

1110.5918 Taronna-EJ : Massless

1203.6578 Lopez-Taronna-EJ : Massive & Partially-massless

Cubic interactions of **massive** and **(partially-)massless** **HS fields in (A)dS**

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HS Interactions

● Massless HS

- Gauge Invariance
- No-go Weinberg'64
Aragon-Deser ...
- Yes-go Berends-Burgers-van Dam
Fradkin-Vasiliev Metsaev ...
- Vasiliev's Equations

● Massive HS

- Reproduce Fierz sys.
- No-go Fierz-Pauli
Velo-Zwanziger ...
- Yes-go Federbush
Argyres-Nappi ...
- String Theory

*Systematic approach to
massless and **massive** HS interactions
in AdS and flat background*

Noether Procedure

$$S_{\text{YM}} = \int \partial^2 A^2 + \textcolor{blue}{g} \int \partial A^3 + \textcolor{blue}{g}^2 \int A^4$$

$$g_2 \int \partial^3 A^3 \rightarrow F^3$$

$$S_{\text{GR}} = \int \partial^2 h^2 + \textcolor{red}{G}^{\frac{1}{2}} \int \partial^2 h^3 + \textcolor{red}{G} \int \partial^2 h^4 + \dots$$

$$g_2 \int \partial^4 h^3 \rightarrow \sqrt{-g} R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma}$$

$$g_3 \int \partial^6 h^3 \rightarrow \sqrt{-g} R^\mu{}_\nu R^\nu{}_\rho R^\rho{}_\mu$$

$$S = \textcolor{green}{S^{(2)}} + \textcolor{green}{S^{(3)}} + \textcolor{green}{S^{(4)}} + \dots$$

$$\delta \varphi = \textcolor{green}{\delta^{(0)} \varphi} + \textcolor{green}{\delta^{(1)} \varphi} + \textcolor{green}{\delta^{(2)} \varphi} + \dots$$

Gauge Invariance

$$\delta S = 0 \Rightarrow \left\{ \begin{array}{l} \delta^{(0)} \textcolor{green}{S^{(2)}} = 0 \\ \delta^{(0)} \textcolor{green}{S^{(3)}} + \cancel{\delta^{(1)} \textcolor{green}{S^{(2)}}} = 0 \\ \vdots \end{array} \right. \approx 0$$

Cubic Interactions

- **Massless HS**

- Gauge Invariance

- **Massive HS**

- Reproduce Fierz sys.

- Correct Propagation of DoF
- Light-Cone(LC) Gauge Metsaev 2006
 - Complete list of couplings
 - Massive & Massless , Cubic , Flat

HS in (A)dS

- Only **Transverse & Traceless (TT) part**
 - $LC \subset TT \subset \text{“Full”}$
 - Key to “Full” Manvelyan-Mkrtchyan-Ruehl Taronna-Sagnotti
 - On-Shell Info (n-pt fns , ...)
- Only Symmetric Bosons
- Only Parity-Invariant Interactions
- Any Dimensions $d \geq 4$
- Metric-like formulation



OUTLINE

- Massless HS in Flat-Space (warm-up ex.)
- Ambient-Space Formalism for (A)dS
 - Free HS fields
 - Interactions
- Cubic Interactions
 - Massive & Massless in (A)dS
 - Partially-Massless in dS

Warm-up Ex

Massless HS in Flat Space

- *Free Action*

$$S^{(2)} = - \int \varphi^\mu \cdots \square \varphi_\mu \cdots + \dots$$

- **Field Eqs**

$$\square \varphi + \dots \approx 0$$

Fronsdal

- Constrained

Francia-Mourad-Sagnotti

- Unconstrained

Buchbinder-Galajinsky-Krykhtin

divergence, trace,
auxiliary fields

- *Cubic Action*

$$S^{(3)} = \int c (\partial^{\cdots} \varphi^{\cdots}) (\partial^{\cdots} \varphi^{\cdots}) (\partial^{\cdots} \varphi \cdots)$$

- *Noether Procedure*

$$\delta^{(0)} \varphi_{\mu\nu\rho\cdots} = \partial_{(\mu} \varepsilon_{\nu\rho\cdots)}$$

$$\delta^{(0)} S^{(3)} = \int c (\partial^{\cdots} \varepsilon^{\cdots}) (\partial^{\cdots} \varphi^{\cdots}) (\partial^{\cdots} [\square \varphi \cdots + \dots])$$

Warm-up Ex

Massless HS in Flat Space

- Focus TT part

$$S^{(3)} = \int c (\partial^{\cdots} \varphi^{\cdots})(\partial^{\cdots} \varphi^{\cdots})(\partial^{\cdots} \varphi \dots)$$

$$= \boxed{TT \text{ in } \varphi} + \boxed{Others}$$

$$\delta^{(0)} S^{(3)} = \int c (\partial^{\cdots} \cdots) (\partial^{\cdots} \varphi^{\cdots}) (\partial^{\cdots} [\square \varphi \dots + \cdots])$$

$$= \boxed{TT \text{ in } \varphi \& \epsilon} \approx_0 + \boxed{Others} \approx_0$$

