

# Quantengravitation und das holographische Universum

Johannes Kepler Universität Linz, Physikkolloquium, März 2017

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## Appetizer, Part I

Physics of the 20<sup>th</sup> century: harmonic oscillator

Simple idea:

Harmonic oscillator: take a physical system and shake it

Amazingly successful:

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- ▶ QFT corrections to Hydrogen atom



Feynman diagrams contributing to Lamb shift

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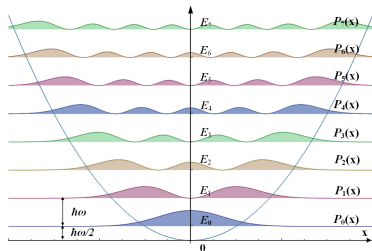
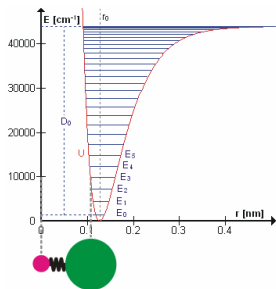
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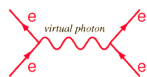
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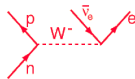
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Electromagnetic

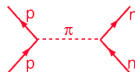


Weak

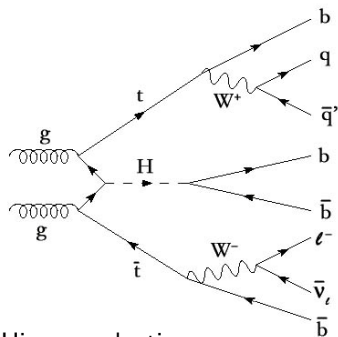


between quarks

Strong Interaction



between nucleons



Higgs production

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- ▶ QFT corrections to Hydrogen atom
- ▶ weakly coupled phonons and electrons in condensed matter
- ▶ Standard Model of particle physics
- ▶ see also the JKU curriculum “Technische Physik”

Lectures in JKU Bachelor curriculum containing harmonic oscillator

- |                                     |                                   |
|-------------------------------------|-----------------------------------|
| ▶ Grundlagen der Physik I-V         | ▶ Theoretische Quantenmechanik I  |
| ▶ Analysis für Physiker(innen) I-II | ▶ Theoretische Thermodynamik      |
| ▶ Mathematische Methoden der Physik | ▶ Theoretische Elektrodynamik I   |
| ▶ Theoretische Mechanik             | ▶ diverse Wahllehrveranstaltungen |

## Appetizer, Part II

Physics of the 21<sup>st</sup> century: black holes? [see colloquium by Strominger at Harvard]

Application of harmonic oscillator limited to perturbative phenomena

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Many physical systems require non-perturbative physics:

- ▶ QCD at low energies
- ▶ High  $T_c$  superconductors
- ▶ Graphene
- ▶ Cold atoms
- ▶ Gravity at high curvature

Generally speaking:

Strongly coupled systems require new techniques



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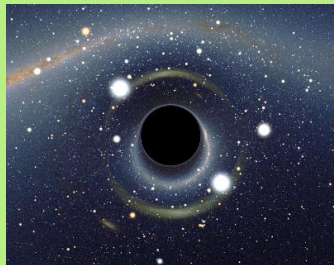
Punch-line of this talk:

Black hole holography can provide such a technique

## Appetizer, Part III

Black holes have apparently paradoxical properties

Black holes: The simplest macroscopic objects in the Universe



Properties determined by:

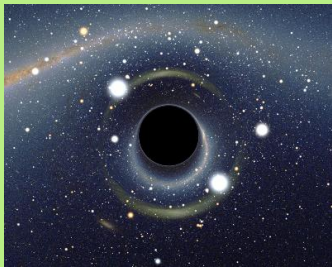
- ▶ Mass  $M$
- ▶ Angular momentum  $J$
- ▶ Charge(s)  $Q$

Black hole  $\sim$  elementary particle!

