

Unitarity in flat space holography

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based on work with
Rosseel and Riegler, [1403.5297](#)
and work with Afshar, Bagchi, Detournay, Fareghbal, Simon

Motivation

- ▶ Holographic principle, if correct, must work beyond AdS/CFT
holographic principle: 't Hooft '93; Susskind '94

AdS/CFT precursor: Brown, Henneaux '86

AdS/CFT: Maldacena '97; Gubser, Klebanov, Polyakov '98; Witten '98

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- ▶ Does it work in flat space?

Polchinski '99

Susskind '99

Giddings '00

Gary, Giddings, Penedones '09; Gary, Giddings '09; ...

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- ▶ Does it work in flat space?
- ▶ Can we find models realizing flat space/field theory correspondences?

Barnich, Compere '06

Barnich et al. '10-'14

Bagchi et al. '10-'14

Strominger et al. '13-'14

...

flat space chiral gravity: Bagchi, Detournay, DG '12

Motivation

- ▶ Holographic principle, if correct, must work beyond AdS/CFT
- ▶ Does it work in flat space?
- ▶ Can we find models realizing flat space/field theory correspondences?
- ▶ Are there higher-spin versions of such models?

Afshar, Bagchi, Fareghbal, DG, Rosseel '13

Gonzalez, Matulich, Pino, Troncoso '13

part of larger program: non-AdS holography in higher spin gravity

Gary, DG, Rashkov '12

Afshar, Gary, DG, Rashkov, Riegler '12

Gutperle, Hijano, Samani '13

Gary, DG, Prohazka, Rey (in prep.) '14

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- ▶ Does this correspondence emerge as limit of (A)dS/CFT?

some aspects: yes; other aspects: no

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- ▶ (When) are these models unitary?

this talk!

see also Barnich, Oblak '14 (induced representations of BMS_3)

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- ▶ Is there an analog of Hawking–Page phase transition?

yes: Bagchi, Detournay, DG, Simon '13
Detournay, DG, Schöller, Simon '14

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Strominger '13

He, Lysov, Mitra, Strominger '14

Banks '14

Cachazo, Strominger '14

- ▶ ...

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Address **unitarity** question in two boundary dimensions!
Particular interest: **unitarity** in flat space higher spin gravity?

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Work mostly on CFT side and study “landscape” of possible flat space asymptotic symmetry algebras:

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In this talk I will address 1.-5., but not 6.!

Main results

1. NO-GO:

Generically (see later) you can have only two out of three:

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- ▶ Flat space
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Example:

Flat space chiral gravity

Bagchi, Detournay, DG, 1208.1658

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Example:

Minimal model holography

Gaberdiel, Gopakumar, 1011.2986, 1207.6697

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Flat space higher spin gravity (Galilean W_3 algebra)

Afshar, Bagchi, Fareghbal, DG and Rosseel, 1307.4768

Gonzalez, Matulich, Pino and Troncoso, 1307.5651

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2. YES-GO:

There is (at least) one counter-example, namely a Vasiliev-type of theory, where you can have all three properties!

Unitary, non-trivial flat space higher spin algebra exists!
Vasiliev-type flat space chiral higher spin gravity?



İnönü–Wigner contraction of Virasoro (Barnich & Compère '06)

BMS₃ and GCA₂/URCA₂

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- ▶ Define superrotations L_n and supertranslations M_n

$$\text{GCA:} \quad L_n := \mathcal{L}_n + \bar{\mathcal{L}}_n \quad M_n := -\epsilon (\mathcal{L}_n - \bar{\mathcal{L}}_n)$$

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- ▶ Galilean limit/ultrarelativistic boost: $\epsilon \sim 1/\ell \rightarrow 0$

$$[L_n, L_m] = (n - m) L_{n+m} + c_L \frac{1}{12} \delta_{n+m,0}$$

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- ▶ Example: URCA TMG: $c_L = 3/\mu G$ and $c_M = 3/G$ [dimensionful!]

(Non-)unitarity in GCAs

Assumptions:

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Conclusions:

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Generalization to higher spin asymptotic symmetry algebras

Facts:

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Example: Galilean W_3 algebra: $[L_n, L_m]$ and $[L_n, M_m]$ as before,

$$[L_n, U_m] = (2n - m)U_{n+m} \quad [L_n, V_m] = [M_n, U_m] = (2n - m)V_{n+m}$$

$$[U_n, U_m] = (n - m)(2n^2 + 2m^2 - nm - 8)L_{n+m} + \frac{192}{c_M} (n - m)\Lambda_{n+m} \\ - \frac{96(c_L + \frac{44}{5})}{c_M^2} (n - m)\Theta_{n+m} + \frac{c_L}{12} n(n^2 - 1)(n^2 - 4) \delta_{n+m,0}$$

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Unitarity leads to further contraction! (Afshar et al. 1307.4768)

Inönü–Wigner contraction of two W -algebras (generators \mathcal{L}, \mathcal{W}):

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Doubly contracted algebra has unitary representations:

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Higher spin states decouple and become null states!