- Noether Procedure \rightarrow $\boxed{\delta^{(0)} S^{(3)}_{TT}} \approx 0$

Warm-up Ex

Massless HS in Flat Space

◦ Cubic Action

Taronna-EJ

$$\begin{aligned} y_1 &:= \partial_{u_1} \cdot \partial_{x_2} & y_2 &:= \partial_{u_2} \cdot \partial_{x_3} & y_3 &:= \partial_{u_3} \cdot \partial_{x_1} \\ z_1 &:= \partial_{u_2} \cdot \partial_{u_3} & z_2 &:= \partial_{u_3} \cdot \partial_{u_1} & z_3 &:= \partial_{u_1} \cdot \partial_{u_2} \end{aligned}$$

$$S^{(3)} \sim \int C(\partial_{u_i} \cdot \partial_{x_j}, \partial_{u_i} \cdot \partial_{u_j}) \varphi(x_1, u_1) \varphi(x_2, u_2) \varphi(x_3, u_3) \Big|_{\substack{x_i=x \\ u_i=0}}$$

Divergent

$$\partial_u \cdot \partial_x \varphi(x, u) = \frac{1}{(s-1)!} \partial^{\mu_s} \varphi_{\mu_s \mu_1 \dots \mu_{s-1}}(x) u^{\mu_1} \dots u^{\mu_{s-1}}$$

Trace

$$\partial_u^2 \varphi(x, u) = \frac{1}{(s-2)!} \varphi'_{\mu_1 \dots \mu_{s-2}}(x) u^{\mu_1} \dots u^{\mu_{s-2}}$$

$$\varphi(x, u) = \frac{1}{s!} \varphi_{\mu_1 \dots \mu_s}(x) u^{\mu_1} \dots u^{\mu_s}$$

Warm-up Ex

Massless HS in Flat Space

◦ Cubic Action

Taronna-EJ

$$\begin{aligned} y_1 &:= \partial_{u_1} \cdot \partial_{x_2} & y_2 &:= \partial_{u_2} \cdot \partial_{x_3} & y_3 &:= \partial_{u_3} \cdot \partial_{x_1} \\ z_1 &:= \partial_{u_2} \cdot \partial_{u_3} & z_2 &:= \partial_{u_3} \cdot \partial_{u_1} & z_3 &:= \partial_{u_1} \cdot \partial_{u_2} \end{aligned}$$

$$S^{(3)} \sim \int C(y_i, z_i) \left| \varphi(x_1, u_1) \varphi(x_2, u_2) \varphi(x_3, u_3) \right|_{\substack{x_i=x \\ u_i=0}}$$

$$\delta^{(0)} S^{(3)} \sim \int (y_3 \partial_{z_2} - y_2 \partial_{z_3}) C(y, z) \varepsilon(x_1, u_1) \varphi(x_2, u_2) \varphi(x_3, u_3) \Big|_{\substack{x_i=x \\ u_i=0}}$$

$$(y_3 \partial_{z_2} - y_2 \partial_{z_3}) C(y, z) = 0$$

Gradient

$$u \cdot \partial_x \varepsilon(x, u) = \frac{1}{s!} \partial_{(\mu_1} \varepsilon_{\mu_2 \dots \mu_s)}(x) u^{\mu_1} \dots u^{\mu_s}$$

Warm-up Ex

Massless HS in Flat Space

- *Cubic Action*

Taronna-EJ

$$\begin{aligned} y_1 &:= \partial_{u_1} \cdot \partial_{x_2} & y_2 &:= \partial_{u_2} \cdot \partial_{x_3} & y_3 &:= \partial_{u_3} \cdot \partial_{x_1} \\ z_1 &:= \partial_{u_2} \cdot \partial_{u_3} & z_2 &:= \partial_{u_3} \cdot \partial_{u_1} & z_3 &:= \partial_{u_1} \cdot \partial_{u_2} \end{aligned}$$

$$S^{(3)} \sim \int C(y, z) \varphi(x_1, u_1) \varphi(x_2, u_2) \varphi(x_3, u_3) \Big|_{\substack{x_i=x \\ u_i=0}}$$

Gauge
invariance

$$(y_i \partial_{z_{i+1}} - y_{i+1} \partial_{z_i}) C(y, z) = 0$$

General
solution

$$C = \mathcal{K}(y_1, y_2, y_3, g)$$

$$g := y_1 z_1 + y_2 z_2 + y_3 z_3$$

- **Generating fn of all consistent couplings**
- **Go to LC → Metsaev's results**
- **Relax “~” → Full vertices**

