# Quantum Null Energy Condition In two dimensions

Daniel Grumiller

Institute for Theoretical Physics TU Wien

ULB/Leuven/Mons/VUB Seminar, November 2018



Based on work in progress with Ecker, Sheikh-Jabbari, Stanzer and van der Schee

# Outline

Inequalities

QNEC

Holographic QNEC in 4d

Holographic QNEC in 2d

Inequalities are a core part of mathematics

Boring inequalities (type 1): true, but could be sharpened

$$p^2 \ge -p^2 \qquad \forall p \in \mathbb{R}$$

Boring inequalities (type 2): true, but actually an equality

$$p^2 \ge p^2 \qquad \forall p \in \mathbb{R}$$

Fine inequalities: true, cannot be sharpened, not always an equality

$$p^2 \ge 0 \qquad \forall p \in \mathbb{R}$$

- Inequalities are a core part of mathematics
- Many inequalities stem from simple observation that squares of real numbers cannot be negative

$$p^2 \ge 0 \qquad \forall p \in \mathbb{R}$$

- Inequalities are a core part of mathematics
- Many inequalities stem from simple observation that squares of real numbers cannot be negative

$$p^2 \ge 0 \qquad \forall p \in \mathbb{R}$$

Example: given two positive real numbers a, b

algebraic mean  $\geq$  geometric mean

- Inequalities are a core part of mathematics
- Many inequalities stem from simple observation that squares of real numbers cannot be negative

$$p^2 \ge 0 \qquad \forall p \in \mathbb{R}$$

Example: given two positive real numbers a, b

algebraic mean  $\geq$  geometric mean

Proof: take p = a - b and get from inequality above

$$(a-b)^2 = a^2 - 2ab + b^2 \ge 0$$

- Inequalities are a core part of mathematics
- Many inequalities stem from simple observation that squares of real numbers cannot be negative

$$p^2 \ge 0 \qquad \forall p \in \mathbb{R}$$

Example: given two positive real numbers  $\boldsymbol{a}, \boldsymbol{b}$ 

algebraic mean  $\geq$  geometric mean

Proof: take p = a - b and get from inequality above

$$(a-b)^2 = a^2 - 2ab + b^2 \ge 0$$

add on both sides 4ab

$$a^2+2ab+b^2=(a+b)^2\geq 4ab$$

- Inequalities are a core part of mathematics
- Many inequalities stem from simple observation that squares of real numbers cannot be negative

$$p^2 \ge 0 \qquad \forall p \in \mathbb{R}$$

Example: given two positive real numbers a, b

algebraic mean  $\geq$  geometric mean

Proof: take p = a - b and get from inequality above

$$(a-b)^2 = a^2 - 2ab + b^2 \ge 0$$

add on both sides 4ab

$$a^2 + 2ab + b^2 = (a+b)^2 \ge 4ab$$

take square root and then divide by 2

$$\frac{a+b}{2} \ge \sqrt{ab}$$

- Inequalities are a core part of mathematics
- Many inequalities stem from simple observation that squares of real numbers cannot be negative

$$p^2 \ge 0 \qquad \forall p \in \mathbb{R}$$

Many inequalities are of Cauchy–Schwarz type

 $|u||v| \geq |u \cdot v|$ 

here u, v are some vector, || is their length and  $\cdot$  the inner product

- Inequalities are a core part of mathematics
- Many inequalities stem from simple observation that squares of real numbers cannot be negative

$$p^2 \ge 0 \qquad \forall p \in \mathbb{R}$$

Many inequalities are of Cauchy–Schwarz type

 $|u||v| \ge |u \cdot v|$ 

here u, v are some vector, || is their length and  $\cdot$  the inner product



- Inequalities are a core part of mathematics
- Many inequalities stem from simple observation that squares of real numbers cannot be negative

$$p^2 \ge 0 \qquad \forall p \in \mathbb{R}$$

Many inequalities are of Cauchy–Schwarz type

$$|u||v| \geq |u \cdot v|$$

here u, v are some vector, || is their length and  $\cdot$  the inner product Many inequalities from convexity (Jensen's inequality)

- Inequalities are a core part of mathematics
- Many inequalities stem from simple observation that squares of real numbers cannot be negative

 $p^2 \ge 0 \qquad \forall p \in \mathbb{R}$ 

Many inequalities are of Cauchy–Schwarz type

 $|u||v| \geq |u \cdot v|$ 

here u, v are some vector, || is their length and  $\cdot$  the inner product Many inequalities from convexity (Jensen's inequality)



special case of Jensen's inequality: secant always above convex curve between intersection points  $x_1$ ,  $x_2$ 

Interesting physical consequences from mathematical inequalities

• Positivity inequalities: probabilities non-negative,  $P \ge 0$ 

Interesting physical consequences from mathematical inequalities

• Positivity inequalities: probabilities non-negative,  $P \ge 0$ 



Example: unitarity constraints on physical parameters in quark mixing matrix if Standard Model correct then measurements must reproduce unitarity triangle

- $\blacktriangleright$  Positivity inequalities: probabilities non-negative,  $P \geq 0$
- ▶ Cauchy–Schwarz inequalities: Heisenberg uncertainty,  $\Delta x \Delta p \ge \frac{1}{2}$

Interesting physical consequences from mathematical inequalities

- ▶ Positivity inequalities: probabilities non-negative,  $P \ge 0$
- Cauchy–Schwarz inequalities: Heisenberg uncertainty,  $\Delta x \Delta p \geq \frac{1}{2}$



green: localized in coordinate space (x), delocalized in momentum space (p) blue: mildly (de-)localized in coordinate and momentum space orange: delocalized in coordinate space (x), localized in momentum space (p)

- ▶ Positivity inequalities: probabilities non-negative,  $P \ge 0$
- ▶ Cauchy–Schwarz inequalities: Heisenberg uncertainty,  $\Delta x \Delta p \ge \frac{1}{2}$
- Convexity inequalities: second law of thermodynamics,  $\delta S \ge 0$

- $\blacktriangleright$  Positivity inequalities: probabilities non-negative,  $P \geq 0$
- ▶ Cauchy–Schwarz inequalities: Heisenberg uncertainty,  $\Delta x \Delta p \ge \frac{1}{2}$
- Convexity inequalities: second law of thermodynamics,  $\delta S \ge 0$
- In gravitational context: energy inequalities