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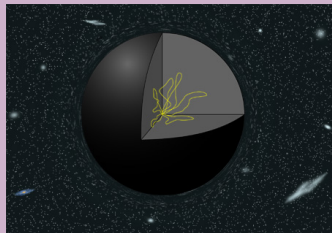


Properties determined by:

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Black holes: The most complicated objects conceivable



Quantum mechanics:

- ▶ Black holes radiate
- ▶ Black holes have entropy
- ▶ Black holes are holographic

Bekenstein–Hawking:

$$S_{\text{BH}} \sim A_{\text{hor}}/4$$

# Outline

Brief history of black holes and observations

Black holes as key to quantum gravity

Black holes and the holographic principle

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Applications of holography

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Simulation of accretion disk around black hole  
(data by K. Thorne et. al. used in movie "Interstellar")

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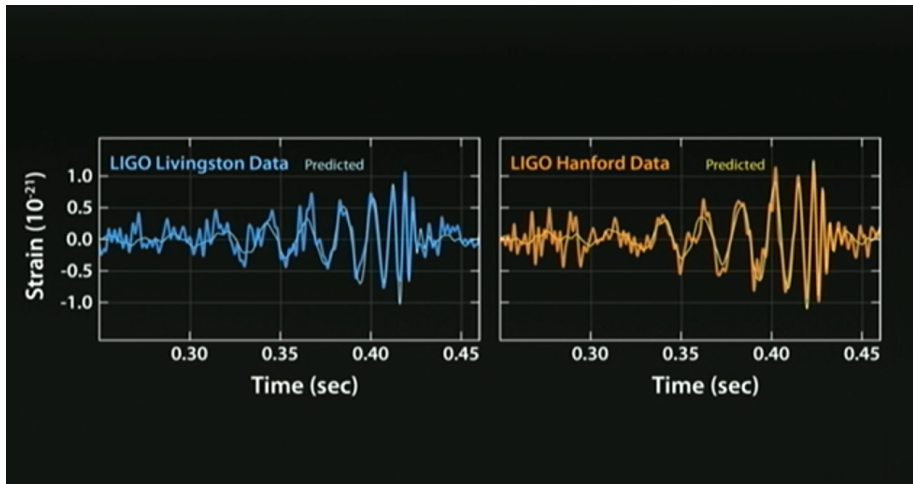
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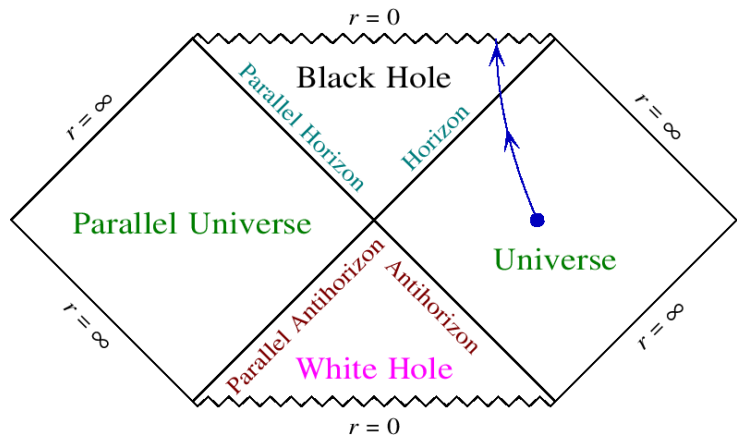
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Gravitational wave signals detected by LIGO in September 2015  
source was a black hole merger ( $36M_{\odot} + 29M_{\odot} = 62M_{\odot} + \text{energy}$ )

## Schwarzschild black hole

Experimental evidence: perihelion shifts, light-bending, GPS, ...



Schwarzschild line-element (horizon at  $r = 2M$ ):

$$ds^2 = -\left(1 - \frac{2M}{r}\right) dt^2 + \frac{dr^2}{1 - \frac{2M}{r}} + r^2 d\theta^2 + r^2 \sin^2\theta d\phi^2$$

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Brief history of black holes and observations

**Black holes as key to quantum gravity**

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## Thermodynamics and black holes — black hole thermodynamics?

### Thermodynamics

Zeroth law:

$T = \text{const.}$  in equilibrium

$T$ : temperature

### Black hole mechanics

Zeroth law:

$\kappa = \text{const.}$  f. stationary black holes

$\kappa$ : surface gravity

## Thermodynamics and black holes — black hole thermodynamics?

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First law:

$dE \sim TdS + \text{work terms}$

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$M$ : mass

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Formal analogy or actual physics?