Generalization to $(A)dS$

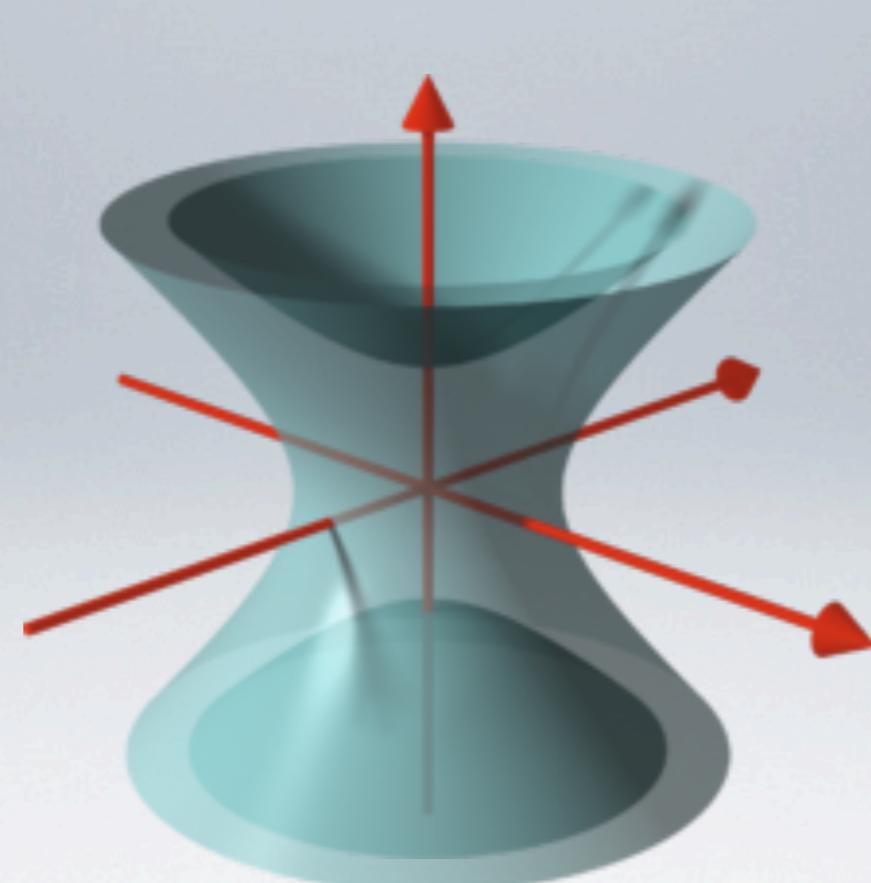
$$\boxed{\partial_x \Rightarrow \nabla_x} ?$$

AMBIENT SPACE

Fronsdal Metsaev Biswas-Siegel ...

de Sitter

$$X^2 = L^2$$



$$\mathbb{R}^{1,d}$$

Anti de Sitter

$$X^2 = -L^2$$

AMBIENT SPACE

Fronsdal Metsaev Biswas-Siegel ...

- Ambient-Space HS Fields $\Phi(X, U) = R^{\Delta_h} \varphi(x, u)$

- Homogeneity

$$X \cdot \partial_X \Phi = \Delta_h \Phi$$

$$X = (R, x)$$

- Tangentiality

$$X \cdot \partial_U \Phi = 0$$

$$U = (v, u)$$

- **MASSLESS** equation in **Ambient Space**

- **massive** & **(partially-)massless** in **(A)dS**

depending on Δ_h

- Gauge symmetry

$$\delta^{(0)} \Phi = U \cdot \partial_X E$$

AMBIENT SPACE

- Consistency of Gauge sym. with Tangentiality & Homogeneity

Lopez-Taronna-EJ

$$\delta^{(0)} \Phi = U \cdot \partial_X E \left[\begin{array}{l} X \cdot \partial_U \Phi = 0 \\ (X \cdot \partial_X - U \cdot \partial_U + \mu + 2) \Phi = 0 \\ (X \cdot \partial_X - U \cdot \partial_U + \mu) E = 0 \end{array} \right]$$

- Massless : $\mu = 0$ $X \cdot \partial_U E = 0$
- Massive : $\mu \notin \{0, \dots, s-1\}$ $E = 0$
- Partially-massless : $\mu \in \{1, \dots, s-1\}$ $E = (U \cdot \partial_X)^\mu \Omega$
 $X \cdot \partial_U \Omega = 0$

AMBIENT SPACE

- Gauge Symmetry $(X \cdot \partial_X - U \cdot \partial_U + \mu + 2) \Phi = 0$
- Massless $(A)dS : \mu = 0$
$$\delta^{(0)}\Phi = U \cdot \partial_X E$$
- Massive $(A)dS : \mu \notin \{0, \dots, s-1\}$
(no gauge sym.)
- Partially-
-massless $dS : \mu \in \{1, \dots, s-1\}$
$$\delta^{(0)}\Phi = (U \cdot \partial_X)^{\mu+1} \Omega$$

How to deal with the **ACTION** ?