How general is this decoupling mechanism? (DG, Riegler, Rosseel '14)

Same conclusions — all higher spin states become null states — for

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How general is this decoupling mechanism? (DG, Riegler, Rosseel '14)

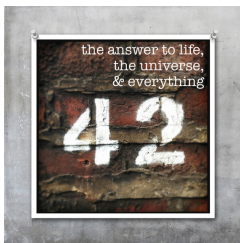
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Upper bound from current algebra part:

$$\hat{\mathfrak{u}}(1) : \quad [O_n, O_m] = \frac{2(54 - c_L)}{3} n \delta_{n+m,0}$$

$$\hat{\mathfrak{su}}(2) : \quad [Q_n^a, Q_m^b] = (a - b) Q_{n+m}^{a+b} + \frac{42 - c_L}{6} n \delta_{a+b,0} \delta_{n+m,0}$$

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Lower bound: non-negativity of c_{bare} in

$$c_L = c_{\hat{u}(1)} + c_{\hat{\mathfrak{su}}(2)} + c_{\text{bare}}$$

with $c_{\dots} = k \dim g / (k + h^\vee)$ with level $k = (42 - c_L)/6$, thus:

$$c_{\hat{u}(1)} = 1 \quad c_{\hat{\mathfrak{su}}(2)} = 3(42 - c_L)/(54 - c_L)$$

$\Rightarrow c_{\text{bare}} \geq 0$ implies $c_L \geq 29 - \sqrt{661} \approx 3.29$

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In all cases above direct consequence of non-linearity in W-algebra!

General NO-GO result

See 1403.5297 for details!

Idea of no-go proof:

- ▶ Assume non-linearity in W-algebra, e.g. $(\lim_{c \rightarrow \infty} f(c) \rightarrow 0)$

$$[W_n, W_m] = \dots + f(c) :AB:_{n+m} + \dots + \omega(c) \prod_{j=-(s-1)}^{s-1} (n+j) \delta_{n+m,0}$$

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- ▶ Define Galilean contraction in usual way, e.g.

$$U_n := W_n + \bar{W}_n$$

$$C_n := A_n + \bar{A}_n$$

$$E_n := B_n + \bar{B}_n$$

$$V_n := -\epsilon (W_n - \bar{W}_n)$$

$$D_n := -\epsilon (A_n - \bar{A}_n)$$

$$F_n := -\epsilon (B_n - \bar{B}_n)$$

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No central terms in higher spin generators:

$$[V_n, V_m] = 0 \quad (\text{trivial})$$

$$[U_n, V_m] = \dots + \# c_M \delta_{n+m,0} \quad (\text{dimensions!})$$

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$$\begin{aligned} [U_n, U_m] = & \dots + \mathcal{O}(1/c_M) (:CF:_{n+m} + :DE:_{n+m}) \\ & + \underbrace{\mathcal{O}\left(\frac{1}{c_M^2}\right)}_{\neq 0} :DF:_{n+m} + \tilde{\omega}(c_L) \prod_{j=-(s-1)}^{s-1} (n+j) \delta_{n+m,0}, \end{aligned}$$

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- ▶ Central terms do not survive this rescaling
- ▶ This is why higher spin states become null states and decouple!

YES-GO

... every no-go result is only as good as its premises!

Drop assumption of non-linearity!

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- ▶ Linear higher spin algebra: Pope–Romans–Shen W_∞ algebra!

$$[V_m^i, V_n^j] = \sum_{r=0}^{\lfloor \frac{i+j}{2} \rfloor} g_{2r}^{ij}(m, n) V_{m+n}^{i+j-2r} + c^i(m) \delta^{ij} \delta_{m+n,0}$$

note: wedge algebra is $\mathfrak{hs}(1)$

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- ▶ Ultra-relativistic contraction of generators

$$\mathcal{V}_m^i = V_m^i - \bar{V}_{-m}^i, \quad \mathcal{W}_m^i = \epsilon (V_m^i + \bar{V}_{-m}^i)$$

and central charges

$$c_{\mathcal{V}} = c - \bar{c}, \quad c_{\mathcal{W}} = \epsilon (c + \bar{c})$$

The Treachery of Algebras — Ceci n'est pas une théorie.



Ceci n'est pas une pipe.

Magritte

YES-GO

Asymptotic symmetry algebra of flat space **chiral** higher spin gravity

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- ▶ If it exists, this must be its asymptotic symmetry algebra:

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where

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- ▶ Vacuum descendants $\mathcal{W}_m^i |0\rangle$ are null states for all i and m !
- ▶ AdS parent theory: no trace anomaly, but **gravitational anomaly** (Like in conformal Chern–Simons gravity \rightarrow Vasiliev type analogue?)

Summary and outlook

Main results:

- ▶ NO-GO: Generically, only two out of three are possible: unitarity, flat space, non-trivial higher spin states
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- ▶ Same argument eliminates also multi-graviton excitations

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Interpretation: Perhaps flat space + unitarity allows only (specific) Vasiliev-type of theories and no truncation to finite spin?

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- ▶ Other definitions of vacuum/hermitian conjugation that lead to unitarity? (suggestive: Barnich, Oblak, '14)
- ▶ Other linear higher spin algebras? ($W_{1+\infty}$; what else?)
- ▶ Construction of unitary flat space chiral higher spin gravity ($FS_{\chi}HSG_3$)?

Evidence so far for unitary $FS_{\chi}HSG_3$:
constructed its asymptotic symmetry algebra!

... hopefully shed **Empire of Light** on existence of $FS_{\chi}HSG_3$ in the future!



Thanks for your attention!

-  D. Grumiller, M. Riegler and J. Rosseel, “Unitarity in three-dimensional flat space higher spin theories,” [arXiv:1403.5297](#).
-  H. Afshar, A. Bagchi, R. Fareghbal, D. Grumiller and J. Rosseel, Phys. Rev. Lett. **111** (2013) 121603 [[arXiv:1307.4768](#)].
-  A. Bagchi, S. Detournay, D. Grumiller and J. Simon, Phys. Rev. Lett. **111** (2013) 181301 [[arXiv:1305.2919](#)].
-  A. Bagchi, S. Detournay and D. Grumiller, Phys. Rev. Lett. **109** (2012) 151301 [[arXiv:1208.1658](#)].

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