- $\blacktriangleright$  Positivity inequalities: probabilities non-negative,  $P \geq 0$
- ▶ Cauchy–Schwarz inequalities: Heisenberg uncertainty,  $\Delta x \Delta p \ge \frac{1}{2}$
- Convexity inequalities: second law of thermodynamics,  $\delta S \ge 0$
- In gravitational context: energy inequalities
  - Definition: (local) inequalities on the stress tensor T<sub>μν</sub> e.g. Null Energy Condition (NEC)

$$T_{kk} = T_{\mu\nu} \, k^{\mu} k^{\nu} \ge 0 \qquad \forall k^{\mu} k_{\mu} = 0$$

Interesting physical consequences from mathematical inequalities

- $\blacktriangleright$  Positivity inequalities: probabilities non-negative,  $P \geq 0$
- Cauchy–Schwarz inequalities: Heisenberg uncertainty,  $\Delta x \Delta p \geq \frac{1}{2}$
- Convexity inequalities: second law of thermodynamics,  $\delta S \ge 0$
- In gravitational context: energy inequalities
  - Definition: (local) inequalities on the stress tensor T<sub>μν</sub> e.g. Null Energy Condition (NEC)

$$T_{kk} = T_{\mu\nu} \, k^{\mu} k^{\nu} \ge 0 \qquad \forall k^{\mu} k_{\mu} = 0$$

For instance: Penrose singularity theorem from Raychaudhuri eq.

$$\frac{\mathrm{d}^2 \mathrm{area}}{\mathrm{d}k^2} = -\left(\frac{\mathrm{d}\operatorname{area}}{\mathrm{d}k}\right)^2 - \mathrm{shear}^2 - 8\pi G T_{kk} \leq -8\pi G T_{kk} \stackrel{\mathrm{NEC}}{\leq} 0$$

If  $T_{kk} \ge 0$  (NEC)  $\Rightarrow$  focussing! (negative acceleration of area)

For experts:  $\frac{d \operatorname{area}}{dk} = \theta$  is null expansion



Inequalities

Interesting physical consequences from mathematical inequalities

- $\blacktriangleright$  Positivity inequalities: probabilities non-negative,  $P \geq 0$
- ▶ Cauchy–Schwarz inequalities: Heisenberg uncertainty,  $\Delta x \Delta p \ge \frac{1}{2}$
- Convexity inequalities: second law of thermodynamics,  $\delta S \ge 0$
- In gravitational context: energy inequalities
  - Definition: (local) inequalities on the stress tensor T<sub>μν</sub> e.g. Null Energy Condition (NEC)

$$T_{kk} = T_{\mu\nu} \, k^{\mu} k^{\nu} \ge 0 \qquad \forall k^{\mu} k_{\mu} = 0$$

Physically plausible (positivity of energy fluxes)

- $\blacktriangleright$  Positivity inequalities: probabilities non-negative,  $P \geq 0$
- ▶ Cauchy–Schwarz inequalities: Heisenberg uncertainty,  $\Delta x \Delta p \ge \frac{1}{2}$
- Convexity inequalities: second law of thermodynamics,  $\delta S \ge 0$
- In gravitational context: energy inequalities
  - Definition: (local) inequalities on the stress tensor T<sub>μν</sub> e.g. Null Energy Condition (NEC)

$$T_{kk} = T_{\mu\nu} \, k^{\mu} k^{\nu} \ge 0 \qquad \forall k^{\mu} k_{\mu} = 0$$

- Physically plausible (positivity of energy fluxes)
- ▶ Mathematically useful (singularity theorem, area theorem [2<sup>nd</sup> law])

Interesting physical consequences from mathematical inequalities

- $\blacktriangleright$  Positivity inequalities: probabilities non-negative,  $P\geq 0$
- Cauchy–Schwarz inequalities: Heisenberg uncertainty,  $\Delta x \Delta p \geq \frac{1}{2}$
- Convexity inequalities: second law of thermodynamics,  $\delta S \ge 0$
- In gravitational context: energy inequalities
  - Definition: (local) inequalities on the stress tensor T<sub>μν</sub> e.g. Null Energy Condition (NEC)

$$T_{kk} = T_{\mu\nu} \, k^{\mu} k^{\nu} \ge 0 \qquad \forall k^{\mu} k_{\mu} = 0$$

- Physically plausible (positivity of energy fluxes)
- Mathematically useful (singularity theorem, area theorem [2<sup>nd</sup> law])

However: all classical energy inequalities violated by quantum effects!

NEC violated by Casimir energy, accelerated mirrors, Hawking radiation, ...

Interesting physical consequences from mathematical inequalities

- $\blacktriangleright$  Positivity inequalities: probabilities non-negative,  $P \geq 0$
- ▶ Cauchy–Schwarz inequalities: Heisenberg uncertainty,  $\Delta x \Delta p \ge \frac{1}{2}$
- Convexity inequalities: second law of thermodynamics,  $\delta S \ge 0$
- In gravitational context: energy inequalities
  - Definition: (local) inequalities on the stress tensor T<sub>μν</sub> e.g. Null Energy Condition (NEC)

$$T_{kk} = T_{\mu\nu} \, k^{\mu} k^{\nu} \ge 0 \qquad \forall k^{\mu} k_{\mu} = 0$$

- Physically plausible (positivity of energy fluxes)
- Mathematically useful (singularity theorem, area theorem [2<sup>nd</sup> law])

However: all classical energy inequalities violated by quantum effects!

Are there quantum energy conditions? [How is  $2^{nd}$  law saved?]

