $(100 10^6 M_{\odot})$
- primordial black holes: no signatures from cosmology
- particle-produced black holes: no signatures from LHC
 - overwhelming evidence for stellar and supermassive black holes
 - \blacktriangleright confirmed mass ranges so far: $3-100 M_{\odot}$ and $10^6-10^{10} M_{\odot}$
 - black holes could in principle exist for any mass > $M_{\rm Planck}$

Outline

Black hole experiments

Black hole theory

Black hole holography

defining property: black hole horizon

- defining property: black hole horizon
- mathematical definition: horizon = boundary of past of null infinity



- defining property: black hole horizon
- mathematical definition: horizon = boundary of past of null infinity
- physical definition: horizon = point of no return



- defining property: black hole horizon
- mathematical definition: horizon = boundary of past of null infinity
- physical definition: horizon = point of no return
- black hole formation:

gravitational quadrupoles and higher radiated away electromagnetic dipoles and higher radiated away

- defining property: black hole horizon
- mathematical definition: horizon = boundary of past of null infinity
- physical definition: horizon = point of no return
- black hole formation:

gravitational quadrupoles and higher radiated away electromagnetic dipoles and higher radiated away

- classical black hole observables:
 - electric monopole charge = charge Q (irrelevant so far in observations)
 - magnetic monopole charge = probably irrelevant
 - electric monopole mass = mass M
 - magnetic monopole mass = probably irrelevant
 - electric angular momentum = angular momentum J
 - magnetic angular momentum = probably irrelevant

- defining property: black hole horizon
- mathematical definition: horizon = boundary of past of null infinity
- physical definition: horizon = point of no return
- black hole formation:

gravitational quadrupoles and higher radiated away electromagnetic dipoles and higher radiated away

- classical black hole observables:
 - electric monopole charge = charge Q (irrelevant so far in observations)
 - magnetic monopole charge = probably irrelevant
 - electric monopole mass = mass M
 - magnetic monopole mass = probably irrelevant
 - electric angular momentum = angular momentum J
 - magnetic angular momentum = probably irrelevant

Essentially same observables as for elementary particles: charge Q, mass M, spin $J \Rightarrow$ amazingly simple!

Thermodynamics and black holes — black hole thermodynamics?

Thermodynamics

Zeroth law:

T = const. in equilibrium

T: temperature

Black hole mechanics Zeroth law:

 $\kappa = {\rm const.}$ f. stationary black holes

 κ : surface gravity

Thermodynamics and black holes — black hole thermodynamics?

Thermodynamics

Zeroth law:

T = const. in equilibrium

First law: $dE \sim TdS +$ work terms

T: temperature

E: energy

S: entropy

Black hole mechanics Zeroth law: $\kappa = \text{const. f. stationary black holes}$ First law:

 $dM \sim \kappa dA +$ work terms

 κ : surface gravity M: mass A: area (of event horizon)

Thermodynamics and black holes — black hole thermodynamics?

Thermodynamics

Zeroth law:

T = const. in equilibrium

First law: $dE \sim TdS +$ work terms

Second law: dS > 0

- T: temperature
- E: energy
- S: entropy

Black hole mechanicsZeroth law:
 $\kappa = \text{const. f. stationary black holes}$ First law:
 $dM \sim \kappa dA + \text{ work terms}$ Second law:
 $dA \ge 0$

 κ : surface gravity M: mass A: area (of event horizon)
Thermodynamics and black holes — black hole thermodynamics?

Thermodynamics

Zeroth law:

T = const. in equilibrium

First law: $dE \sim TdS +$ work terms

Second law: $dS \ge 0$

Third law: $T \rightarrow 0$ impossible

T: temperature

- E: energy
- $S: \ {\rm entropy}$

Black hole mechanics Zeroth law: $\kappa = \text{const. f. stationary black holes}$ First law: $dM \sim \kappa dA +$ work terms Second law: dA > 0Third law: $\kappa \to 0$ impossible κ : surface gravity

M: mass

A: area (of event horizon)

Formal analogy or actual physics?

Hawking effect (QFT under external conditions)

Black holes evaporate due to semi-classical effects!