## Bekenstein's argument

Assume first black holes have no entropy

Simple Gedankenexperiment:

- ▶ Take empty spacetime with a black hole and a cup of tea

Total entropy in Universe:

- ▶  $S_i = S_{\text{tea cup}}$

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- ▶  $S_f < S_i$



## Bekenstein's argument

Assume now black holes have entropy proportional to area of event horizon

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Total entropy in Universe:

- ▶  $S_i = S_{\text{tea cup}} + S_{\text{BH}}$
- ▶  $S = S_{\text{tea cup}} + S_{\text{BH}}$
- ▶  $S_f = 0 + S_{\text{BH}} + \Delta S_{\text{BH}}$
- ▶  $S_f < S_i$   $S_f = S_i$

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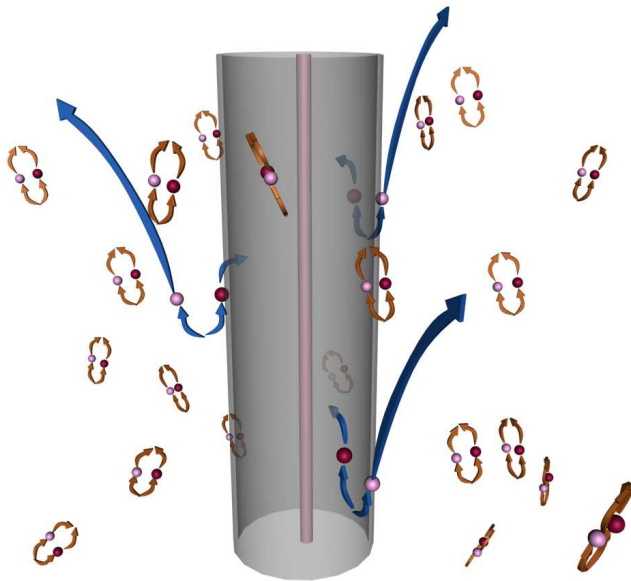
$$S_{\text{BH}} \propto A_{\text{horizon}}$$

Issue above resolved — black hole gets bigger if you throw something in it:

$$S_{\text{total}} = S_{\text{BH}} + S_{\text{tea cup}} = S_{\text{BH}} + \Delta S_{\text{BH}}$$

# Hawking effect confirms Bekenstein's entropy proposal

Black holes evaporate due to quantum effects!



Natural units:

$$T_{\text{H}} = \frac{\kappa}{2\pi}$$

$$S_{\text{BH}} = \frac{A}{4}$$

Schwarzschild  
(SI units):

$$T_{\text{H}} = \frac{\hbar c^3}{8\pi G k_B M}$$

$$S_{\text{BH}} = \frac{c^3 A}{4G\hbar}$$

## Semi-classical puzzles with black holes

### Black holes as the hydrogen atom of quantum gravity

ok, black holes do not violate the second law, but...

- ▶ how can smallest astro-ph black hole have huge entropy

$$S_{\text{BH}} \approx 10^{77}$$

if black holes so simple?

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- ▶ if information paradox resolved like in cond-mat, what are black hole microstates? (and why so many?  $e^{10^{77}}$ )

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- ▶ if black holes thermal states, contradict unitarity of quantum mechanics? (information paradox)
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## Semi-classical puzzles with black holes

### Black holes as the hydrogen atom of quantum gravity

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Understanding quantum behavior of black holes  
crucial milestone on road to quantum gravity!



# Outline

Brief history of black holes and observations

Black holes as key to quantum gravity

**Black holes and the holographic principle**

Evidence for holography

Applications of holography

## Simple motivation of holographic principle

- ▶ entropy in quantum (field) theory

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idea by 't Hooft and Susskind in 1990ies: Holographic Principle

Quantum gravity in  $d+1$  dimensions equivalent to ordinary quantum (field) theory in  $d$  dimensions

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Promising idea — but is it realized in Nature?

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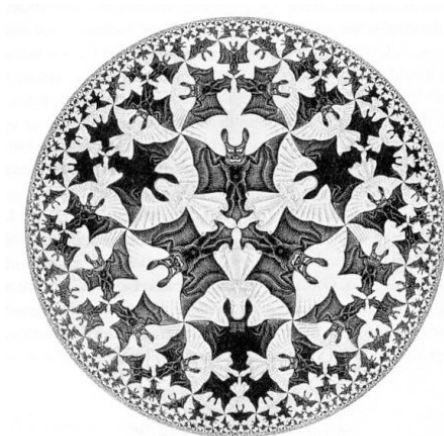
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## AdS/CFT [Maldacena 1997]

Motivating Anti-de Sitter/Conformal Field Theory (AdS/CFT) correspondence

Best studied realization of holography is AdS/CFT correspondence:

- ▶ AdS is a negatively curved spacetime (maximally symmetric)



Open Universe Looking from inside, boundary at infinity  
Limit Circle IV, by M. C. Escher

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Most general line-element compatible with symmetries:

$$ds^2 = (E/L)^2 \eta_{\mu\nu} dx^\mu dx^\nu + (L/E)^2 dE^2$$

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This is precisely the line element of AdS in 1 dimension higher!