▶ Definition: quantum energy condition = convexity condition for  $\langle T_{\mu\nu} \rangle$  valid for any state and any (reasonable) quantum field theory (QFT)

- ▶ Definition: quantum energy condition = convexity condition for  $\langle T_{\mu\nu} \rangle$  valid for any state and any (reasonable) quantum field theory (QFT)
- Example: Averaged Null Energy Condition (ANEC)

$$\int \mathrm{d}x^{\lambda} k_{\lambda} \left\langle T_{\mu\nu} k^{\mu} k^{\nu} \right\rangle \ge 0$$

- ▶ Definition: quantum energy condition = convexity condition for  $\langle T_{\mu\nu} \rangle$  valid for any state and any (reasonable) quantum field theory (QFT)
- Example: Averaged Null Energy Condition (ANEC)

$$\int \mathrm{d}x^{\lambda} k_{\lambda} \left\langle T_{\mu\nu} k^{\mu} k^{\nu} \right\rangle \ge 0$$

valid  $\forall k^{\mu}(\text{with }k^{\mu}k_{\mu}=0)$  and  $\forall$  states  $|\rangle$  in any (reasonable) QFT

ANEC proved under rather generic assumptions

Faulkner, Leigh, Parrikar and Wang 1605.08072 Hartman, Kundu and Tajdini 1610.05308

- ▶ Definition: quantum energy condition = convexity condition for  $\langle T_{\mu\nu} \rangle$  valid for any state and any (reasonable) quantum field theory (QFT)
- Example: Averaged Null Energy Condition (ANEC)

$$\int \mathrm{d}x^{\lambda} k_{\lambda} \left\langle T_{\mu\nu} k^{\mu} k^{\nu} \right\rangle \ge 0$$

- ANEC proved under rather generic assumptions
- ANEC sufficient for focussing properties used in singularity theorems

- ▶ Definition: quantum energy condition = convexity condition for  $\langle T_{\mu\nu} \rangle$  valid for any state and any (reasonable) quantum field theory (QFT)
- Example: Averaged Null Energy Condition (ANEC)

$$\int \mathrm{d}x^{\lambda} k_{\lambda} \left\langle T_{\mu\nu} k^{\mu} k^{\nu} \right\rangle \ge 0$$

- ANEC proved under rather generic assumptions
- ANEC sufficient for focussing properties used in singularity theorems
- ANEC compatible with quantum interest conjecture

- ▶ Definition: quantum energy condition = convexity condition for  $\langle T_{\mu\nu} \rangle$  valid for any state and any (reasonable) quantum field theory (QFT)
- Example: Averaged Null Energy Condition (ANEC)

$$\int \mathrm{d}x^{\lambda} k_{\lambda} \left\langle T_{\mu\nu} k^{\mu} k^{\nu} \right\rangle \ge 0$$

- ANEC proved under rather generic assumptions
- ANEC sufficient for focussing properties used in singularity theorems
- ANEC compatible with quantum interest conjecture
- However: ANEC is non-local  $(\int dx^+)$

- ▶ Definition: quantum energy condition = convexity condition for  $\langle T_{\mu\nu} \rangle$  valid for any state and any (reasonable) quantum field theory (QFT)
- Example: Averaged Null Energy Condition (ANEC)

$$\int \mathrm{d}x^{\lambda} k_{\lambda} \left\langle T_{\mu\nu} k^{\mu} k^{\nu} \right\rangle \ge 0$$

valid  $\forall k^{\mu}(\text{with }k^{\mu}k_{\mu}=0)$  and  $\forall$  states  $|\rangle$  in any (reasonable) QFT

- ANEC proved under rather generic assumptions
- ANEC sufficient for focussing properties used in singularity theorems
- ANEC compatible with quantum interest conjecture
- However: ANEC is non-local  $(\int dx^+)$

Is there a local quantum energy condition?

# Outline

Inequalities

# QNEC

Holographic QNEC in 4d

Holographic QNEC in 2d

Proposed by Bousso, Fisher, Leichenauer and Wall in 1506.02669

QNEC (in 
$$D>2$$
) is the following inequality 
$$\langle T_{kk}\rangle\geq \frac{\hbar}{2\pi\sqrt{\gamma}}\,S''$$

Physical motivation from focussing properties and second law: Replace area by area + 4G (entanglement entropy) Modified Raychaudhuri eq., schematically:

$$\frac{\mathrm{d}^2 \mathrm{area}}{\mathrm{d}k^2} + 4G\,S'' = -8\pi G\,T_{kk} + 4G\,S'' \stackrel{\mathrm{QNEC}}{\leq} 0$$

requires for focussing property (=2<sup>nd</sup> law) QNEC

fineprint: above we set expansion to zero,  $\frac{d \operatorname{area}}{dk} = 0$ , and shear to zero; we also set the area to unity,  $\sqrt{\gamma} = 1$  thus, QNEC is implied from quantum focussing for special congruences

Proposed by Bousso, Fisher, Leichenauer and Wall in 1506.02669

QNEC (in D>2) is the following inequality  $\langle T_{kk}\rangle \geq \frac{\hbar}{2\pi\sqrt{\gamma}}\,S''$ 

Obvious observations:

- ▶ if r.h.s. vanishes: semi-classical version of NEC
- if r.h.s. negative: weaker condition than NEC (NEC can be violated while QNEC holds)
- if r.h.s. positive: stronger condition than NEC (if QNEC holds also NEC holds)

Proposed by Bousso, Fisher, Leichenauer and Wall in 1506.02669





•  $T_{kk} = T_{\mu\nu}k^{\mu}k^{\nu}$  with  $k_{\mu}k^{\mu} = 0$  and  $\langle \rangle$  denotes expectation value

Proposed by Bousso, Fisher, Leichenauer and Wall in 1506.02669





•  $T_{kk} = T_{\mu\nu}k^{\mu}k^{\nu}$  with  $k_{\mu}k^{\mu} = 0$  and  $\langle \rangle$  denotes expectation value • S'': 2<sup>nd</sup> variation of EE for entangling surface deformations along  $k_{\mu}$
### Quantum null energy condition (QNEC)

Proposed by Bousso, Fisher, Leichenauer and Wall in 1506.02669





•  $T_{kk} = T_{\mu\nu}k^{\mu}k^{\nu}$  with  $k_{\mu}k^{\mu} = 0$  and  $\langle \rangle$  denotes expectation value • S'': 2<sup>nd</sup> variation of EE for entangling surface deformations along  $k_{\mu}$ •  $\sqrt{\gamma}$ : induced volume form of entangling region (black boundary curve)

# Proofs (D > 2)

- For free QFTs: Bousso, Fisher, Koeller, Leichenauer and Wall, 1509.02542
- For holographic CFTs: Koeller and Leichenauer, 1512.06109
- For general CFTs: Balakrishnan, Faulkner, Khandker and Wang, 1706.09432

Ongoing work with Ecker, Sheikh-Jabbari, Stanzer and van der Schee

QNEC (in CFT\_2) is the following inequality  $\langle T_{kk}\rangle\geq S''+\frac{6}{c}\,S'^2$  c>0 is the central charge of the CFT\_2

Ongoing work with Ecker, Sheikh-Jabbari, Stanzer and van der Schee

QNEC (in CFT\_2) is the following inequality 
$$\langle T_{kk}\rangle\geq S''+\frac{6}{c}\,S'^2$$
  $c>0$  is the central charge of the CFT\_2