Radial dependence of the Lagrangian

$$\mathsf{L}_{(s,\mu)}^{(2)} = R^{2(s-\mu-3)} \mathcal{L}_{(s,\mu)}^{(2)}$$

$$\mathsf{L}_{(s,\mu)}^{(2)} = \Phi \cdot \square \Phi$$

$$\mathcal{L}_{(s,\mu)}^{(2)} = \varphi \cdot (\square_{(A)dS} + m_\mu^2) \varphi$$

$$S_{(s,\mu)}^{(2)} = \left(\int_0^\infty dR R^{1+2(s-\mu-3)} \right) S_{(s,\mu)}^{(2)}$$

HS Cubic Interactions in (A)dS

- ❖ How to deal with the R -dependence?
 - ❖ Remove it by introducing $\delta(\sqrt{\epsilon X^2} - L)$
- ❖ Then, what are the differences w.r.t. the (A)dS intrinsic way?
 - ❖ We still work with ∂_X not ∇_x
- ❖ Only simplifications? No subtleties?
 - ❖ A subtlety arises in “TOTAL DERIVATIVES”

- Cubic Action in (A)dS

$$\begin{array}{lll} Y_1 = \partial_{U_1} \cdot \partial_{X_2} & Y_2 = \partial_{U_2} \cdot \partial_{X_3} & Y_3 = \partial_{U_3} \cdot \partial_{X_1} \\ Z_1 = \partial_{U_2} \cdot \partial_{U_3} & Z_2 = \partial_{U_3} \cdot \partial_{U_1} & Z_3 = \partial_{U_1} \cdot \partial_{U_2} \end{array}$$

$$S^{(3)} \sim \int d^{d+1}X \; \delta\left(\sqrt{\epsilon X^2} - L\right) C(Y_i, Z_i) \times \\ \times \Phi(X_1, U_1) \Phi(X_2, U_2) \Phi(X_3, U_3) \Big|_{\substack{X_i=X \\ U_i=0}}$$

- Massless Gauge Symmetry

$$\delta^{(0)} S^{(3)} \sim \int d^{d+1}X \; \delta\left(\sqrt{\epsilon X^2} - L\right) \left[C(Y_i, Z_i), \boldsymbol{U}_1 \cdot \boldsymbol{\partial}_{\boldsymbol{X}_1} \right] \times \\ \times \boldsymbol{E}(\boldsymbol{X}_1, \boldsymbol{U}_1) \Phi(X_2, U_2) \Phi(X_3, U_3) \Big|_{\substack{X_i=X \\ U_i=0}}$$

- Partially-massless Gauge Symmetry

$$\delta^{(0)} S^{(3)} \sim \int d^{d+1}X \; \delta\left(\sqrt{\epsilon X^2} - L\right) \left[C(Y_i, Z_i), (\boldsymbol{U}_1 \cdot \boldsymbol{\partial}_{\boldsymbol{X}_1})^{r+1} \right] \times \\ \times \boldsymbol{\Omega}(\boldsymbol{X}_1, \boldsymbol{U}_1) \Phi(X_2, U_2) \Phi(X_3, U_3) \Big|_{\substack{X_i=X \\ U_i=0}}$$

• Cubic Action in (A)dS

$$\begin{array}{lll} Y_1 = \partial_{U_1} \cdot \partial_{X_2} & Y_2 = \partial_{U_2} \cdot \partial_{X_3} & Y_3 = \partial_{U_3} \cdot \partial_{X_1} \\ Z_1 = \partial_{U_2} \cdot \partial_{U_3} & Z_2 = \partial_{U_3} \cdot \partial_{U_1} & Z_3 = \partial_{U_1} \cdot \partial_{U_2} \end{array}$$

$$S^{(3)} \sim \int d^{d+1}X \, \delta(\sqrt{\epsilon X^2} - L) \, C(Y_i, Z_i) \times \\ \times \Phi(X_1, U_1) \Phi(X_2, U_2) \Phi(X_3, U_3) \Big|_{\substack{X_i=X \\ U_i=0}}$$

• Massless Gauge Symmetry

$$\delta^{(0)} S^{(3)} \left[C(Y, Z), U_1 \cdot \partial_{X_1} \right] = (Y_3 \partial_{Z_2} - Y_2 \partial_{Z_3}) C(Y, Z) \times \\ \times E(X_1, U_1 + (\text{total derivatives}), U_3) \Big|_{\substack{X_i=X \\ U_i=0}}$$

• Partially-massless Gauge Symmetry

$$\delta^{(0)} S^{(3)} \left[C(Y, Z), (U_1 \cdot \partial_{X_1})^{r+1} \right] = (Y_3 \partial_{Z_2} - Y_2 \partial_{Z_3})^{r+1} C(Y, Z) \times \\ \times \Omega(X_1, U_1 + (\text{total derivatives}), U_3) \Big|_{\substack{X_i=X \\ U_i=0}}$$