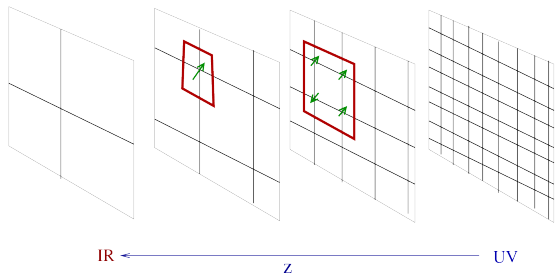
# AdS/CFT

Understanding AdS/CFT as an RG flow [McGreevy 2009]

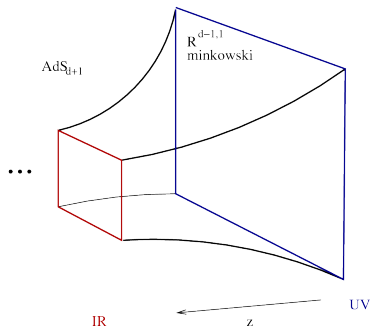
Convenient coordinate trafo:  $z = L^2/E$

$$ds^2 = (L/z)^2 (\eta_{\mu\nu} dx^\mu dx^\nu + dz^2)$$

Field theoretic interpretation: RG-flow!



Left: series of block-spin transformations



Right: cartoon of AdS spacetime

UV in field theory  $\sim$  IR in gravity theory!

## AdS/CFT

AdS<sub>3</sub>/CFT<sub>2</sub> as precursor [Brown, Henneaux 1986]

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- ▶ global physical degrees of freedom, depending on boundary conditions
- ▶ asymptotically AdS<sub>3</sub>: physical Hilbert space falls into representations of two copies of the Virasoro algebra

$$[L_n^\pm, L_m^\pm] = (n - m) L_{n+m}^\pm + \frac{c}{12} (n^3 - n) \delta_{n+m,0}$$

with Brown–Henneaux central charge ( $\ell$  is the AdS radius,  
 $\Lambda = -1/\ell^2$ )

$$c = \frac{3\ell}{2G}$$

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### Conclusion

Any consistent theory of quantum gravity in AdS<sub>3</sub> (compatible with Brown–Henneaux boundary conditions) must be dual to a CFT<sub>2</sub>!



AdS/CFT [Maldacena 1997; Gubser, Klebanov, Polyakov 1997; Witten 1998]

Precise formulation of the conjectured correspondence

Precise statement of AdS/CFT conjecture [Maldacena 1997]:

Type IIB superstring theory on  $\text{AdS}_5 \times S^5$  is equivalent to  $\mathcal{N} = 4$  super-Yang–Mills theory in 3+1 dimensions with gauge group  $U(N)$

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Reformulation of conjecture as equivalence of all correlation functions:

$$\langle \exp \int d^4x \phi_0(x) \mathcal{O}(x) \rangle_{\text{CFT}} = Z_{\text{string}} \left[ \phi(x, z) \Big|_{z=0} = \phi_0(x) \right]$$

l.h.s.: generating function of correlation functions in  $\text{CFT}_4$  for operator  $\mathcal{O}$

r.h.s.: string theory partition function w. condition  $\phi = \phi_0$  at  $\text{AdS}_5$  bdry