▶ S like anomalous operator with conformal weights (0,0)⇒ construct vertex operator  $V = \exp\left[\frac{6}{c}S\right]$ 

Ongoing work with Ecker, Sheikh-Jabbari, Stanzer and van der Schee

QNEC (in CFT\_2) is the following inequality 
$$\langle T_{kk}\rangle\geq S''+\frac{6}{c}\,S'^2$$
  $c>0$  is the central charge of the CFT\_2

- ► S like anomalous operator with conformal weights (0,0)⇒ construct vertex operator  $V = \exp\left[\frac{6}{c}S\right]$
- QNEC saturation equivalent to vertex operator solving Hill's equation

$$V'' - \mathcal{L}V = 0 \qquad \qquad \mathcal{L} \sim \langle T_{kk} \rangle$$

Ongoing work with Ecker, Sheikh-Jabbari, Stanzer and van der Schee

QNEC (in CFT\_2) is the following inequality 
$$\langle T_{kk}\rangle\geq S''+\frac{6}{c}\,S'^2$$
  $c>0$  is the central charge of the CFT\_2

- ► S like anomalous operator with conformal weights (0,0)⇒ construct vertex operator  $V = \exp\left[\frac{6}{c}S\right]$
- QNEC saturation equivalent to vertex operator solving Hill's equation

$$V'' - \mathcal{L}V = 0 \qquad \qquad \mathcal{L} \sim \langle T_{kk} \rangle$$

QNEC saturated for vacuum, thermal states and their descendants

Ongoing work with Ecker, Sheikh-Jabbari, Stanzer and van der Schee

QNEC (in CFT\_2) is the following inequality 
$$\langle T_{kk}\rangle\geq S''+\frac{6}{c}\,S'^2$$
  $c>0$  is the central charge of the CFT\_2

- ► S like anomalous operator with conformal weights (0,0)⇒ construct vertex operator  $V = \exp\left[\frac{6}{c}S\right]$
- QNEC saturation equivalent to vertex operator solving Hill's equation

$$V'' - \mathcal{L}V = 0 \qquad \qquad \mathcal{L} \sim \langle T_{kk} \rangle$$

QNEC saturated for vacuum, thermal states and their descendants
QNEC not saturated in hol. CFT<sub>2</sub> with positive bulk energy fluxes

Ongoing work with Ecker, Sheikh-Jabbari, Stanzer and van der Schee

QNEC (in CFT\_2) is the following inequality 
$$\langle T_{kk}\rangle\geq S''+\frac{6}{c}\,S'^2$$
  $c>0$  is the central charge of the CFT\_2

- ► S like anomalous operator with conformal weights (0,0)⇒ construct vertex operator  $V = \exp\left[\frac{6}{c}S\right]$
- QNEC saturation equivalent to vertex operator solving Hill's equation

$$V'' - \mathcal{L}V = 0 \qquad \qquad \mathcal{L} \sim \langle T_{kk} \rangle$$

- QNEC saturated for vacuum, thermal states and their descendants
- ▶ QNEC not saturated in hol. CFT<sub>2</sub> with positive bulk energy fluxes
- QNEC can be violated in hol.  $CFT_2$  with negative bulk energy fluxes

Calculating QNEC holographically

calculating CFT observable holographically = some gravity calculation

AdS/CFT:

Maldacena hep-th/9711200 Gubser, Klebanov and Polyakov hep-th/9802109 Witten hep-th/9802150

holographic stress tensor:

Henningson and Skenderis hep-th/9806087 Balasubramanian and Kraus hep-th/9902121 Emparan, Johnson and Myers hep-th/9903238 de Haro, Solodukhin and Skenderis hep-th/0002230

holographic entanglement entropy (HEE): Ryu and Takayanagi hep-th/0603001 Hubeny, Rangamani and Takayanagi 0705.0016

### Calculating QNEC holographically

calculating CFT observable holographically = some gravity calculation

• need holographic computation of  $\langle T_{kk} \rangle$ 

well-known AdS/CFT prescription: extract boundary stress tensor from bulk metric expanded near AdS boundary

Example:  $AdS_3/CFT_2$ 

$$ds^{2} = \frac{\ell^{2}}{z^{2}} \left( dz^{2} + 2 dx^{+} dx^{-} \right) + \langle T_{++} \rangle \left( dx^{+} \right)^{2} + \langle T_{--} \rangle \left( dx^{-} \right)^{2} + \mathcal{O}(z^{2})$$

AdS<sub>3</sub> boundary:  $z \to 0$  $\mathcal{O}(1)$  terms in metric: flux components of stress tensor  $\langle T_{\pm\pm} \rangle$ (trace vanishes,  $\langle T_{+-} \rangle = 0$ )  $\ell$ : so-called AdS-radius (cosmological constant  $\Lambda = -1/\ell^2$ )

### Calculating QNEC holographically

calculating CFT observable holographically = some gravity calculation

- need holographic computation of  $\langle T_{kk} \rangle$
- need holographic computation of (deformations of) EE

 $\mathsf{HEE} = \mathsf{area} \ \mathsf{of} \ \mathsf{extremal} \ \mathsf{surface}$ 

### simple to calculate!



also: simple proof of strong subadditivity inequalities



# Outline

Inequalities

QNEC

### Holographic QNEC in 4d

Holographic QNEC in 2d

see work with Ecker, Stanzer and van der Schee 1710.09837

# thermal states in $\mathsf{CFT}_4 = \mathsf{black}$ holes in $\mathsf{AdS}_5$

paper-and-pencil calculation starts with Schwarzschild black brane

$$ds^{2} = \frac{1}{z^{2}} \left( -f(z) dt^{2} + \frac{dz^{2}}{f(z)} + dy^{2} + dx_{1}^{2} + dx_{2}^{2} \right)$$

with  $f(z) = 1 - \pi^4 T^4 z^4$ 

t = const.