Outline

Black hole experiments

Black hole theory

Black hole holography

Currently template for experimental data in quantum gravity

$$S_{\rm BH} = \frac{A}{4}$$

Currently template for experimental data in quantum gravity

$$S_{\rm BH} = rac{A}{4} \sim {\rm length}^{d-1}$$

entropy not extensive (d: number of spatial dimensions)

Currently template for experimental data in quantum gravity

$$S_{\rm BH} = \frac{A}{4} \sim {\rm length}^{d-1} \sim {\rm volume} \big|_{d-1}$$

- entropy not extensive (d: number of spatial dimensions)
- entropy would be extensive if theory was in 1 lower dimension

Currently template for experimental data in quantum gravity

$$S_{\rm BH} = \frac{A}{4} \sim {\rm length}^{d-1} \sim {\rm volume} \big|_{d-1}$$

- entropy not extensive (d: number of spatial dimensions)
- entropy would be extensive if theory was in 1 lower dimension

- Holographic principle 't Hooft '93; Susskind '95 -

Gravity in d spatial dimensions equivalent to QFT in d-1 spatial dimensions

Currently template for experimental data in quantum gravity

$$S_{\rm BH} = \frac{A}{4} \sim {\rm length}^{d-1} \sim {\rm volume} \big|_{d-1}$$

- entropy not extensive (d: number of spatial dimensions)
- entropy would be extensive if theory was in 1 lower dimension

- Holographic principle 't Hooft '93; Susskind '95 -

Gravity in d spatial dimensions equivalent to QFT in d-1 spatial dimensions

WTF? How can this be true?

Best studied realization of holography is AdS/CFT correspondence:

AdS is a negatively curved spacetime (maximally symmetric)



Open Universe Looking from inside, boundary at infinity Limit Circle IV, by M. C. Escher

Best studied realization of holography is AdS/CFT correspondence:

- AdS is a negatively curved spacetime (maximally symmetric)
- CFT is a field theory with conformal symmetry

Conformal symmetry includes scaling symmetry

coordinates: $x^{\mu} \to \lambda x^{\mu}$ energy: $E \to E/\lambda$

Best studied realization of holography is AdS/CFT correspondence:

- AdS is a negatively curved spacetime (maximally symmetric)
- CFT is a field theory with conformal symmetry

Conformal symmetry includes scaling symmetry

coordinates: $x^{\mu} \to \lambda x^{\mu}$ energy: $E \to E/\lambda$

Idea: treat energy as the fifth coordinate (RG-flow)

Best studied realization of holography is AdS/CFT correspondence:

- AdS is a negatively curved spacetime (maximally symmetric)
- CFT is a field theory with conformal symmetry

Conformal symmetry includes scaling symmetry

coordinates: $x^{\mu} \to \lambda x^{\mu}$ energy: $E \to E/\lambda$

Idea: treat energy as the fifth coordinate (RG-flow) Most general line-element compatible with symmetries:

$$ds^{2} = (E/L)^{2} \eta_{\mu\nu} dx^{\mu} dx^{\nu} + (L/E)^{2} dE^{2}$$

 \boldsymbol{L} sets physical scales

Best studied realization of holography is AdS/CFT correspondence:

- AdS is a negatively curved spacetime (maximally symmetric)
- CFT is a field theory with conformal symmetry

Conformal symmetry includes scaling symmetry

coordinates: $x^{\mu} \to \lambda x^{\mu}$ energy: $E \to E/\lambda$

Idea: treat energy as the fifth coordinate (RG-flow) Most general line-element compatible with symmetries:

$$ds^{2} = (E/L)^{2} \eta_{\mu\nu} dx^{\mu} dx^{\nu} + (L/E)^{2} dE^{2}$$

L sets physical scales and is called "AdS-radius"

This is precisely the line element of AdS in 1 dimension higher!

$\begin{array}{l} \text{AdS}_3/\text{CFT}_2 \\ \text{Focus on symmetries} \end{array}$

2d CFTs allow to apply powerful methods; let us focus on CFT_2

- Symmetries of CFT_d : SO(d, 2) if d > 2
- Symmetries of AdS_{d+1} : SO(d, 2) if d > 1

$\begin{array}{l} \mathsf{AdS}_3/\mathsf{CFT}_2 \\ \mathsf{Focus \ on \ symmetries} \end{array}$

2d CFTs allow to apply powerful methods; let us focus on CFT_2

- Symmetries of CFT_d : SO(d, 2) if d > 2
- Symmetries of AdS_{d+1} : SO(d, 2) if d > 1
- matches for d > 2, but what about d = 2?