- * Total Derivatives

$$\delta(\sqrt{X^2} - L) \partial_X \cdot (-) \Big|_{x_i=x} = -\frac{1}{L} \delta'(\sqrt{X^2} - L) X \cdot (-) \Big|_{x_i=x}$$

The diagram illustrates the derivation of a total derivative. The expression is $\delta(\sqrt{X^2} - L) \partial_X \cdot (-) \Big|_{x_i=x}$. A red circle highlights the term $X \cdot (-)$. Two red arrows point from the terms $X_i \cdot \partial_{X_i}$ and $X_i \cdot \partial_{U_i}$ below it towards the highlighted term, indicating they are being moved to act on it.

- * Homogeneity and Tangentiality
- * Move them to the right, and Act them on fields
 - * Use commutation relations with Y_i and Z_i

- Cubic Action in (A)dS

$$\begin{array}{lll} Y_1 = \partial_{U_1} \cdot \partial_{X_2} & Y_2 = \partial_{U_2} \cdot \partial_{X_3} & Y_3 = \partial_{U_3} \cdot \partial_{X_1} \\ Z_1 = \partial_{U_2} \cdot \partial_{U_3} & Z_2 = \partial_{U_3} \cdot \partial_{U_1} & Z_3 = \partial_{U_1} \cdot \partial_{U_2} \end{array}$$

$$S^{(3)} \sim \int d^{d+1} X \; \delta\big(\sqrt{\epsilon \, X^2} - L\big) \; C(Y_i,Z_i) \times \\ \qquad \qquad \qquad \times \Phi(X_1,U_1) \, \Phi(X_2,U_2) \, \Phi(X_3,U_3) \; \Big|_{\substack{X_i=X \\ U_i=0}}$$

- Massless Gauge Symmetry

$$\left[Y_2 \partial_{Z_3} - Y_3 \partial_{Z_2} + \frac{\hat{\delta}}{L} \left(Y_2 \partial_{Y_2} - Y_3 \partial_{Y_3} - \tfrac{\mu_2-\mu_3}{2} \right) \partial_{Y_1} \right] C(Y,Z) = 0$$

$$\delta^{(n)}\big(\sqrt{\epsilon \, X^2} - L\big) \equiv \delta\big(\sqrt{\epsilon \, X^2} - L\big) \, (\epsilon \, \hat{\delta})^n$$

- Partially-massless Gauge Symmetry

$$\sum_{l_1+l_2+l_3=r_1+1} \binom{r_1+1}{l_1 \, l_2 \, l_2} \left(Y_3 \, \partial_{Y_3} - Y_2 \, \partial_{Y_2} - Z_3 \, \partial_{Z_3} + \tfrac{1}{2} \, (\mu_2 - \mu_3) \right)_{l_1} \\ \qquad \qquad \qquad \times \; \hat{\delta}^{l_1} \, Y_3^{l_2} \, \left(-Y_2 + \hat{\delta} \, Z_3 \, \partial_{Y_1} \right)^{l_3} \, \partial_{Y_1}^{l_1} \, \partial_{Z_2}^{l_2} \, \partial_{Z_3}^{l_3} \, C(Y,Z) = 0$$

- How to solve

$$\sum_{l_1+l_2+l_3=r_1+1} \binom{r_1+1}{l_1 l_2 l_3} \left(Y_3 \partial_{Y_3} - Y_2 \partial_{Y_2} - Z_3 \partial_{Z_3} + \frac{1}{2} (\mu_2 - \mu_3) \right)_{l_1} \\ \times \hat{\delta}^{l_1} Y_3^{l_2} \left(-Y_2 + \hat{\delta} Z_3 \partial_{Y_1} \right)^{l_3} \partial_{Y_1}^{l_1} \partial_{Z_2}^{l_2} \partial_{Z_3}^{l_3} C(Y, Z) = 0$$

- Straightforward for given spins: $s_1 - s_2 - s_3$

- Expand the coupling in $Y_1^{m_1} Y_2^{m_2} Y_3^{m_3} Z_1^{n_1} Z_2^{n_2} Z_3^{n_3}$ with arbitrary coeff.