## AdS/CFT [see e.g. Aharony, Gubser, Maldacena, Ooguri, Oz 1999]

### Selected checks of the AdS/CFT correspondence

- ▶ perturbative symmetries match (isometries and supersymmetries): supergroup  $SU(2, 2|4)$  (bosonic part:  $SO(4, 2) \times SU(4)$ )
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simple  $AdS_3/CFT_2$  example: all correlation functions of stress energy tensor [[Bagchi, DG, Merbis 2015](#)]

$$\langle T_{\mu_1\nu_1}(z_1)T_{\mu_2\nu_2}(z_2)\dots T_{\mu_n\nu_n}(z_n)\rangle_{CFT_2} = \frac{\delta\Gamma_{AdS_3}}{\delta g^{\mu_1\nu_1}\delta g^{\mu_2\nu_2}\dots\delta g^{\mu_n\nu_n}} \Big|_{EOM}$$

in particular ( $z_{ij} := z_i - z_j$ )

$$\langle T_1 T_2 \rangle = \frac{c}{2z_{12}} \quad c : \text{central charge}$$

$$\langle T_1 T_2 T_3 \rangle = \frac{c}{z_{12}^2 z_{23}^2 z_{13}^2}$$

$$\langle T_1 T_2 \dots T_n T_{n+1} \rangle = \sum_{i=2}^n \left( \frac{2}{z_{1i}^2} + \frac{1}{z_{1i}} \partial_{z_i} \right) \langle T_1 T_2 \dots T_n \rangle$$

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- ▶ anomalies match

e.g. anomaly associated with global  $SU(4)_R$  currents

$$(\mathcal{D}^\mu J_\mu)^a = \frac{N^2 - 1}{384\pi^2} i d^a{}_{bc} \epsilon^{\mu\nu\kappa\lambda} F_{\mu\nu}^b F_{\kappa\lambda}^c$$

gravity computation (valid in large  $N$  limit) yields

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$$\begin{aligned} \Delta = & 4 + 12g^2 - 48g^4 + 336g^6 + 96(-26 + 6\zeta(3) - 15\zeta(5))g^8 \\ & - 96(-158 - 72\zeta(3) + 54\zeta^2(3) + 90\zeta(5) - 315\zeta(7))g^{10} \\ & - 48(160 + 432\zeta^2(3) - 2340\zeta(5) - 72\zeta(3)[-76 + 45\zeta(5)] - 1575\zeta(7) \\ & + 10206\zeta(9))g^{12} + 48(-44480 - 8784\zeta^2(3) + 2592\zeta^3(3) - 4776\zeta(5) \\ & - 20700\zeta^2(5) + 24\zeta(3)[4540 + 357\zeta(5) - 1680\zeta(7)] - 26145\zeta(7) \\ & - 17406\zeta(9) + 152460\zeta(11))g^{14} + \mathcal{O}(g^{16}) \end{aligned}$$

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uses map to spin-chain system and techniques for integrable systems: for instance, magnon-dispersion relation

$$E(p) = \sqrt{1 + \frac{\lambda}{\pi^2} \sin^2(p/2)}$$

or cusp-anomalous dimension  $\pi D_{\text{cusp}} = \sqrt{\lambda} - 3 \ln 2 - \beta(2)/\sqrt{\lambda} + \dots$

$\beta(2)$  is Catalan's constant,  $\beta(2) = \sum_{n=0}^{\infty} (-1)^n / (2n + 1)^2$



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e.g. charge- $k$  type IIB D-instanton action coincides with charge- $k$  super Yang–Mills instanton action [see e.g. Bianchi 2001]

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- ▶ matching of entropy between gravity and field theory side

e.g. Bekenstein–Hawking entropy from Cardy formula for  $CFT_2$

$$S_{\text{BH}} = \frac{A}{4} = 2\pi \sum_{\pm} \sqrt{c\Delta^{\pm}/6} = S_{\text{CFT}}$$

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[see [Beisert et al., 2010 for review](#)]
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- ▶ matching of entropy between gravity and field theory side
- ▶ conceptual checks

e.g. UV/IR connection [[Susskind, Witten 1998](#)]; high  $T$  behavior of entropy in string theory  $S \sim T^3$ , like in  $CFT_4$

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- ▶ holographic entanglement entropy [Ryu, Takayanagi 2006]

remarkable proposal:  $HEE = \text{area of extremal surface}$   
later proved [Lewkowycz, Maldacena 2013]

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- ▶ holographic entanglement entropy [[Ryu, Takayanagi 2006](#)]
- ▶ holographic checks of various inequalities (e.g. holographic entropy bound “ $S \leq S_{\text{BH}}$ ”, quantum null energy condition, quantum focussing conjecture, strong subadditivity, ...)

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... is a bit like 'non-elephant biology'!