2d CFTs allow to apply powerful methods; let us focus on CFT_2

- Symmetries of CFT_d : SO(d, 2) if d > 2
- Symmetries of AdS_{d+1} : SO(d,2) if d > 1
- matches for d > 2, but what about d = 2?
- CFT₂ has infinitely many symmetries!

$$[L_n, L_m] = (n-m) L_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0} \qquad n, m \in \mathbb{Z}$$

symmetry algebra is called 'Virasoro algebra with central charge c'

fineprint: there are actually two copies of this algebra; to reduce clutter we focus on one of them

2d CFTs allow to apply powerful methods; let us focus on CFT_2

- Symmetries of CFT_d : SO(d, 2) if d > 2
- Symmetries of AdS_{d+1} : SO(d,2) if d > 1
- matches for d > 2, but what about d = 2?
- CFT₂ has infinitely many symmetries!

$$[L_n, L_m] = (n-m) L_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0} \qquad n, m \in \mathbb{Z}$$

symmetry algebra is called 'Virasoro algebra with central charge c^\prime

fineprint: there are actually two copies of this algebra; to reduce clutter we focus on one of them

how can this possibly be true in AdS₃?

2d CFTs allow to apply powerful methods; let us focus on CFT_2

- Symmetries of CFT_d : SO(d, 2) if d > 2
- Symmetries of AdS_{d+1} : SO(d,2) if d > 1
- matches for d > 2, but what about d = 2?
- CFT₂ has infinitely many symmetries!

$$[L_n, L_m] = (n-m) L_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0} \qquad n, m \in \mathbb{Z}$$

symmetry algebra is called 'Virasoro algebra with central charge c'

fineprint: there are actually two copies of this algebra; to reduce clutter we focus on one of them

- how can this possibly be true in AdS₃?
- how can any gravity theory have an infinite amount of symmetries?!

at most D(D+1)/2 Killing vectors ξ in D spacetime dimensions

$$\mathcal{L}_{\xi}g_{\mu\nu} = \xi^{\alpha}\partial_{\alpha}g_{\mu\nu} + g_{\mu\alpha}\partial_{\nu}\xi^{\alpha} + g_{\nu\alpha}\partial_{\mu}\xi^{\alpha} = 0$$

e.g. in D = 4 for $g_{\mu\nu} = \eta_{\mu\nu}$ get $\xi^{\alpha} = \xi^{\alpha}_{(0)} + \Lambda^{\alpha}{}_{\beta}x^{\beta}$ = Poincaré trafos

2d CFTs allow to apply powerful methods; let us focus on CFT_2

- Symmetries of CFT_d : SO(d, 2) if d > 2
- Symmetries of AdS_{d+1} : SO(d,2) if d > 1
- matches for d > 2, but what about d = 2?
- CFT₂ has infinitely many symmetries!

$$[L_n, L_m] = (n-m) L_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0} \qquad n, m \in \mathbb{Z}$$

symmetry algebra is called 'Virasoro algebra with central charge c'

fineprint: there are actually two copies of this algebra; to reduce clutter we focus on one of them

how can this possibly be true in AdS₃?

how can any gravity theory have an infinite amount of symmetries?!

Last question was resolved already in 1960ies by Bondi, van der Burgh, Metzner and Sachs (BMS)

The limit of general relativity at small curvature is special relativity. Right?

• consider metrics $g_{\mu\nu}$ that asymptote to Minkowski metric $\eta_{\mu\nu}$

$$g_{\mu\nu} = \eta_{\mu\nu} + \delta g_{\mu\nu}$$

where $\delta g_{\mu\nu}$ is chosen suitable at infinity

The limit of general relativity at small curvature is special relativity. Right? Right?!

• consider metrics $g_{\mu\nu}$ that asymptote to Minkowski metric $\eta_{\mu\nu}$

$$g_{\mu\nu} = \eta_{\mu\nu} + \delta g_{\mu\nu}$$

where $\delta g_{\mu\nu}$ is chosen suitable at infinity

determine all asymptotic Killing vectors

$$\mathcal{L}_{\xi}g_{\mu\nu} = \mathcal{O}(\delta g_{\mu\nu})$$

The limit of general relativity at small curvature is special relativity. Right? Wrong!

