

Conformal Carroll scalars with boosts

Stefano Baiguera
Ben Gurion University of the Negev

Based on
[SB, G. Oling, W. Sybesma, B. Søggaard, arXiv: 2207.03468]

14 September 2022

Find explicit actions for conformal Carroll theories

- Conformal Carroll isomorphic to BMS
- Few examples of flat space candidate dual field theories
- Include local Carroll boost symmetry

Use a small c expansion to build the Carroll equivalent of

$$S = -\frac{1}{2} \int d^d x \sqrt{-g} \left[g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi + \frac{(d-2)}{4(d-1)} R \phi^2 \right] \quad (1)$$

[Gupta, Suryanarayana, 2020][Bagchi, Banerjee, Dutta, Kolekar, Sharma, 2022][Rivera-Betancour, Vilatte, 2022][Saha, 2022]

Outline

- 1 Carroll geometry
- 2 Conformal Carroll actions
- 3 Conclusions and perspectives

Carroll symmetry

Lorentz boosts

$$t \rightarrow t + \beta x, \quad x \rightarrow x + \beta t \quad (2)$$

Carroll boosts ($c \rightarrow 0$ limit) [Levi-Leblond, 1965][Sen Gupta, 1966]

$$t \rightarrow t + \lambda x, \quad x \rightarrow x \quad (3)$$

Boost Ward identity implies [Henneaux, Salgado-Rebodello, 2021][De Boer, Hartong, Obers, Sybesma, Vandoren, 2021]

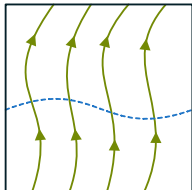
$$\langle \phi(t, x) \phi(0, 0) \rangle = \begin{cases} f(t) \delta(x) & \text{timelike} \\ g(x) & \text{spacelike} \end{cases} \quad (4)$$

Carroll geometry

Geometry defined by a timelike vector field v^μ and a spatial metric $h_{\mu\nu}$
 [Duval, Gibbons, Horvathy, Zhang, 2014][Hartong, 2015][Ciambelli, Marteau, Petkou, Petrououlos, Siampos, 2018][Hansen, Obers, Oling, Sogaard, 2021]

Complemented with τ_μ and $h^{\mu\nu}$ transforming under local boosts

$$v^\mu h_{\mu\nu} = 0, \quad \tau_\mu h^{\mu\nu} = 0, \quad v^\mu \tau_\mu = -1, \quad h^{\mu\rho} h_{\rho\nu} = \delta_\nu^\mu + v^\mu \tau_\nu \quad (5)$$



Extrinsic curvature and acceleration

$$K_{\mu\nu} = -\frac{1}{2} \mathcal{L}_v h_{\mu\nu}, \quad a_\mu = 2v^\rho \partial_{[\mu} \tau_{\rho]} \quad (6)$$

Connection satisfying $\tilde{\nabla}_\mu v^\nu = 0$ and $\tilde{\nabla}_\rho h_{\mu\nu} = 0$
 [Bekaert, Morand, 2015][Hartong, 2015]

$$\tilde{\Gamma}_{\mu\nu}^\rho = -v^\rho \partial_{(\mu} \tau_{\nu)} - v^\rho \tau_{(\mu} \mathcal{L}_v \tau_{\nu)} + \frac{1}{2} h^{\rho\sigma} (\partial_\mu h_{\nu\sigma} + \partial_\nu h_{\sigma\mu} - \partial_\sigma h_{\mu\nu}) - h^{\rho\sigma} \tau_\nu K_{\mu\sigma}$$

Energy-momentum tensor

$$T_\mu^v \equiv -\frac{1}{e} \frac{\delta S}{\delta v^\mu}, \quad T_{\mu\nu}^h \equiv -\frac{2}{e} \frac{\delta S}{\delta h^{\mu\nu}}, \quad e = \det(\tau_\mu, h_{\mu\nu}) \quad (7)$$

Boost-invariant combination [Hartong, Kiritsis, Obers, 2014]

$$T^\mu{}_\nu = -v^\mu T_\nu^v - h^{\mu\rho} T_{\rho\nu}^h \quad (8)$$

Ward identities

- Weyl

$$T^\mu{}_\mu = 0 \quad (9)$$

- Boost (in adapted coordinates (t, x^i))

$$T_0^i = 0 \quad (10)$$

Pre-ultra local parametrization of the geometry

Get Carroll geometry from $c \rightarrow 0$ expansion of Lorentzian geometry (method based on [Hansen, Hartong, Obers, 2020][Hansen, Obers, Oling, SØgaard, 2021])

$$g_{\mu\nu} = -c^2 T_\mu T_\nu + \Pi_{\mu\nu}, \quad g^{\mu\nu} = -\frac{1}{c^2} V^\mu V^\nu + \Pi^{\mu\nu} \quad (11)$$

Carroll geometry appears at leading order in the expansion

$$\begin{aligned} V^\mu &= v^\mu + \mathcal{O}(c^2), & \Pi_{\mu\nu} &= h_{\mu\nu} + \mathcal{O}(c^2) \\ T_\mu &= \tau_\mu + \mathcal{O}(c^2), & \Pi^{\mu\nu} &= h^{\mu\nu} + \mathcal{O}(c^2) \end{aligned} \quad (12)$$