- System of linear equations for those coeff.

$$\begin{bmatrix} \cdot & \cdots & \cdot \\ \vdots & \ddots & \vdots \\ \cdot & \cdots & \cdot \end{bmatrix}^N_M \quad N, M \sim s_1 s_2 s_3$$

- Numerical methods **Mathematica**

* Ex) 4-4-2 Couplings

- * Spin 4 : first partially-massless point ($\mu=1$)
- * Spin 2 : massless ($\mu=0$)

$$\begin{aligned}
C_1 &= Y_1^4 Y_2^4 Y_3^2 - 12 \hat{\delta}^2 Y_1^2 Y_2^2 (Y_1 Z_1 + Y_2 Z_2)^2 + 48 \hat{\delta}^3 Y_1 Y_2 (Y_1 Z_1 + Y_2 Z_2) Z_3 (2 Y_1 Z_1 + 2 Y_2 Z_2 + Y_3 Z_3) \\
&\quad - 24 \hat{\delta}^4 Z_3^2 [6 Y_1^2 Z_1^2 + 6 Y_2^2 Z_2^2 + 4 Y_2 Y_3 Z_2 Z_3 + Y_3^2 Z_3^2 + 2 Y_1 Z_1 (7 Y_2 Z_2 + 2 Y_3 Z_3)] + 96 \hat{\delta}^5 Z_1 Z_2 Z_3^3, \\
C_2 &= Y_1^3 Y_2^3 Y_3^2 Z_3 - 3 \hat{\delta} Y_1^2 Y_2^2 (Y_1 Z_1 + Y_2 Z_2)^2 + 12 \hat{\delta}^2 Y_1 Y_2 (Y_1 Z_1 + Y_2 Z_2) Z_3 (2 Y_1 Z_1 + 2 Y_2 Z_2 + Y_3 Z_3) \\
&\quad - 6 \hat{\delta}^3 Z_3^2 [6 Y_1^2 Z_1^2 + 6 Y_2^2 Z_2^2 + 4 Y_2 Y_3 Z_2 Z_3 + Y_3^2 Z_3^2 + 2 Y_1 Z_1 (7 Y_2 Z_2 + 2 Y_3 Z_3)] + 24 \hat{\delta}^4 Z_1 Z_2 Z_3^3, \\
C_3 &= Y_1^3 Y_2^3 Y_3 (Y_1 Z_1 + Y_2 Z_2) + \hat{\delta} Y_1^2 Y_2^2 (6 Y_1^2 Z_1^2 + 11 Y_1 Y_2 Z_1 Z_2 + 6 Y_2^2 Z_2^2) \\
&\quad - 18 \hat{\delta}^2 Y_1 Y_2 (Y_1 Z_1 + Y_2 Z_2) Z_3 (2 Y_1 Z_1 + 2 Y_2 Z_2 + Y_3 Z_3) \\
&\quad + 6 \hat{\delta}^3 Z_3^2 [6 Y_1^2 Z_1^2 + 2 Y_2 Z_2 (3 Y_2 Z_2 + Y_3 Z_3) + Y_1 Z_1 (15 Y_2 Z_2 + 2 Y_3 Z_3)] - 12 \hat{\delta}^4 Z_1 Z_2 Z_3^3, \\
C_4 &= - Y_1^2 Y_2^2 (Y_1^2 Z_1^2 + 2 Y_1 Y_2 Z_1 Z_2 + Y_2^2 Z_2^2 - Y_3^2 Z_3^2) \\
&\quad + 4 \hat{\delta} Y_1 Y_2 (Y_1 Z_1 + Y_2 Z_2) Z_3 (2 Y_1 Z_1 + 2 Y_2 Z_2 + Y_3 Z_3) \\
&\quad - 2 \hat{\delta}^2 Z_3^2 [6 Y_1^2 Z_1^2 + 6 Y_2^2 Z_2^2 + 4 Y_2 Y_3 Z_2 Z_3 + Y_3^2 Z_3^2 + 2 Y_1 Z_1 (7 Y_2 Z_2 + 2 Y_3 Z_3)] + 8 \hat{\delta}^3 Z_1 Z_2 Z_3^3, \\
C_5 &= Y_1^2 Y_2^2 (Y_1 Z_1 + Y_2 Z_2) (Y_1 Z_1 + Y_2 Z_2 + Y_3 Z_3) \\
&\quad - \hat{\delta} Y_1 Y_2 Z_3 [6 Y_1^2 Z_1^2 + 2 Y_2 Z_2 (3 Y_2 Z_2 + 2 Y_3 Z_3) + Y_1 Z_1 (13 Y_2 Z_2 + 4 Y_3 Z_3)] \\
&\quad + 2 \hat{\delta}^2 Z_3^2 [3 Y_1^2 Z_1^2 + Y_2 Z_2 (3 Y_2 Z_2 + Y_3 Z_3) + Y_1 Z_1 (8 Y_2 Z_2 + Y_3 Z_3)] - 2 \hat{\delta}^3 Z_1 Z_2 Z_3^3, \\
C_6 &= Y_1 Y_2 Z_3 (Y_1 Z_1 + Y_2 Z_2 + Y_3 Z_3)^2 \\
&\quad - \hat{\delta} Z_3^2 [3 Y_1^2 Z_1^2 + 3 Y_2^2 Z_2^2 + 4 Y_2 Y_3 Z_2 Z_3 + Y_3^2 Z_3^2 + 4 Y_1 Z_1 (2 Y_2 Z_2 + Y_3 Z_3)] \\
&\quad + 4 \hat{\delta}^2 Z_1 Z_2 Z_3^3.
\end{aligned}$$

- ✿ Can we simplify the results?
- ✿ Can we get a general formula valid for any $s_1-s_2-s_3$?