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- ▶ Success of AdS/CFT motivates to take holography seriously and test it in more generality
- ▶ Study non-AdS holography to test generality of holographic principle (and also for potential new applications)

## Flat space holography

Ongoing collaboration with [Bagchi et al](#) since 2012 on Flat Space<sub>3</sub>/Galilean CFT<sub>2</sub>

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Brief summary: it may work

- ▶ concrete proposal for holographic correspondence in 3 dimensions

Flat space chiral gravity [Bagchi, DG, Detournay 2012](#)

$$I = \frac{k}{4\pi} \int (\Gamma \wedge d\Gamma + \frac{2}{3} \Gamma \wedge \Gamma \wedge \Gamma)$$

conjectured to be dual to chiral CFT<sub>2</sub> with  $c = 24k$

(for  $k = 1$  conjecturally dual to monster CFT, see [Witten 2007](#))

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Ongoing collaboration with [Bagchi et al](#) since 2012 on Flat Space<sub>3</sub>/Galilean CFT<sub>2</sub>

Brief summary: it may work

- ▶ concrete proposal for holographic correspondence in 3 dimensions
- ▶ perturbative symmetries match

Brown–Henneaux-like pre-cursor [Barnich, Compère 2006](#)

$$[L_n, L_m] = (n - m) L_{n+m} + \frac{c_L}{12} \delta_{n+m,0}$$

$$[L_n, M_m] = (n - m) M_{n+m} + \frac{c_M}{12} \delta_{n+m,0}$$

$$[M_n, M_m] = 0$$

algebra known as BMS<sub>3</sub> or GCA<sub>2</sub> [[Bagchi 2010](#)]

## Flat space holography

Ongoing collaboration with Bagchi et al since 2012 on Flat Space<sub>3</sub>/Galilean CFT<sub>2</sub>

Brief summary: it may work

- ▶ concrete proposal for holographic correspondence in 3 dimensions
- ▶ perturbative symmetries match
- ▶ non-perturbative symmetries match

analogue of S-transformation works on both sides

Barnich 2012; Bagchi, Detournay, Fareghbal, Simon 2012



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e.g. gravitational anomaly  $c - \bar{c} = \frac{3}{\mu G}$  Bagchi, DG, Detournay 2012

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analogous to AdS<sub>3</sub>/CFT<sub>2</sub> calculation [[Bagchi, DG, Merbis 2015](#)]

$$\langle M^1 N^2 \rangle = \frac{c_M}{2s_{12}^4}$$

$$\langle N^1 N^2 \rangle = \frac{c_L - 2c_M \tau_{12}}{2s_{12}^4}$$

$$\langle M^1 N^2 \dots N^n \rangle = \sum_{i=2}^n \left( \frac{2}{s_{1i}^2} + \frac{c_{1i}}{2} \partial_{\varphi_i} \right) \langle M^2 N^3 \dots N^n \rangle$$

$$\langle N^1 N^2 \dots N^n \rangle = \frac{c_L}{c_M} \langle M^1 N^2 \dots N^n \rangle + \sum_{i=1}^n u_i \partial_{\varphi_i} \langle M^1 N^2 \dots N^n \rangle$$

$N, M$ : Galilean/Carrollian conformal analogue of stress-tensor components

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$$S_{\text{BH}} = \frac{A}{4} = 2\pi \sqrt{c_L \Delta_L / 6} + 2\pi \Delta_L \sqrt{c_M / (2\Delta_M)} = S_{\text{GCFT}}$$

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Bagchi, Basu, DG, Riegler 2015; Basu, Riegler 2016

$$S_{\text{HEE}} = \frac{c_L}{6} \ln \frac{\ell_\varphi}{a} + \frac{c_M}{6} \frac{\ell_u}{\ell_\varphi}$$

$c_{L,M}$ : central charges in BMS<sub>3</sub>

$a$ : cut-off

$\ell_{\varphi,u}$ : define size and orientation of entangling region

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- ▶ First tests of flat space holography work
- ▶ Encouraging to pursue flat space holography
- ▶ Numerous further tests possible/desired
- ▶ Numerous conceptual issues (harder than AdS/CFT!)