 \blacktriangleright consider metrics $g_{\mu
u}$ that asymptote to Minkowski metric $\eta_{\mu
u}$

$$g_{\mu\nu} = \eta_{\mu\nu} + \delta g_{\mu\nu}$$

where $\delta g_{\mu\nu}$ is chosen suitable at infinity

determine all asymptotic Killing vectors

$$\mathcal{L}_{\xi}g_{\mu\nu} = \mathcal{O}(\delta g_{\mu\nu})$$

- shocking discovery by BMS:
 - asymptotic symmetry algebra much larger than Poincaré
 - infinitely many symmetries ('supertranslations')
 - do not just get special relativity

The limit of general relativity at small curvature is special relativity. Right? Wrong!

• consider metrics $g_{\mu\nu}$ that asymptote to Minkowski metric $\eta_{\mu\nu}$

$$g_{\mu\nu} = \eta_{\mu\nu} + \delta g_{\mu\nu}$$

where $\delta g_{\mu\nu}$ is chosen suitable at infinity

determine all asymptotic Killing vectors

$$\mathcal{L}_{\xi}g_{\mu\nu} = \mathcal{O}(\delta g_{\mu\nu})$$

- shocking discovery by BMS:
 - asymptotic symmetry algebra much larger than Poincaré
 - infinitely many symmetries ('supertranslations')
 - do not just get special relativity

Two important lessons

- 1. BMS symmetries of relevance for QFTs (soft theorems)
- 2. more generally, gravity theories can have infinitely many symmetries

consider asymptotically AdS₃ metrics

$$g_{\mu\nu} = \bar{g}_{\mu\nu}^{\rm AdS} + \delta g_{\mu\nu}$$

with suitable* choice of $\delta g_{\mu\nu}$

* if you care about details read Max Riegler's DKPI PhD thesis 1609.02733

consider asymptotically AdS₃ metrics

$$g_{\mu\nu} = \bar{g}_{\mu\nu}^{\rm AdS} + \delta g_{\mu\nu}$$

with suitable choice of $\delta g_{\mu\nu}$

determine all asymptotic Killing vectors

$$\mathcal{L}_{\xi}g_{\mu\nu} = \mathcal{O}(\delta g_{\mu\nu})$$

consider asymptotically AdS₃ metrics

$$g_{\mu\nu} = \bar{g}_{\mu\nu}^{\rm AdS} + \delta g_{\mu\nu}$$

with suitable choice of $\delta g_{\mu\nu}$

determine all asymptotic Killing vectors

$$\mathcal{L}_{\xi}g_{\mu\nu} = \mathcal{O}(\delta g_{\mu\nu})$$

▶ for choice of $\delta g_{\mu\nu}$ proposed by Brown, Henneaux '86

 $\xi = \epsilon(z)\partial_z + \text{subleading}$

consider asymptotically AdS₃ metrics

$$g_{\mu\nu} = \bar{g}_{\mu\nu}^{\rm AdS} + \delta g_{\mu\nu}$$

determine all asymptotic Killing vectors

$$\mathcal{L}_{\xi}g_{\mu\nu} = \mathcal{O}(\delta g_{\mu\nu})$$

▶ for choice of $\delta g_{\mu\nu}$ proposed by Brown, Henneaux '86

$$\xi = \epsilon(z)\partial_z + \text{subleading}$$

▶ Fourier modes $l_n \sim \oint e^{inz} \epsilon(z) \, dz$ yield asymptotic symmetry algebra

$$[l_n, l_m]_{\text{Lie}} = (n-m) \, l_{n+m}$$

consider asymptotically AdS₃ metrics

$$g_{\mu\nu} = \bar{g}_{\mu\nu}^{\rm AdS} + \delta g_{\mu\nu}$$

determine all asymptotic Killing vectors

$$\mathcal{L}_{\xi}g_{\mu\nu} = \mathcal{O}(\delta g_{\mu\nu})$$

▶ for choice of $\delta g_{\mu\nu}$ proposed by Brown, Henneaux '86

$$\xi = \epsilon(z)\partial_z + \text{subleading}$$

▶ Fourier modes $l_n \sim \oint e^{inz} \epsilon(z) \, dz$ yield asymptotic symmetry algebra

$$[l_n, l_m]_{\text{Lie}} = (n-m) \, l_{n+m}$$

recall CFT_2 symmetries

$$[L_n, L_m] = (n-m) L_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0}$$

Almost Virasoro, so almost there — but central charge c missing

matching of symmetries so far works only up to central term

Physical phase space of AdS3 Einstein gravity

- matching of symmetries so far works only up to central term
- reason: did not yet consider specific gravity theory