Daniel Grumiller — Quantum Null Energy Condition

see work with Ecker, Stanzer and van der Schee 1710.09837

### thermal states in $CFT_4 = black$ holes in $AdS_5$

paper-and-pencil calculation starts with Schwarzschild black brane

$$ds^{2} = \frac{1}{z^{2}} \left( -f(z) dt^{2} + \frac{dz^{2}}{f(z)} + dy^{2} + dx_{1}^{2} + dx_{2}^{2} \right)$$

with  $f(z)=1-\pi^4T^4z^4$ 

• determine area of minimal surfaces for small temperature,  $T\ell \ll 1$ , and extract HEE ( $\ell$  = width of strip)

$$\frac{1}{2\pi} S'' \approx -\frac{0.065}{\ell^4} + 0.019 \,\pi^4 T^4 - 0.083 \,\ell^4 \pi^8 T^8 + \mathcal{O}\big(\ell^8 T^{12}\big)$$

see work with Ecker, Stanzer and van der Schee 1710.09837

# thermal states in $CFT_4 = black$ holes in $AdS_5$

paper-and-pencil calculation starts with Schwarzschild black brane

$$ds^{2} = \frac{1}{z^{2}} \left( -f(z) dt^{2} + \frac{dz^{2}}{f(z)} + dy^{2} + dx_{1}^{2} + dx_{2}^{2} \right)$$

with  $f(z)=1-\pi^4T^4z^4$ 

• determine area of minimal surfaces for small temperature,  $T\ell \ll 1$ , and extract HEE ( $\ell$  = width of strip)

$$\frac{1}{2\pi} S'' \approx -\frac{0.065}{\ell^4} + 0.019 \,\pi^4 T^4 - 0.083 \,\ell^4 \pi^8 T^8 + \mathcal{O}\big(\ell^8 T^{12}\big)$$

▶ do same for large temperatures,  $T\ell \gg 1$ 

$$\frac{1}{2\pi} S'' \approx -0.364 \,\pi^4 T^4 \, e^{-\sqrt{6}\ell\pi T} + \mathcal{O}\left(e^{-2\sqrt{6}\ell\pi T}\right)$$

see work with Ecker, Stanzer and van der Schee 1710.09837

# thermal states in $\mathsf{CFT}_4 = \mathsf{black}$ holes in $\mathsf{AdS}_5$

paper-and-pencil calculation starts with Schwarzschild black brane

$$ds^{2} = \frac{1}{z^{2}} \left( -f(z) dt^{2} + \frac{dz^{2}}{f(z)} + dy^{2} + dx_{1}^{2} + dx_{2}^{2} \right)$$

with  $f(z)=1-\pi^4T^4z^4$ 

• determine area of minimal surfaces for small temperature,  $T\ell \ll 1$ , and extract HEE ( $\ell$  = width of strip)

$$\frac{1}{2\pi} S'' \approx -\frac{0.065}{\ell^4} + 0.019 \,\pi^4 T^4 - 0.083 \,\ell^4 \pi^8 T^8 + \mathcal{O}\big(\ell^8 T^{12}\big)$$

 $\blacktriangleright$  do same for large temperatures,  $T\ell\gg 1$ 

$$\frac{1}{2\pi} S'' \approx -0.364 \,\pi^4 T^4 \, e^{-\sqrt{6}\ell\pi T} + \mathcal{O}\left(e^{-2\sqrt{6}\ell\pi T}\right)$$

use numerics for intermediate values of temperature

see work with Ecker, Stanzer and van der Schee 1710.09837

#### thermal states in $CFT_4$ = black holes in $AdS_5$



notational alert: L in the plot corresponds to width  $\ell$ 

Daniel Grumiller — Quantum Null Energy Condition

- paper-and-pencil calculations with Romatschke 0803.3226
  - $\delta$ -like shocks
  - particle production in forward lightcone of shocks
  - ▶ shortly after collision anisotropic pressure:  $P_L/E = -3$ ,  $P_T/E = +2$  confirmed numerically for thin shocks by Casalderrey-Solana, Heller, Mateos and van der Schee 1305.4919
  - Close to shockwaves negative energy fluxes ⇒ NEC violation! confirmed numerically and interpreted as absence of local rest frame by Arnold, Romatschke and van der Schee 1408.2518

plasma formation in  $CFT_4$  = colliding gravitational shock waves in  $AdS_5$  toy model for quark-gluon plasma formation in heavy ion collisions

- paper-and-pencil calculations with Romatschke 0803.3226
  - $\delta$ -like shocks
  - particle production in forward lightcone of shocks
  - Shortly after collision anisotropic pressure: P<sub>L</sub>/E = −3, P<sub>T</sub>/E = +2 confirmed numerically for thin shocks by Casalderrey-Solana, Heller, Mateos and van der Schee 1305.4919
  - Close to shockwaves negative energy fluxes ⇒ NEC violation! confirmed numerically and interpreted as absence of local rest frame by Arnold, Romatschke and van der Schee 1408.2518

 consider finite width gravitational shockwaves (pioneered numerically by Chesler and Yaffe 1011.3562)

- paper-and-pencil calculations with Romatschke 0803.3226
  - $\delta$ -like shocks
  - particle production in forward lightcone of shocks
  - Shortly after collision anisotropic pressure: P<sub>L</sub>/E = −3, P<sub>T</sub>/E = +2 confirmed numerically for thin shocks by Casalderrey-Solana, Heller, Mateos and van der Schee 1305.4919
  - Close to shockwaves negative energy fluxes ⇒ NEC violation! confirmed numerically and interpreted as absence of local rest frame by Arnold, Romatschke and van der Schee 1408.2518
- consider finite width gravitational shockwaves (pioneered numerically by Chesler and Yaffe 1011.3562)
- extract metric, holographic stress tensor and HEE numerically

- paper-and-pencil calculations with Romatschke 0803.3226
  - $\delta$ -like shocks
  - particle production in forward lightcone of shocks
  - Shortly after collision anisotropic pressure: P<sub>L</sub>/E = −3, P<sub>T</sub>/E = +2 confirmed numerically for thin shocks by Casalderrey-Solana, Heller, Mateos and van der Schee 1305.4919
  - Close to shockwaves negative energy fluxes ⇒ NEC violation! confirmed numerically and interpreted as absence of local rest frame by Arnold, Romatschke and van der Schee 1408.2518
- consider finite width gravitational shockwaves (pioneered numerically by Chesler and Yaffe 1011.3562)
- extract metric, holographic stress tensor and HEE numerically
- check QNEC and its saturation, particularly in region of NEC violation





# Outline

Inequalities

QNEC

Holographic QNEC in 4d

Holographic QNEC in 2d

- Definition:

A state is in quantum equilibrium when QNEC saturates for all times and all entangling regions

Consequences:



Consequences:

Far-from-(thermal)-equilibrium state can be in quantum equilibrium
Figure 2 from 1311.3655 (Nature Phys.)
Bhaseen, Doyon, Lucas, Schalm



- Far from equilibrium transport in strongly coupled CFT
- Long-time energy transport universally via steady-state
- In AdS<sub>3</sub>/CFT<sub>2</sub>: specific Bañados geometry with step function
- Our results imply QNEC saturation at all times

Holographic QNEC in 2d



Consequences:

Far-from-(thermal)-equilibrium state can be in quantum equilibrium



Daniel Grumiller — Quantum Null Energy Condition

- Definition:

A state is in quantum equilibrium when QNEC saturates for all times and all entangling regions

Consequences:

- Far-from-(thermal)-equilibrium state can be in quantum equilibrium
- All states dual to Bañados geometries are in quantum equilibrium

- Definition:

A state is in quantum equilibrium when QNEC saturates for all times and all entangling regions

Consequences:

- Far-from-(thermal)-equilibrium state can be in quantum equilibrium
- All states dual to Bañados geometries are in quantum equilibrium
- Natural to introduce "quantum equilibration time": For a given separation of the entangling interval quantum equilibration time = smallest time after which normalized QNEC non-saturation lower than prescribed bound (e.g. 1%)

- Definition:

A state is in quantum equilibrium when QNEC saturates for all times and all entangling regions

Consequences:

- Far-from-(thermal)-equilibrium state can be in quantum equilibrium
- All states dual to Bañados geometries are in quantum equilibrium
- Natural to introduce "quantum equilibration time": For a given separation of the entangling interval quantum equilibration time = smallest time after which normalized QNEC non-saturation lower than prescribed bound (e.g. 1%)

Quantum equilibrium hopefully a useful notion

• Vaidya = simple model for bulk matter; mass function M(t)

$$ds^{2} = \frac{1}{z^{2}} \left( -(1 - M(t)z^{2}) dt^{2} - 2 dt dz + dx^{2} \right)$$

• Vaidya = simple model for bulk matter; mass function M(t)

$$ds^{2} = \frac{1}{z^{2}} \left( -(1 - M(t)z^{2}) dt^{2} - 2 dt dz + dx^{2} \right)$$

Numerical studies show curious "half-saturation" for large entangling regions l

$$\lim_{k \gg 1} \frac{S'' + \frac{6}{c} \left(S'\right)^2}{\langle T_{kk} \rangle} \approx \frac{1}{2}$$

• Vaidya = simple model for bulk matter; mass function M(t)

$$ds^{2} = \frac{1}{z^{2}} \left( -(1 - M(t)z^{2}) dt^{2} - 2 dt dz + dx^{2} \right)$$

Numerical studies show curious "half-saturation" for large entangling regions l

$$\lim_{k \gg 1} \frac{S'' + \frac{6}{c} \, (S')^2}{\langle T_{kk} \rangle} \approx \frac{1}{2}$$

 $\blacktriangleright$  Can be derived perturbatively for  $M(t)=\epsilon\theta(t)$  with  $\epsilon\ll 1$ 

• Vaidya = simple model for bulk matter; mass function M(t)

$$ds^{2} = \frac{1}{z^{2}} \left( -(1 - M(t)z^{2}) dt^{2} - 2 dt dz + dx^{2} \right)$$

Numerical studies show curious "half-saturation" for large entangling regions l

$$\lim_{l\gg 1} \frac{S'' + \frac{6}{c} \, (S')^2}{\langle T_{kk} \rangle} \approx \frac{1}{2}$$

- $\blacktriangleright$  Can be derived perturbatively for  $M(t)=\epsilon\theta(t)$  with  $\epsilon\ll 1$
- ▶ If size of entangling region much larger than time,  $l \gg t_0$  we find QNEC half-saturation

$$\lim_{l \gg t_0} \frac{S'' + \frac{6}{c} \, (S')^2}{\langle T_{kk} \rangle} = \frac{1}{2} \pm \frac{t_0}{l} + \mathcal{O}(t_0^2/l^2) + \mathcal{O}(\epsilon)$$

• Vaidya = simple model for bulk matter; mass function M(t)

$$ds^{2} = \frac{1}{z^{2}} \left( -(1 - M(t)z^{2}) dt^{2} - 2 dt dz + dx^{2} \right)$$

Numerical studies show curious "half-saturation" for large entangling regions l

$$\lim_{l\gg 1} \frac{S'' + \frac{6}{c} \, (S')^2}{\langle T_{kk} \rangle} \approx \frac{1}{2}$$

- $\blacktriangleright$  Can be derived perturbatively for  $M(t)=\epsilon\theta(t)$  with  $\epsilon\ll 1$
- ▶ If size of entangling region much larger than time,  $l \gg t_0$  we find QNEC half-saturation

$$\lim_{l \gg t_0} \frac{S'' + \frac{6}{c} (S')^2}{\langle T_{kk} \rangle} = \frac{1}{2} \pm \frac{t_0}{l} + \mathcal{O}(t_0^2/l^2) + \mathcal{O}(\epsilon)$$

If time is much larger than entangling region we find QNEC saturation

$$\lim_{t_0 \gg l} \frac{S'' + \frac{6}{c} (S')^2}{\langle T_{kk} \rangle} = 1 + \mathcal{O}(\epsilon)$$
Consider global  $AdS_3$  with massive scalar field and take into account quantum backreactions (Belin, Iqbal, Lokhande 1805.08782)

$$ds^{2} = -(r^{2} + G_{1}(r)^{2}) dt^{2} + \frac{dr^{2}}{r^{2} + G_{2}(r)^{2}} + r^{2} d\varphi^{2}$$

with Newton's constant  $G=3/(2c){\rm ,}\ {\rm mass}\ m^2=4h(h-1)$  and

$$G_1(r) = 1 - 8Gh + \mathcal{O}(G^2) \qquad G_2(r) = 1 - 8Gh(1 - 1/(r^2 + 1)^{2h-1}) + \mathcal{O}(G^2)$$

Consider global  $AdS_3$  with massive scalar field and take into account quantum backreactions (Belin, Iqbal, Lokhande 1805.08782)

$$ds^{2} = -(r^{2} + G_{1}(r)^{2}) dt^{2} + \frac{dr^{2}}{r^{2} + G_{2}(r)^{2}} + r^{2} d\varphi^{2}$$

with Newton's constant G = 3/(2c), mass  $m^2 = 4h(h-1)$  and  $G_1(r) = 1-8Gh + \mathcal{O}(G^2)$   $G_2(r) = 1-8Gh(1-1/(r^2+1)^{2h-1}) + \mathcal{O}(G^2)$ 