GENERATING FUNCTION OF ALL MASSIVE & MASSLESS CUBIC INTERACTIONS

- RECALL THE MASSLESS FLAT RESULTS

$$C = \mathcal{K}(y_1, y_2, y_3, g)$$

- IS IT POSSIBLE TO WRITE THE (A)DS INTERACTIONS IN THE SAME WAY?

$$C = \mathcal{K}(?, ?, ?, ?, ?)$$

■ [PROBLEM]

**COUPLINGS ARE NOT HOMOGENEOUS
IN THE NUMBER OF DERIVATIVES**

■ DIFFERENT DERIVATIVE PARTS ARE SIZED BY POWERS OF $\hat{\delta}$ (OR $1/L$)

EX: MINIMUM-DERIVATIVE 4-4-4 COUPLING

$$(Y_1 Z_1 + Y_2 Z_2 + Y_3 Z_3)^4 - 12 \hat{\delta} Z_1 Z_2 Z_3 (Y_1 Z_1 + Y_2 Z_2 + Y_3 Z_3)^2 + 12 \hat{\delta}^2 Z_1^2 Z_2^2 Z_3^2$$

$$\delta^{(n)} \left(\sqrt{\epsilon X^2} - L \right) \equiv \delta \left(\sqrt{\epsilon X^2} - L \right) (\epsilon \hat{\delta})^n$$

■ BUILDING BLOCKS WITH

TOTAL DERIVATIVES !

CUBIC INTERACTIONS OF **MASSIVE** & **MASSLESS** HS

- **3 MASSIVE**

$$C = \mathcal{K}(Y_1, Y_2, Y_3, Z_1, Z_2, Z_3)$$

- **2 MASSIVE** ($\mu_1=\mu_2$)
& 1 MASSLESS

$$C = \mathcal{K}(\tilde{Y}_1, \tilde{Y}_2, \tilde{Y}_3, Z_1, \tilde{G})$$

- **2 MASSIVE** ($\mu_1 \neq \mu_2$)
& 1 MASSLESS

$$C = \mathcal{K}(Y_2, Y_3, Z_1, \tilde{H}_2, \tilde{H}_3)$$

- **1 MASSIVE**
& 2 MASSLESS

$$C = \mathcal{K}(Y_3, \tilde{H}_1, \tilde{H}_2, \tilde{H}_3)$$

- **3 MASSLESS**

$$C = \mathcal{K}(\tilde{Y}_1, \tilde{Y}_2, \tilde{Y}_3, \tilde{G})$$

$$\tilde{Y}_i = Y_i + \alpha_i \partial_X \cdot \partial_{U_i}$$

$$\tilde{G} = \sum_i (Y_i + \beta_i \partial_X \cdot \partial_{U_i}) Z_i$$

$$\begin{aligned} \alpha_1 &= \alpha, & \alpha_2 &= -\frac{1}{\alpha+1}, & \alpha_3 &= -\frac{\alpha+1}{\alpha} \\ \beta_1 &= \beta, & \beta_2 &= -\frac{\beta+1}{\alpha+1}, & \beta_3 &= -\frac{\alpha-\beta}{\alpha} \end{aligned}$$

$$\tilde{H}_i = Y_{i+1} (Y_{i-1} - \partial_X \cdot \partial_{U_{i-1}}) - \frac{1}{2} \partial_X \cdot (\partial_{X_i} - \partial_{X_{i+1}} - \partial_{X_{i-1}}) Z_i$$

EX: MINIMUM-DERIVATIVE 4-4-4 COUPLING

$$(Y_1 Z_1 + Y_2 Z_2 + Y_3 Z_3)^4 - 12 \hat{\delta} Z_1 Z_2 Z_3 (Y_1 Z_1 + Y_2 Z_2 + Y_3 Z_3)^2 + 12 \hat{\delta}^2 Z_1^2 Z_2^2 Z_3^2 = \tilde{G}^4$$