Carroll connection from expansion of Levi-Civita connection compatible with $g_{\mu\nu}$

Conformal Carroll actions

Pre-ultra local parametrization of the action

Relativistic scalar with conformal coupling

$$S = -\frac{1}{2} \int d^d x \sqrt{-g} \left[g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi + \frac{(d-2)}{4(d-1)} R \phi^2 \right] \quad (13)$$

Perform the pre-ultra local parametrization

$$S = -\frac{c}{2} \int d^d x E \left[\left(-\frac{1}{c^2} V^\mu V^\nu + \Pi^{\mu\nu} \right) \partial_\mu \phi \partial_\nu \phi \right. \\ \left. + \frac{(d-2)}{4(d-1)} \left(\frac{1}{c^2} [\mathcal{K}^{\mu\nu} \mathcal{K}_{\mu\nu} + \mathcal{K}^2 - 2V^\mu \partial_\mu \mathcal{K}] \right. \right. \\ \left. \left. + \Pi^{\mu\nu} \tilde{R}_{\mu\nu} - \tilde{\nabla}_\mu (V^\nu \Pi^{\mu\rho} T_{\nu\rho}) + \frac{c^2}{4} \Pi^{\mu\nu} \Pi^{\rho\sigma} T_{\mu\rho} T_{\nu\sigma} \right) \phi^2 \right] \quad (14)$$

No expansion yet!

Timelike action

In the Carroll limit $c \rightarrow 0$, leading order terms give **timelike action**

$$S_t = -\frac{1}{2} \int d^d x e \left[-(v^\mu \partial_\mu \phi)^2 + \frac{(d-2)}{4(d-1)} (K^{\mu\nu} K_{\mu\nu} + K^2 - 2v^\rho \partial_\rho K) \phi^2 \right] \quad (15)$$

Flat space propagator $\sim |t| \delta(\vec{x})$ of timelike form

Symmetries:

- Diffeomorphisms
- Weyl transformations
- Local Carroll boosts

Energy-momentum tensor satisfies

$$T_0^i = 0, \quad T_\mu^\mu = 0 \quad (16)$$

Spacelike action

Alternative $c \rightarrow 0$ limit using Lagrange multipliers gives **spacelike action**

$$S_s = -\frac{1}{2} \int d^d x e \left[h^{\mu\nu} \partial_\mu \phi \partial_\nu \phi + \frac{(d-2)}{4(d-1)} \left(h^{\mu\nu} \tilde{R}_{\mu\nu} - \tilde{\nabla}_\mu a^\mu \right) \phi^2 \right. \\ \left. + \chi \left(v^\mu \partial_\mu \phi - \frac{(d-2)}{2(d-1)} K \phi \right) + \chi^{\mu\nu} K_{\mu\nu} \phi \right] \quad (17)$$

Constraints \Rightarrow Extrinsic curvature must be pure trace

$$K_{\mu\nu} = \frac{h_{\mu\nu}}{d-1} K \quad (18)$$

Timelike action vanishes on the constraint surface

\Rightarrow The two actions cannot coexist at the same time!

(As opposed to [Gupta, Suryanarayana, 2020] [Rivera-Betancour, Vilatte, 2022])

Properties of the spacelike action

Symmetries (**crucial to use constraints!**):

- Diffeomorphisms
- Weyl transformations
- Local Carroll boosts

Energy-momentum tensor satisfies

$$T^i_0 = 0, \quad T^\mu_\mu = 0 \quad (19)$$

Dimensional reduction

Integrate out the constraints

$$S_s = -\frac{1}{2} \int d^{d-1}x \sqrt{h} \left[h^{ij} \partial_i \hat{\phi} \partial_j \hat{\phi} + \frac{(d-3)}{4(d-2)} \hat{R} \hat{\phi}^2 + \frac{1}{4(d-1)(d-2)} A^{-2} \hat{R}_{A^{-2} h_{ij}} \hat{\phi}^2 \right]$$

$$A \equiv \int_v d\tau, \quad \hat{\phi} \equiv A^{1/2} \phi \quad (20)$$

\hat{R} Ricci scalar of Levi-Civita connection compatible with $h_{\mu\nu}$

Euclidean Weyl symmetry manifest!

Conclusions

- Timelike and spacelike conformal Carroll scalar actions
- Well-defined stress tensor satisfying

$$T^i_0 = 0, \quad T^\mu_\mu = 0$$

Future developments

- Build conformal Carroll \Leftrightarrow CCFT dictionary [Donnay, Fiorucci, Herfray, Ruzziconi, 2022][Bagchi, Banerjee, Basu, Dutta, 2022]
- Explicit computation of Weyl anomalies in scalar theories
- General classification of Weyl anomalies

Timelike anomalies (parity-even):

[SB, Oling, Sybesma, SØgaard, in progress][Arav, Chapman, Oz, 2014]

$$\langle T_{\mu}^{\mu} \rangle = \begin{cases} \emptyset & d = 2 \\ \emptyset & d = 3 \\ b_1(-7\text{Tr}(K^4) + \frac{1}{3}K\text{Tr}(K^3) + \text{Tr}(K^2)(\mathcal{L}_v K) + (\mathcal{L}_v K)^2) + b_2(\dots) & d = 4 \\ \dots & \end{cases}$$

Thank you!