Just using (H)RT to calculate HEE yields

$$S'' + \frac{6}{c} (S')^2 = -\frac{c}{24} + h - \frac{h\sqrt{\pi} \Gamma(2h+2)}{4 \Gamma(2h+\frac{3}{2})} \sin^{4h-2} \frac{\Delta\varphi}{2} + \mathcal{O}(1/c)$$

Interpretation for small interval,  $\Delta \varphi \ll 1$ : QNEC saturation up to polynomially suppressed terms

$$\langle T_{kk} \rangle = -\frac{c}{24} + h + \mathcal{O}(1/c)$$

Consider global  $AdS_3$  with massive scalar field and take into account quantum backreactions (Belin, Iqbal, Lokhande 1805.08782)

$$ds^{2} = -(r^{2} + G_{1}(r)^{2}) dt^{2} + \frac{dr^{2}}{r^{2} + G_{2}(r)^{2}} + r^{2} d\varphi^{2}$$

with Newton's constant G=3/(2c), mass  $m^2=4h(h-1)$  and

 $G_1(r) = 1 - 8Gh + \mathcal{O}(G^2)$   $G_2(r) = 1 - 8Gh(1 - 1/(r^2 + 1)^{2h-1}) + \mathcal{O}(G^2)$ 

Just using (H)RT to calculate HEE yields

$$S'' + \frac{6}{c} (S')^2 = -\frac{c}{24} + h - \frac{h\sqrt{\pi} \Gamma(2h+2)}{4 \Gamma(2h+\frac{3}{2})} \sin^{4h-2} \frac{\Delta\varphi}{2} + \mathcal{O}(1/c)$$

▶ For the half-interval  $(\Delta \varphi = \pi)$  large-h expansion interesting

Consider global  $AdS_3$  with massive scalar field and take into account quantum backreactions (Belin, Iqbal, Lokhande 1805.08782)

$$ds^{2} = -(r^{2} + G_{1}(r)^{2}) dt^{2} + \frac{dr^{2}}{r^{2} + G_{2}(r)^{2}} + r^{2} d\varphi^{2}$$

with Newton's constant G = 3/(2c), mass  $m^2 = 4h(h-1)$  and  $G_1(r) = 1 - 8Gh + \mathcal{O}(G^2)$   $G_2(r) = 1 - 8Gh(1 - 1/(r^2 + 1)^{2h-1}) + \mathcal{O}(G^2)$ 

 $G_1(r) = 1 - \delta G h + O(G') \qquad G_2(r) = 1 - \delta G h (1 - 1/(r' + 1))$ 

Just using (H)RT to calculate HEE yields

$$S'' + \frac{6}{c} (S')^2 = -\frac{c}{24} + h - \frac{h\sqrt{\pi} \Gamma(2h+2)}{4 \Gamma(2h+\frac{3}{2})} \sin^{4h-2} \frac{\Delta\varphi}{2} + \mathcal{O}(1/c)$$

For the half-interval (Δφ = π) large-h expansion interesting
 Just using (H)RT to calculate HEE yields

$$S'' + \frac{6}{c} (S')^2 = -\frac{c}{24} - \#h^{3/2} + \dots$$

Consider global  $AdS_3$  with massive scalar field and take into account quantum backreactions (Belin, Iqbal, Lokhande 1805.08782)

$$ds^{2} = -(r^{2} + G_{1}(r)^{2}) dt^{2} + \frac{dr^{2}}{r^{2} + G_{2}(r)^{2}} + r^{2} d\varphi^{2}$$

with Newton's constant G = 3/(2c), mass  $m^2 = 4h(h-1)$  and

 $G_1(r) = 1 - 8Gh + \mathcal{O}(G^2) \qquad G_2(r) = 1 - 8Gh(1 - 1/(r^2 + 1)^{2h-1}) + \mathcal{O}(G^2)$ 

Just using (H)RT to calculate HEE yields

$$S'' + \frac{6}{c} (S')^2 = -\frac{c}{24} + h - \frac{h\sqrt{\pi} \Gamma(2h+2)}{4 \Gamma(2h+\frac{3}{2})} \sin^{4h-2} \frac{\Delta\varphi}{2} + \mathcal{O}(1/c)$$

For the half-interval (Δφ = π) large-h expansion interesting
 Just using (H)RT to calculate HEE yields

$$S'' + \frac{6}{c} (S')^2 = -\frac{c}{24} - \#h^{3/2} + \dots$$

From CFT perspective hard to understand where  $h^{3/2}$  comes from

$$S_{\rm EE} = \frac{\rm Area}{4G} + S_{\rm bulk}$$

Faulkner, Lewcowycz, Maldacena 1307.2892

$$S_{\rm EE} = \frac{\rm Area}{4G} + S_{\rm bulk}$$

HEE bulk corrections for AdS<sub>3</sub> with massive scalar field backreactions calculated in Belin, Iqbal, Lokhande 1805.08782

$$S_{\rm EE} = \frac{\rm Area}{4G} + S_{\rm bulk}$$

- HEE bulk corrections for AdS<sub>3</sub> with massive scalar field backreactions calculated in Belin, Iqbal, Lokhande 1805.08782
- For QNEC essentially need to boost their results appropriately

$$S_{\rm EE} = \frac{\rm Area}{4G} + S_{\rm bulk}$$

- ► HEE bulk corrections for AdS<sub>3</sub> with massive scalar field backreactions calculated in Belin, Iqbal, Lokhande 1805.08782
- ► For QNEC essentially need to boost their results appropriately
- In case of half-interval get precise cancellation of offending h<sup>3/2</sup>-term!

$$S_{\rm EE} = \frac{\rm Area}{4G} + S_{\rm bulk}$$

- HEE bulk corrections for AdS<sub>3</sub> with massive scalar field backreactions calculated in Belin, Iqbal, Lokhande 1805.08782
- ► For QNEC essentially need to boost their results appropriately
- ▶ In case of half-interval get precise cancellation of offending  $h^{3/2}$ -term!
- ► Full result for the half-interval:

$$S'' + \frac{6}{c} (S')^2 = -\frac{c}{24} + \frac{3}{4} h + \mathcal{O}(\sqrt{h}) + \mathcal{O}(1/c)$$