- matching of symmetries so far works only up to central term
- reason: did not yet consider specific gravity theory
- focus on Einstein gravity with negative cosmological constant

$$I_{\rm EH} = -\frac{1}{16\pi G} \int d^3x \sqrt{-g} \left(R + \frac{2}{\ell^2} \right) + \text{boundary terms}$$

- matching of symmetries so far works only up to central term
- reason: did not yet consider specific gravity theory
- focus on Einstein gravity with negative cosmological constant

$$I_{\rm EH} = -\frac{1}{16\pi G} \int d^3x \sqrt{-g} \left(R + \frac{2}{\ell^2} \right) + \text{boundary terms}$$

canonical realization of asymptotic symmetries Brown, Henneaux '86

$$i\{Q[l_n], Q[l_m]\} = (n-m) Q[l_{n+m}] + \frac{\ell}{8G} (n^3 - n) \delta_{n+m,0}$$

- matching of symmetries so far works only up to central term
- reason: did not yet consider specific gravity theory
- focus on Einstein gravity with negative cosmological constant

$$I_{\rm EH} = -\frac{1}{16\pi G} \int d^3x \sqrt{-g} \left(R + \frac{2}{\ell^2} \right) + \text{boundary terms}$$

canonical realization of asymptotic symmetries Brown, Henneaux '86

$$i\{Q[l_n], Q[l_m]\} = (n-m) Q[l_{n+m}] + \frac{\ell}{8G} (n^3 - n) \delta_{n+m,0}$$

▶ calling $Q[l_n] = L_n$ and replacing $i\{,\} \rightarrow [,]$ yields

$$[L_n, L_m] = (n-m) L_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0}$$

Virasoro algebra with central charge

$$c = \frac{3\ell}{2G}$$

- matching of symmetries so far works only up to central term
- reason: did not yet consider specific gravity theory
- focus on Einstein gravity with negative cosmological constant

$$I_{\rm EH} = -\frac{1}{16\pi G} \int d^3x \sqrt{-g} \left(R + \frac{2}{\ell^2} \right) + \text{boundary terms}$$

canonical realization of asymptotic symmetries Brown, Henneaux '86

$$i\{Q[l_n], Q[l_m]\} = (n-m) Q[l_{n+m}] + \frac{\ell}{8G} (n^3 - n) \delta_{n+m,0}$$

▶ calling $Q[l_n] = L_n$ and replacing $i\{,\} \rightarrow [,]$ yields

$$[L_n, L_m] = (n-m) L_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0}$$

Virasoro algebra with central charge

$$c = \frac{3\ell}{2G}$$

matches perfectly with symmetries of CFT₂!

Consequences of symmetry matching

What we have learned so far about AdS_3 is already impressive:

3d Einstein gravity with negative cosmological constant and Brown-Henneaux boundary conditions, if it exists, is dual to a CFT_2 since the physical Hilbert space falls into representations of two Virasoro algebras

Consequences of symmetry matching

What we have learned so far about AdS_3 is already impressive:

3d Einstein gravity with negative cosmological constant and Brown-Henneaux boundary conditions, if it exists, is dual to a CFT_2 since the physical Hilbert space falls into representations of two Virasoro algebras

Consequence: all observables must match between AdS_3 and CFT_2 sides!
What we have learned so far about AdS_3 is already impressive:

3d Einstein gravity with negative cosmological constant and Brown-Henneaux boundary conditions, if it exists, is dual to a CFT_2 since the physical Hilbert space falls into representations of two Virasoro algebras

Consequence: all observables must match between AdS_3 and CFT_2 sides!