CUBIC INTERACTIONS OF **MASSIVE** & **MASSLESS** HS

- **3 MASSIVE**

$$C = \mathcal{K}(Y_1, Y_2, Y_3, Z_1, Z_2, Z_3)$$

- **2 MASSIVE** ($\mu_1=\mu_2$)
& 1 MASSLESS

$$C = \mathcal{K}(\tilde{Y}_1, \tilde{Y}_2, \tilde{Y}_3, Z_1, \tilde{G})$$

- **2 MASSIVE** ($\mu_1 \neq \mu_2$)
& 1 MASSLESS

$$C = \mathcal{K}(Y_2, Y_3, Z_1, \tilde{H}_2, \tilde{H}_3)$$

- **1 MASSIVE**
& 2 MASSLESS

$$C = \mathcal{K}(Y_3, \tilde{H}_1, \tilde{H}_2, \tilde{H}_3)$$

- **3 MASSLESS**

$$C = \mathcal{K}(\tilde{Y}_1, \tilde{Y}_2, \tilde{Y}_3, \tilde{G})$$

$$\tilde{Y}_i = Y_i + \alpha_i \partial_X \cdot \partial_{U_i}$$

$$\tilde{G} = \sum_i (Y_i + \beta_i \partial_X \cdot \partial_{U_i}) Z_i$$

$$\begin{aligned} \alpha_1 &= \alpha, & \alpha_2 &= -\frac{1}{\alpha+1}, & \alpha_3 &= -\frac{\alpha+1}{\alpha} \\ \beta_1 &= \beta, & \beta_2 &= -\frac{\beta+1}{\alpha+1}, & \beta_3 &= -\frac{\alpha-\beta}{\alpha} \end{aligned}$$

$$\tilde{H}_i = Y_{i+1} (Y_{i-1} - \partial_X \cdot \partial_{U_{i-1}}) - \frac{1}{2} \partial_X \cdot (\partial_{X_i} - \partial_{X_{i+1}} - \partial_{X_{i-1}}) Z_i$$

FLAT LIMIT → METSAEV'S RESULTS

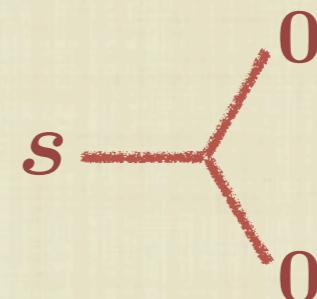
- “HIGHEST-DERIVATIVE” TYPE INTERACTIONS OF **PARTIALLY-MASSLESS HS**

$$C = \mathcal{K}(\tilde{Y}_1, \tilde{Y}_2, \tilde{Y}_3)$$

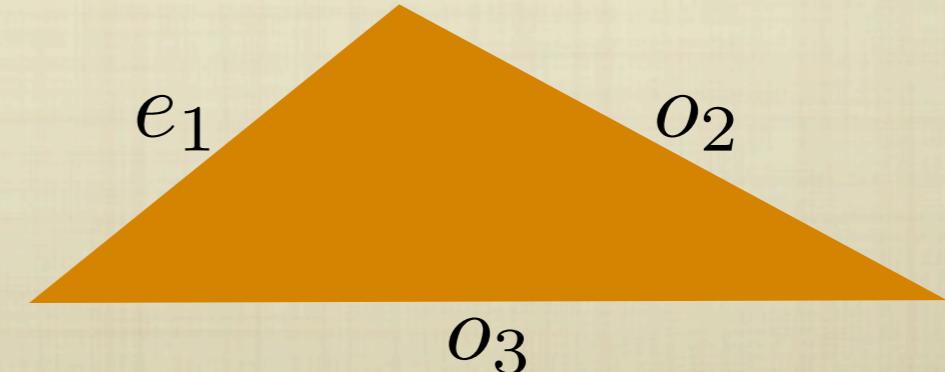
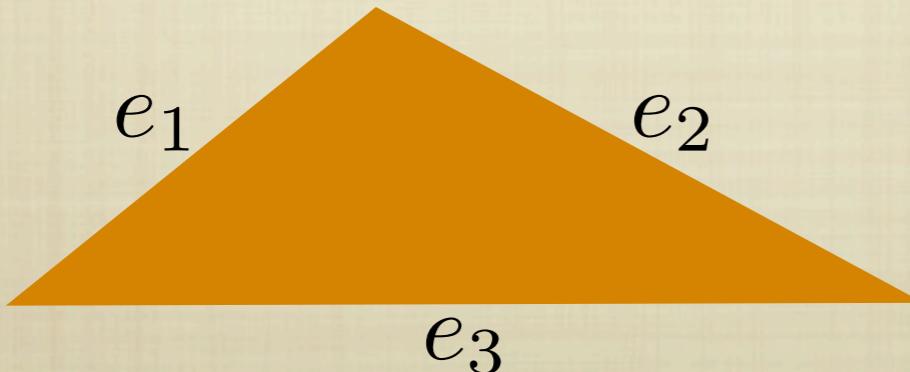
- THIS TYPE OF INTERACTIONS EXISTS IF AND ONLY IF

$$\mu_1 - |\mu_2 - \mu_3| \in 2\mathbb{N}$$

- RESTRICTION ON COUPLINGS OF **PM HS** TO TWO SCALARS



- THREE **PM HS** : TRIANGULAR INEQUALITY



SUMMARY & OUTLOOK

- PDE for consistent interactions in (A)dS
 - Massive & Massless : *generating function*
 - Partially-massless : for given spins
- [Stückelberg formulation & Massless limit]
- Fermions and Mixed-symmetry
- Deformations of HS Gauge Algebra
- AdS/CFT computations