Faulkner, Lewcowycz, Maldacena 1307.2892

$$S_{\rm EE} = \frac{\rm Area}{4G} + S_{\rm bulk}$$

- HEE bulk corrections for AdS<sub>3</sub> with massive scalar field backreactions calculated in Belin, Iqbal, Lokhande 1805.08782
- ► For QNEC essentially need to boost their results appropriately
- ▶ In case of half-interval get precise cancellation of offending  $h^{3/2}$ -term!
- ► Full result for the half-interval:

$$S'' + \frac{6}{c} (S')^2 = -\frac{c}{24} + \frac{3}{4} h + \mathcal{O}(\sqrt{h}) + \mathcal{O}(1/c)$$

Backreacted boundary stress tensor flux component

$$\langle T_{kk} \rangle = -\frac{c}{24} + h + \mathcal{O}(1/c)$$

Faulkner, Lewcowycz, Maldacena 1307.2892

$$S_{\rm EE} = \frac{\rm Area}{4G} + S_{\rm bulk}$$

- HEE bulk corrections for AdS<sub>3</sub> with massive scalar field backreactions calculated in Belin, Iqbal, Lokhande 1805.08782
- ► For QNEC essentially need to boost their results appropriately
- ▶ In case of half-interval get precise cancellation of offending  $h^{3/2}$ -term!
- ► Full result for the half-interval:

$$S'' + \frac{6}{c} (S')^2 = -\frac{c}{24} + \frac{3}{4} h + \mathcal{O}(\sqrt{h}) + \mathcal{O}(1/c)$$

Backreacted boundary stress tensor flux component

$$\langle T_{kk} \rangle = -\frac{c}{24} + h + \mathcal{O}(1/c)$$

QNEC non-saturation:  

$$\langle T_{kk} \rangle - S'' - \frac{6}{c} (S')^2 |_{\Delta \varphi = \pi, c \gg h \gg 1} = \frac{1}{4} h$$

Daniel Grumiller — Quantum Null Energy Condition

Holographic QNEC in 2d

Quantum Null Energy Condition (QNEC)

 QNEC in words: Expectation value of null projection of stress tensor bigger or equal than second derivative of entanglement entropy with respect to null deformations of entangling surface

- QNEC in words: Expectation value of null projection of stress tensor bigger or equal than second derivative of entanglement entropy with respect to null deformations of entangling surface
- Only known local energy condition that could be universally true

- QNEC in words: Expectation value of null projection of stress tensor bigger or equal than second derivative of entanglement entropy with respect to null deformations of entangling surface
- Only known local energy condition that could be universally true
- Various proofs of QNEC exist

- QNEC in words: Expectation value of null projection of stress tensor bigger or equal than second derivative of entanglement entropy with respect to null deformations of entangling surface
- Only known local energy condition that could be universally true
- Various proofs of QNEC exist
- Provided first numerical studies of QNEC in AdS<sub>5</sub>/CFT<sub>4</sub> context 1710.09837

- QNEC in words: Expectation value of null projection of stress tensor bigger or equal than second derivative of entanglement entropy with respect to null deformations of entangling surface
- Only known local energy condition that could be universally true
- Various proofs of QNEC exist
- Provided first numerical studies of QNEC in AdS<sub>5</sub>/CFT<sub>4</sub> context 1710.09837
- Shockwaves can saturate QNEC in far from equilibrium regime

- QNEC in words: Expectation value of null projection of stress tensor bigger or equal than second derivative of entanglement entropy with respect to null deformations of entangling surface
- Only known local energy condition that could be universally true
- Various proofs of QNEC exist
- Provided first numerical studies of QNEC in AdS<sub>5</sub>/CFT<sub>4</sub> context 1710.09837
- Shockwaves can saturate QNEC in far from equilibrium regime
- QNEC sharper in AdS<sub>3</sub>/CFT<sub>2</sub>

$$\langle T_{kk} \rangle \ge S'' + \frac{6}{c} S'^2$$

Quantum Null Energy Condition (QNEC)

- QNEC in words: Expectation value of null projection of stress tensor bigger or equal than second derivative of entanglement entropy with respect to null deformations of entangling surface
- Only known local energy condition that could be universally true
- Various proofs of QNEC exist
- Provided first numerical studies of QNEC in AdS<sub>5</sub>/CFT<sub>4</sub> context 1710.09837
- Shockwaves can saturate QNEC in far from equilibrium regime
- QNEC sharper in AdS<sub>3</sub>/CFT<sub>2</sub>

$$\langle T_{kk} \rangle \ge S'' + \frac{6}{c} S'^2$$

 QNEC saturation iff extremal surface does not cross bulk matter — "quantum equilibrium"

- QNEC in words: Expectation value of null projection of stress tensor bigger or equal than second derivative of entanglement entropy with respect to null deformations of entangling surface
- Only known local energy condition that could be universally true
- Various proofs of QNEC exist
- Provided first numerical studies of QNEC in AdS<sub>5</sub>/CFT<sub>4</sub> context 1710.09837
- Shockwaves can saturate QNEC in far from equilibrium regime
- QNEC sharper in AdS<sub>3</sub>/CFT<sub>2</sub>

$$\langle T_{kk} \rangle \ge S'' + \frac{6}{c} \, S'^2$$

- QNEC saturation iff extremal surface does not cross bulk matter "quantum equilibrium"
- Curious QNEC "half-saturation" in Vaidya

- QNEC in words: Expectation value of null projection of stress tensor bigger or equal than second derivative of entanglement entropy with respect to null deformations of entangling surface
- Only known local energy condition that could be universally true
- Various proofs of QNEC exist
- Provided first numerical studies of QNEC in AdS<sub>5</sub>/CFT<sub>4</sub> context 1710.09837
- Shockwaves can saturate QNEC in far from equilibrium regime
- QNEC sharper in AdS<sub>3</sub>/CFT<sub>2</sub>

$$\langle T_{kk} \rangle \ge S'' + \frac{6}{c} \, S'^2$$

- QNEC saturation iff extremal surface does not cross bulk matter "quantum equilibrium"
- Curious QNEC "half-saturation" in Vaidya
- 1/c corrections can spoil saturation and require to take into account bulk corrections to entanglement entropy

## Much to be learned about QNEC and its potential applications



Thanks for your attention!