

Motivation (i)

-) The infrared structure of gravity in asymptotically flat (AF) spacetimes is remarkably rich.
-) Gravity in AF spacetimes gives rise to an infinite-dimensional symmetry group (at null infinity):

$$BMS \cong SO(3,1) \ltimes \mathcal{T}$$

\uparrow Lorentz group \uparrow $C^\infty(S^2)$ supertranslations

-) But vacuum (= Minkowski space) only invariant under Poincaré group \implies spontaneous symmetry breaking

vacuum is infinitely degenerate! [Geroch '77, Ashtekar '81]

Motivation (ii)

$$\text{(original) } BMS \approx \underbrace{SO(3,1)}_{\text{global conformal symmetries of } S^2} \ltimes \mathcal{T}$$

\Downarrow

$$\text{extended } BMS = \underbrace{Vir \otimes Vir}_{\text{local conformal symmetries of } S^2} \ltimes \mathcal{T} \quad [\text{Barnich, Troessaert}]$$

more recently: include $Diff(S^2)$ superrotations and/or $Weyl$

[Campiglia, Laddha; Barnich, Lambert; Ciambelli, Leigh, ...]

\Rightarrow vacuum highly degenerate \Rightarrow spontaneous symmetry breaking

\Rightarrow

Goldstone bosons

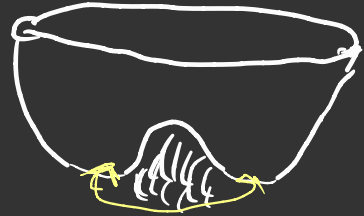
Motivation (iii)

Quick reminder:

spontaneously broken global $U(1)$ symmetry: scalar field
w/ potential

$$\int (\partial\phi)^2 \quad \phi \rightarrow \phi + \alpha$$

$\phi = \text{const} \Rightarrow \text{zero modes}$

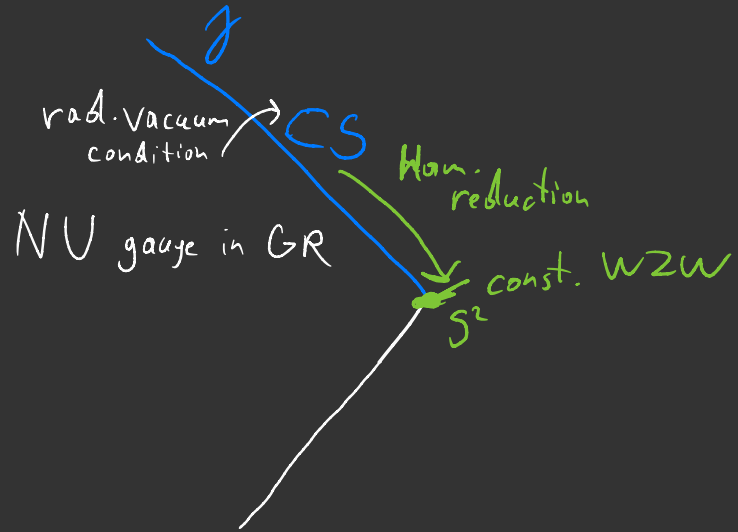


This talk

- 1) find an effective action for superrotation vacuum
- 2) find an action for Goldstone bosons of superrotation

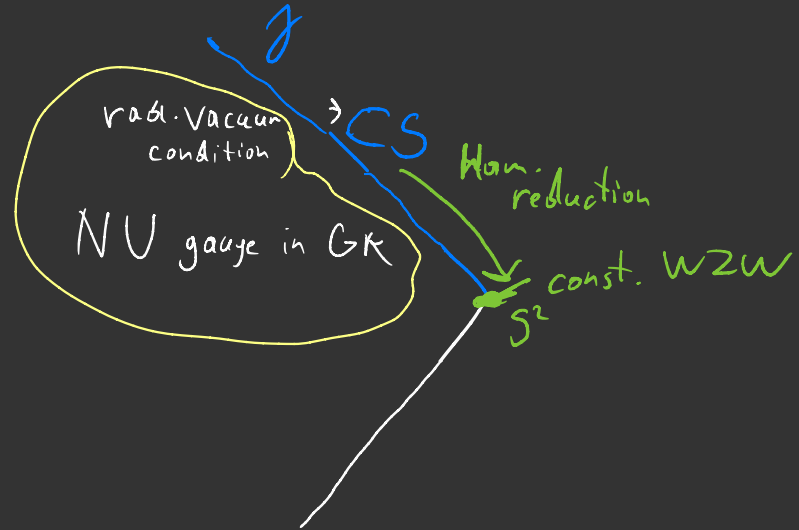
Outline

- 1) Asymptotically flat gravity
- 2) The vacuum conditions & dynamical Carroll geometry on \mathcal{I}^+
- 3) The effective action of superrotation modes
- 4) Comments & Outlook