## Rindler holography/Near horizon holography

Soft Heisenberg hair [Afshar, Detournay, DG, Merbis, Perez, Tempo, Troncoso 2016]

### Main idea

Impose boundary conditions that ensure existence of regular horizon

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- ▶ Surprise 3: algebra can be used to construct all black hole microstates

$$|\text{BTZ} - \text{micro}\rangle \sim \prod_{0 < n_i^\pm < N^\pm} J_{-n_i^\pm}^\pm |0\rangle$$

$J_n^\pm$ : linear combinations of  $X_n, P_n$ ;  $N^\pm$ : mass/angular momentum

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- ▶ Boltzmann's formula yields Bekenstein–Hawking entropy!

$$S = \ln p(N^+) + \ln p(N^-) = 2\pi \sum_{\pm} \sqrt{c\Delta^\pm/6} + \dots = \frac{A}{4} + \dots$$

... denote subleading (log-) corrections

# Outline

Brief history of black holes and observations

Black holes as key to quantum gravity

Black holes and the holographic principle

Evidence for holography

Applications of holography

Why should I care?

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We can expect many new applications in the next decade(s)!

Example: shear viscosity of strongly coupled non-Abelian plasma  
...like the one generated in heavy ion collisions at RHIC and LHC!

Observable of interest is shear viscosity over entropy density,  $\eta/s$

► Perturbative QCD:

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[Policastro, Son, Starinets 2001; Kovtun, Son, Starinets 2005]

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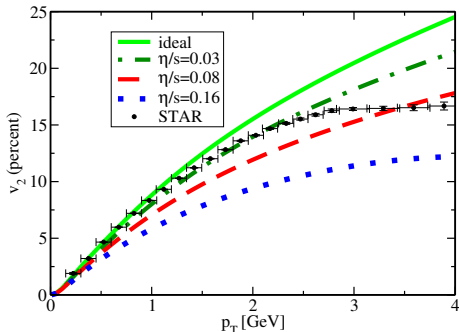
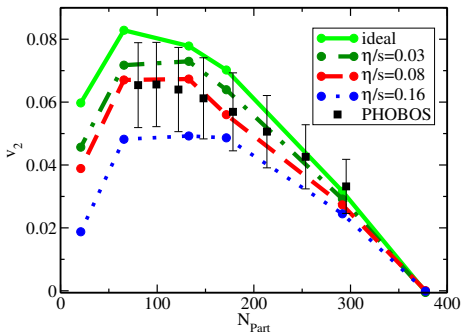
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Simple and sharp prediction from holography for  $\eta/s$   
in strongly coupled non-Abelian plasma

## Experimental results [see e.g. Romatschke, Romatschke 2007]



## Best fit of data [Luzum, Romatschke 2008]

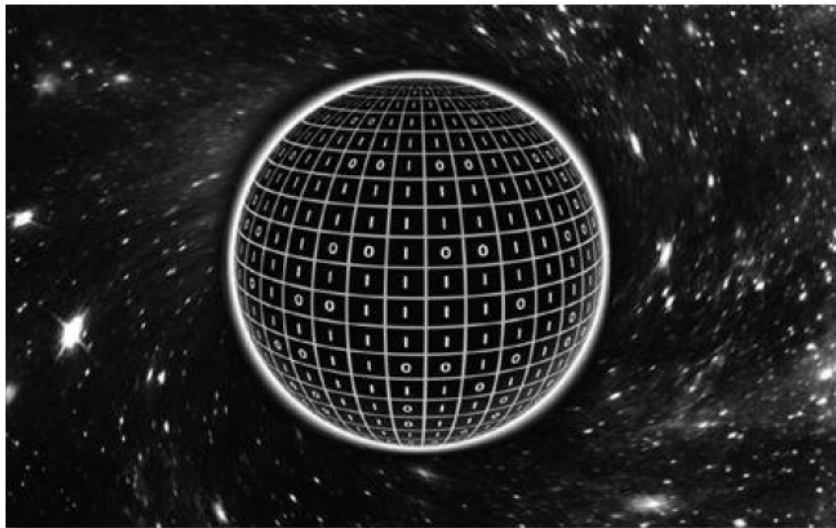
$$\frac{\eta}{s} = 0.10 \pm 0.10(\text{theory}) \pm 0.08(\text{experiment})$$