 CFT₂ correlation functions calculable on gravity side e.g. 5-point function of stress-tensor flux components

CFT₂:
$$\langle T_{++}(z_1)T_{++}(z_2)T_{++}(z_3)T_{++}(z_4)T_{++}(z_5)\rangle = \frac{4c g_5(\gamma, \zeta)}{\prod_{1 \le i \le 5} z_{ij}}$$

$$\gamma = z_{12}z_{34}/(z_{13}z_{24}) \text{, } \zeta = z_{25}z_{34}/(z_{35}z_{24}) \text{, } z_{ij} = z_i - z_j \text{ and }$$

 $g_5(\gamma,\,\zeta) = \frac{\gamma+\zeta}{2(\gamma-\zeta)} - \frac{\gamma^2-\gamma\zeta+\zeta^2}{\gamma(\gamma-1)\zeta(\zeta-1)(\gamma-\zeta)} \left([\gamma(\gamma-1)+1][\zeta(\zeta-1)+1] - \gamma\zeta \right)$

matches precisely with gravity side

$$\frac{\delta^5 I_{\rm EH}[g_{\mu\nu}]}{\delta g^{++}(z_1)\delta g^{++}(z_2)\delta g^{++}(z_3)\delta g^{++}(z_4)\delta g^{++}(z_5)} = \frac{4c\,g_5(\gamma,\,\zeta)}{\prod_{1\le i\le 5}z_{ij}}$$

What we have learned so far about AdS_3 is already impressive:

3d Einstein gravity with negative cosmological constant and Brown-Henneaux boundary conditions, if it exists, is dual to a CFT_2 since the physical Hilbert space falls into representations of two Virasoro algebras

Consequence: all observables must match between AdS_3 and CFT_2 sides!

- CFT₂ correlation functions calculable on gravity side
- CFT₂ entanglement entropy calculable through length of geodesics



What we have learned so far about AdS_3 is already impressive:

3d Einstein gravity with negative cosmological constant and Brown-Henneaux boundary conditions, if it exists, is dual to a CFT_2 since the physical Hilbert space falls into representations of two Virasoro algebras

Consequence: all observables must match between AdS_3 and CFT_2 sides!

- CFT₂ correlation functions calculable on gravity side
- CFT₂ entanglement entropy calculable through length of geodesics
- Bekenstein–Hawking entropy matches CFT₂ entropy

$$S_{\rm AdS_3} = S_{\rm BH} = \frac{A}{4} = 2\pi \sqrt{\frac{c}{6} \left(M + J\right)} + 2\pi \sqrt{\frac{c}{6} \left(M - J\right)} = S_{\rm Cardy} = S_{\rm CFT_2}$$

What we have learned so far about AdS_3 is already impressive:

3d Einstein gravity with negative cosmological constant and Brown-Henneaux boundary conditions, if it exists, is dual to a CFT_2 since the physical Hilbert space falls into representations of two Virasoro algebras

Consequence: all observables must match between AdS_3 and CFT_2 sides!

- CFT₂ correlation functions calculable on gravity side
- CFT₂ entanglement entropy calculable through length of geodesics
- Bekenstein–Hawking entropy matches CFT₂ entropy

$$S_{\rm AdS_3} = S_{\rm BH} = \frac{A}{4} = 2\pi \sqrt{\frac{c}{6} \left(M + J\right)} + 2\pi \sqrt{\frac{c}{6} \left(M - J\right)} = S_{\rm Cardy} = S_{\rm CFT_2}$$

Holographic dictionary between AdS_3 and CFT_2 observables!

 AdS_{d+1}/CFT_d provides plethora of applications:

can choose to describe observables on gravity or field theory side

- can choose to describe observables on gravity or field theory side
- pick the side that is simpler

- can choose to describe observables on gravity or field theory side
- pick the side that is simpler
- AdS/CFT duality is of strong/weak type

- can choose to describe observables on gravity or field theory side
- pick the side that is simpler
- AdS/CFT duality is of strong/weak type
- complicated calculations in AdS often simple in CFT (and vice versa)

 AdS_{d+1}/CFT_d provides plethora of applications:

- can choose to describe observables on gravity or field theory side
- pick the side that is simpler
- AdS/CFT duality is of strong/weak type
- complicated calculations in AdS often simple in CFT (and vice versa)
- idea 1: map problems in quantum gravity (difficult) to problems in weakly coupled CFT (simple)

example: information loss problem in quantum gravity mapped to information gain problem on CFT side

 AdS_{d+1}/CFT_d provides plethora of applications:

- can choose to describe observables on gravity or field theory side
- pick the side that is simpler
- AdS/CFT duality is of strong/weak type
- complicated calculations in AdS often simple in CFT (and vice versa)
- idea 1: map problems in quantum gravity (difficult) to problems in weakly coupled CFT (simple)
- idea 2: map problems in strongly coupled CFT (difficult) to problems in weakly coupled classical gravity (simple)

example: calculation of shear viscosity over entropy density at strong coupling mapped to gravitational wave absorption of black hole

$$\frac{\eta}{s} = \frac{1}{4\pi} + \text{finite coupling corrections}$$

 AdS_{d+1}/CFT_d provides plethora of applications:

- can choose to describe observables on gravity or field theory side
- pick the side that is simpler
- AdS/CFT duality is of strong/weak type
- complicated calculations in AdS often simple in CFT (and vice versa)
- idea 1: map problems in quantum gravity (difficult) to problems in weakly coupled CFT (simple)
- idea 2: map problems in strongly coupled CFT (difficult) to problems in weakly coupled classical gravity (simple)

If you have a big enough hammer every problem looks like a nail

 AdS_{d+1}/CFT_d provides plethora of applications:

- can choose to describe observables on gravity or field theory side
- pick the side that is simpler
- AdS/CFT duality is of strong/weak type
- complicated calculations in AdS often simple in CFT (and vice versa)
- idea 1: map problems in quantum gravity (difficult) to problems in weakly coupled CFT (simple)
- idea 2: map problems in strongly coupled CFT (difficult) to problems in weakly coupled classical gravity (simple)

If you have a big enough hammer every problem looks like a nail

Black hole holography currently seems like a big hammer! Black hole holography harmonic oscillator of 21st century?

Applications:

heavy ion collisions (e.g. ALICE @ LHC)

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)
- chaotic quantum mechanical models (e.g. JT/SYK correspondence)

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)
- chaotic quantum mechanical models (e.g. JT/SYK correspondence)
 Open issues:

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)
- chaotic quantum mechanical models (e.g. JT/SYK correspondence)
 Open issues:
 - how to use holography to construct generic black hole microstates?

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)
- chaotic quantum mechanical models (e.g. JT/SYK correspondence) Open issues:
 - how to use holography to construct generic black hole microstates?
 - how to reconstruct spacetime from CFT data?

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)
- chaotic quantum mechanical models (e.g. JT/SYK correspondence) Open issues:
 - how to use holography to construct generic black hole microstates?
 - how to reconstruct spacetime from CFT data?
 - how to prove AdS/CFT?

Applications:

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)
- chaotic quantum mechanical models (e.g. JT/SYK correspondence) Open issues:
 - how to use holography to construct generic black hole microstates?
 - how to reconstruct spacetime from CFT data?
 - how to prove AdS/CFT?

Generalizations:

Applications:

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)
- chaotic quantum mechanical models (e.g. JT/SYK correspondence) Open issues:
 - how to use holography to construct generic black hole microstates?
 - how to reconstruct spacetime from CFT data?
 - how to prove AdS/CFT?

Generalizations:

How general is holography?

Applications:

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)
- chaotic quantum mechanical models (e.g. JT/SYK correspondence) Open issues:
 - how to use holography to construct generic black hole microstates?
 - how to reconstruct spacetime from CFT data?
 - how to prove AdS/CFT?

Generalizations:

- How general is holography?
- Does it work beyond AdS/CFT?

Applications:

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)
- chaotic quantum mechanical models (e.g. JT/SYK correspondence) Open issues:
 - how to use holography to construct generic black hole microstates?
 - how to reconstruct spacetime from CFT data?
 - how to prove AdS/CFT?

Generalizations:

- How general is holography?
- Does it work beyond AdS/CFT?
- In particular, does it work in flat space or de Sitter?

Applications:

- heavy ion collisions (e.g. ALICE @ LHC)
- strongly correlated cond-mat systems (e.g. strange metals)
- chaotic quantum mechanical models (e.g. JT/SYK correspondence) Open issues:
 - how to use holography to construct generic black hole microstates?
 - how to reconstruct spacetime from CFT data?
 - how to prove AdS/CFT?

Generalizations:

- How general is holography?
- Does it work beyond AdS/CFT?
- In particular, does it work in flat space or de Sitter?

Will need more research on black holes to resolve all issues!

Black holes II prepares you for some of the research directions

Thanks for your attention...











EHT collaboration, April 2019