Outline

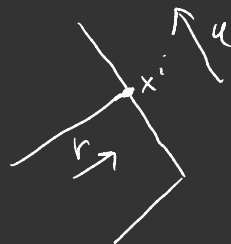
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Asymptotically flat gravity in Newman-Unti gauge (i)

$$ds^2 = g_{uu} du^2 - 2 du dr + g_{ui} du dx^i + g_{ij} dx^i dx^j$$

$$g_{uu} = O(r) \quad g_{ui} = O(1) \quad g_{ij} = r^2 \gamma_{ij} + r C_{ij} + \dots$$



work in first-order formulation

$$ds^2 = \eta_{IJ} E^I \otimes E^J$$

$$E^{\hat{u}} = du$$

$$E^{\hat{r}} = dr + E_{\mu}^{\hat{r}} dx^{\mu}$$

$$E^a = E_{\mu}^a dx^{\mu}$$

$$\eta_{IJ} = \begin{pmatrix} \hat{u} & \hat{r} & \underbrace{a, b} \\ 0 & -1 & \\ -1 & 0 & \\ & & \delta_{ab} \end{pmatrix}_{IJ}$$

AF spacetimes in NU gauge (ii)

tetrad expansion

$$E^{\hat{u}} = du$$

$$E^{\hat{r}} = dr + (r h_{\mu} + h_{\mu}^{(0)} + \dots) dx^{\mu}$$

$$E^a = (r e_{\mu}^a + e_{\mu}^{a(0)} + \dots) dx^{\mu}$$

spin connection expansion

$$\Omega_{\mu}^{ab} = \omega_{\mu}^a \varepsilon^b + O(r^{-2})$$

$$\Omega_{\mu}^{\hat{r}a} = b_{\mu}^a + O(r^{-1})$$

$$\Omega_r^{ab} = 0$$

Asymptotic Symmetries

Local Lorentz (λ, λ^a)

rotation of e^a : $\delta_{\lambda} e^a = \lambda \varepsilon^a_b e^b$

"null rotation": $\delta_{\lambda^a} h = \lambda^a e^a$

$$\delta_{\lambda^a} e^a = \lambda^a \cdot h$$

Supertranslations: do not act on leading order fields

asympt. diffeos ($\eta, \delta^{\mu}_{\nu} e^{\nu}$)

Weyl rescalings: $\delta_{\eta} e^a = \eta e^a$

$$\delta_{\eta} h = d\eta$$

Diff(S^1) superrot: $\delta_{g^i} e^a = \mathcal{L}_g e^a$

Radiative vacua of AF spacetimes

Solve first order Einstein equations order by order in r

•) leading order torsion constraint

$$dh - e^a \wedge b_a = 0 ; d e^a + \omega \varepsilon^a_b \wedge e^b - h \wedge e^a = 0$$

•) scalar curvature $\sim \mathcal{O}(r^{-2})$ [true for most matter fields]

$$d\omega - \varepsilon_{ab} e^a \wedge b^b = 0$$

Additional condition: no radiation @ \mathcal{I}^+

$$d b^a + h \wedge b^a + \omega \varepsilon^a_b \wedge b^b = 0$$

How to interpret these equations?

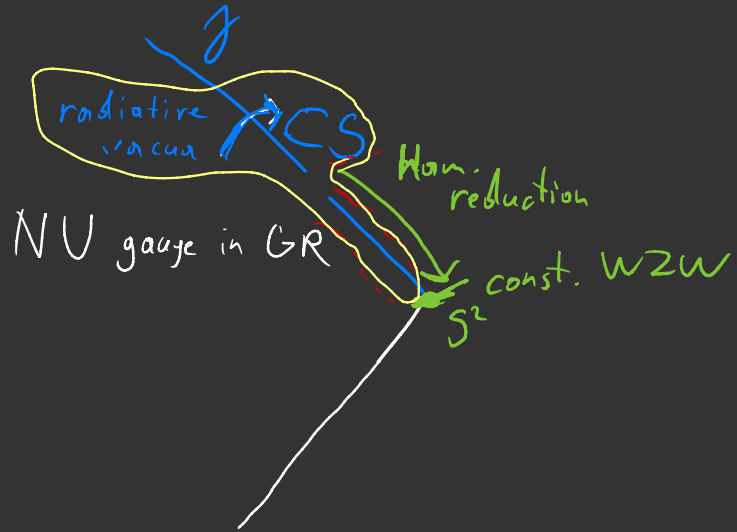
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4) Comments & Conclusion