## Compare with holographic prediction

$$\frac{\eta}{s} = \frac{1}{4\pi} \approx 0.08$$

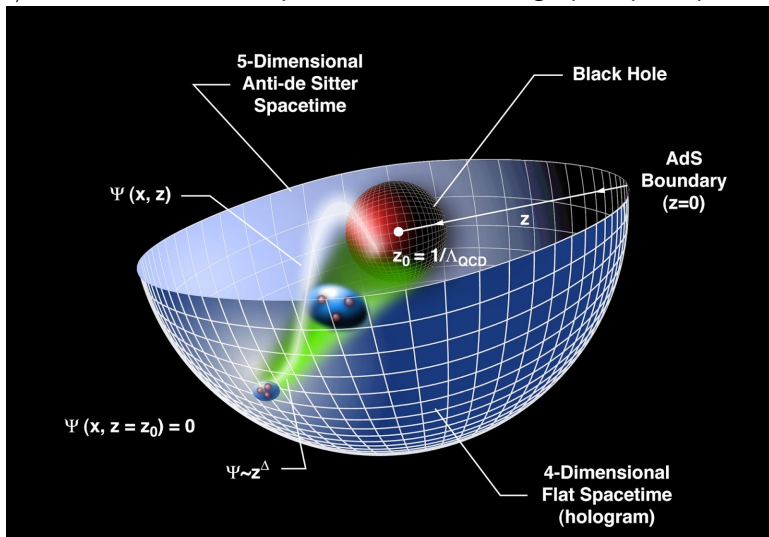
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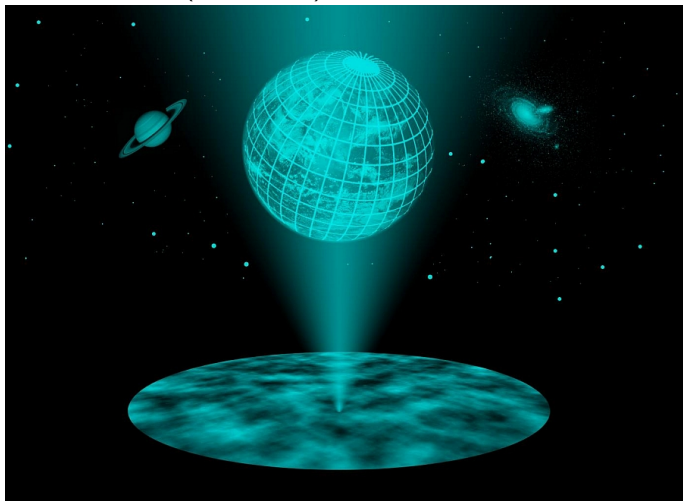
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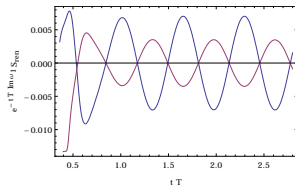
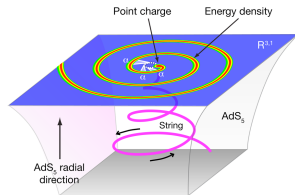
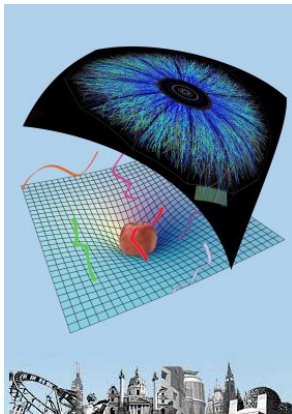
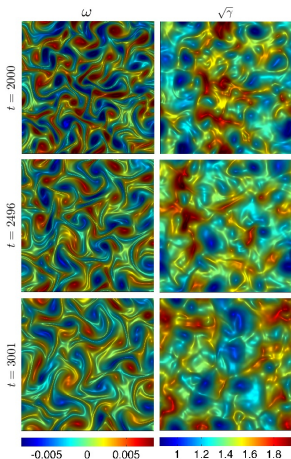
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- ▶ Non-AdS holography (if it exists) of interest for our Universe
- ▶ Numerous applications of AdS/CFT





## My black holes & holography group at TU Wien (postdocs & PhDs)



Wout Merbis



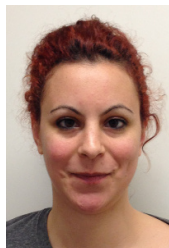
Mirah Gary



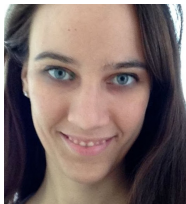
Hernan Gonzalez



Christian Ecker



Maria Irakleidou



Iva Lovrekovic



Stefan Prohazka



Jakob Salzer



Friedrich Schöller



Philipp Stanzer

Thank you for your attention!