Homogeneous Spaces of $ISO(3,1)$

$$\text{Mink}_4 \cong \frac{ISO(3,1)}{SO(3,1)} \quad \begin{array}{l} \swarrow \text{choice of codimension 4} \\ \searrow \text{subgroup} \end{array}$$

Minkowski $\Leftrightarrow (iso(3,1), so(3,1))$ "Klein pair"

natural split of generators T of \mathfrak{g} $\begin{cases} \nearrow \text{"internal"} \Leftrightarrow T \in so(3,1) \\ \searrow \text{"external"} \Leftrightarrow T \in \mathfrak{g}/so(3,1) \end{cases}$

More generally: $(iso(3,1), \mathfrak{h}) \Rightarrow (10 - \dim \mathfrak{h})$ -dimensional homogeneous space of Poincaré

Homogeneous Spaces of $ISO(3,1)$

$$\mathcal{J} = \frac{ISO(3,1)}{(ISO(2) \times \mathbb{R}) \times \mathbb{R}^3}$$

$$\text{light-cone}_3 \approx \frac{SO(3,1)}{ISO(2)} \quad \text{"}\mathcal{J} \text{ after forgetting about supertranslations"}$$

$$[J, B_a] = \varepsilon_a{}^b B_b \quad [H, B_a] = B_a \quad [B_a, P_c] = H \delta_{ab} - J \varepsilon_{ab}$$

$$[J, P_a] = \varepsilon_a{}^b P_b \quad [H, P_a] = -P_a$$

\Rightarrow "gauge the light cone"

[cf. 2112.03317] w/ Figueroa O'Farrill, Havel, Prohazka & Stefan's talk @ CW Vienna]

Effective theory of radiative vacua [also Herfray; see also Hartong]

"Gauge the light-cone" [\Leftrightarrow construct Cartan geometry based on Klein pair]

$$A = h H + e^a P_a + b^a B_a + \omega J$$

Symmetry: $\delta_\Lambda A = d\Lambda + [A, \Lambda]$

$$\Lambda = \eta H + \pi^a P_a + \lambda^a B_a + \lambda J$$

act as: Weyl Diff(S⁴) "null rotation" spatial rotation

Effective theory of radiative vacua [also: Herfray]

Define field strength: $F(A) = dA + A \wedge A$

$$F(A) = 0 \iff \text{rad. vacuum conditions}$$

[NB: Until this point we could have "gauged" γ to obtain also supertranslation vacua; cf. Herfray.]

However, for light-cone this follows from an action principle:

$$S_{CS} = \frac{k}{2\pi} \int_{\gamma} \text{tr} \left(A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right)$$

Chern-Simons action

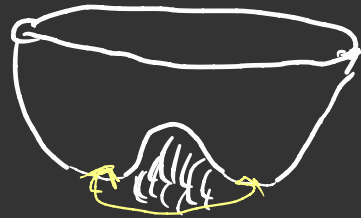
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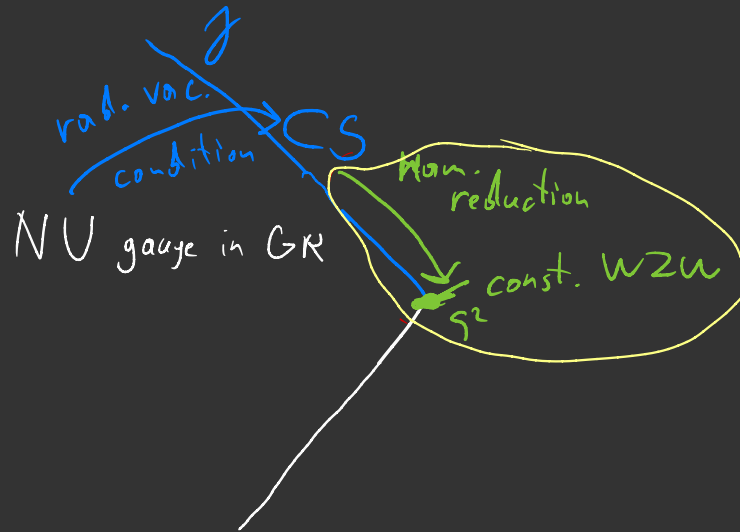


This talk

- 1) find an effective action for superrotation vacua
 \Rightarrow CS action based on $(so(3,1), iso(2))$ ✓
- 2) find an action for Goldstone bosons of superrotation

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- 2) The vacuum conditions of Chern-Simons theory on \mathcal{I}^+
- 3) The effective action of superrotation modes
- 4) Comments & Conclusion

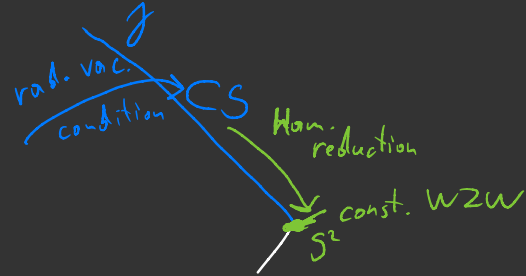


Hamiltonian reduction of CS theory [Coussuert, Henneaux, van Driel; Cotler, Jensen]

•) CS theory is topological \Rightarrow dynamics happen at boundary

•) impose bdy. conditions @ \mathcal{J}^+

•) Hamiltonian reduction of CS:
solve constraint $A_i = g^{-1} \partial_i g$



$$g = e^{\beta_a B^a} e^{\lambda_{WJ} + \lambda_H H} e^{\pi^a P_a}$$

$$S[\pi, \bar{\pi}] = \frac{k}{2\pi} \int_{S^2} dz d\bar{z} \left[\frac{\partial \bar{\partial} \pi \partial \pi}{(\partial \pi)^2} + \frac{\partial \bar{\partial} \bar{\pi} \partial \bar{\pi}}{(\partial \bar{\pi})^2} \right]$$

$\pi = \pi^1 + i \pi^2 \dots$ Goldstone mode of Vir superrotations

2d effective action of superrotations

The effective action of superrotation modes

$$S[\pi, \bar{\pi}] = \frac{k}{2\alpha} \int_{S^2} dz d\bar{z} \left[\frac{\partial_z \partial_{\bar{z}} \pi \partial_z \bar{\pi}}{(\partial_z \pi)^2} + \frac{\partial_{\bar{z}} \partial_z \bar{\pi} \partial_{\bar{z}} \pi}{(\partial_{\bar{z}} \bar{\pi})^2} \right]$$

zero modes: $\pi(z, \bar{z}) = \pi(z)$, $\bar{\pi}(z, \bar{z}) = \bar{\pi}(\bar{z})$

labels $\text{Vir} \otimes \overline{\text{Vir}}$ Vacuum

$SO(3,1)$ gauge symmetry: $\pi \mapsto \frac{a\pi + b}{c\pi + d}$ $ad - bc = 1$

$\pi, \bar{\pi} \in \frac{\text{Vir} \otimes \overline{\text{Vir}}}{\text{PSL}(2, \mathbb{C})} \Rightarrow$ Goldstone modes of broken superrotation

Conclusion

-) Leading order fields in NU gauge obey equations of $SO(3,1)$ CS theory when restricted to radiative vacuum
-) Interpret $SO(3,1)$ -CS theory as "gauging of light-cone"
-) $SO(3,1)$ CS theory for vacuum sector leads to Alekseev-Shatashvili theory on boundary of $\mathcal{J} \implies$ effective 2d description of superrotation vacuum

Outlook

-) "Gauge" Scri \Rightarrow include supertranslations
write down action! [Herfray]
-) physical significance? related to IR divergences? CCFT?
[supertranslation Goldstone bosons related to
IR divergences] [Himwich, Narayanan, Pate, Paul, Strominger; Nguyen JS]
-) homogeneous space: $AdS-Carrdl \approx \frac{SO(3,1)}{SO(2,1)}$
describes "time-like" infinity, similar "space-like" infinity
"gauge" + Lagrangian in 2206.14178 w/ Figueroa-O'Farrill, Havel, Prohazka
 \Rightarrow relate to vacua at space-like infinity

Thank
You!