

Flat space higher spin holography

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TU Wien

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based on work w. Afshar, Bagchi, Basu, Detournay, Fareghbal,
Gary, Riegler, Rosseel, Salzer, Sarkar, Schöller, Simon, ...

Outline

Motivations

Holography basics

Flat space holography

Flat space higher spin holography

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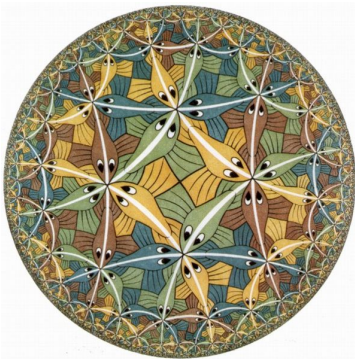
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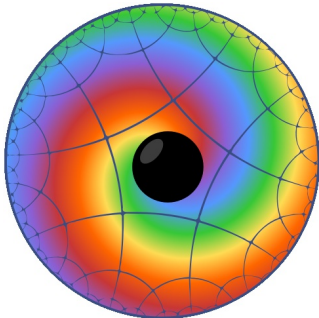
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 - ▶ Holographic principle realized in Nature? (yes/no)



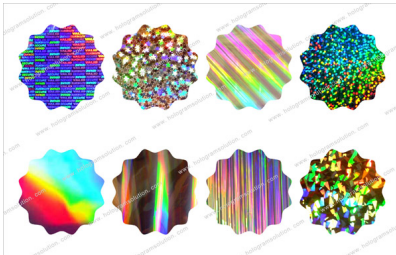
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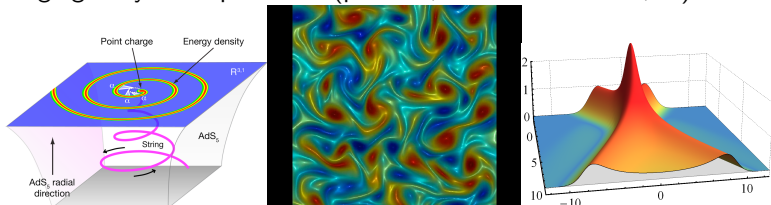
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- ▶ Applications
 - ▶ Gauge gravity correspondence (plasmas, condensed matter, ...)




Specific motivation for 3D

Gravity in 3D is simpler than in higher dimensions

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A quote by Albert Einstein: "Simplicity is the ultimate sophistication". The text is centered on a light gray background with a subtle gradient and a faint diagonal line.

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Specific motivation for flat space higher spin gravity

Massless higher spin theories constrained by no-gos!

- ▶ Coleman–Mandula '67
- ▶ Aragone–Deser '79
- ▶ Weinberg–Witten '80
- ▶ recent summary: [Bekaert, Boulanger, Sundell '12](#)

Conclusion: there are no consistent interacting massless higher spin theories in 4- (or higher-) dimensional flat space

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- ▶ Vasiliev '87-'90: higher spin theories in (A)dS
- ▶ Afshar, Bagchi, Fareghbal, DG, Rosseel '13; Gonzalez, Matulich, Pino, Troncoso '13: flat space higher spin theories in 3d

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Address these issues in 3D!



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AdS₃ gravity

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Interesting generic constraints from CFT₂!

e.g. [Hellerman '09](#), [Hartman, Keller, Stoica '14](#)

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Caveat: while there are many string compactifications with AdS₃ factor, applying holography just to AdS₃ factor does not capture everything!

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle

Example: Einstein gravity with Dirichlet boundary conditions

$$I = -\frac{1}{16\pi G_N} \int d^3x \sqrt{|g|} \left(R + \frac{2}{\ell^2} \right)$$

with $\delta g = \text{fixed}$ at the boundary

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions

Example: asymptotically AdS

$$ds^2 = d\rho^2 + \left(e^{2\rho/\ell} \gamma_{ij}^{(0)} + \gamma_{ij}^{(2)} + \dots \right) dx^i dx^j$$

with $\delta\gamma^{(0)} = 0$ for $\rho \rightarrow \infty$

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions
3. Perform canonical analysis and check consistency of bc's
 - ▶ Find and classify all constraints
 - ▶ Construct canonical gauge generators
 - ▶ Add boundary terms and get (variation of) canonical charges
 - ▶ Check integrability of canonical charges
 - ▶ Check finiteness of canonical charges
 - ▶ Check conservation (in time) of canonical charges
 - ▶ Calculate Dirac bracket algebra of canonical charges

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Example: Brown–Henneaux analysis for 3D Einstein gravity

$$\{Q[\varepsilon], Q[\eta]\} = \delta_\varepsilon Q[\eta]$$

$$Q[\varepsilon] \sim \oint d\varphi \mathcal{L}(\varphi) \varepsilon(\varphi)$$

$$\delta_\varepsilon \mathcal{L} = \mathcal{L} \varepsilon + 2\mathcal{L} \varepsilon' + \frac{\ell}{16\pi G_N} \varepsilon'''$$

Holographic algorithm from gravity point of view

Universal recipe:

1. Identify bulk theory and variational principle
2. Fix background and impose suitable boundary conditions
3. Perform canonical analysis and check consistency of bc's
4. Derive (classical) asymptotic symmetry algebra and central charges

Example: Two copies of Virasoro algebra

$$[\mathcal{L}_n, \mathcal{L}_m] = (n - m) \mathcal{L}_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0}$$

with Brown–Henneaux central charge

$$c = \frac{3\ell}{2G_N}$$

Reminder: ASA = quotient algebra of asymptotic symmetries by 'trivial' asymptotic symmetries with zero canonical charges

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5. Improve to quantum ASA

Example: semi-classical ASA in spin-3 gravity (Henneaux, Rey '10; Campoleoni, Pfenninger, Fredenhagen, Theisen '10)

$$[W_n, W_m] = \frac{16}{5c} \sum_p L_p L_{n+m-p} + \dots$$

quantum ASA

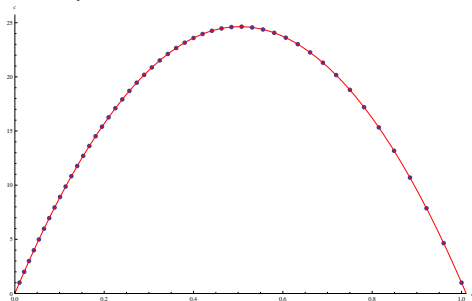
$$[W_n, W_m] = \frac{16}{5c + 22} \sum_p : L_p L_{n+m-p} : + \dots$$

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6. Study unitary representations of quantum ASA

Example:



Afshar et al '12

Discrete set of Newton
constant values compatible
with unitarity
(3D spin-N gravity in
next-to-principal embedding)

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7. Identify/constrain dual field theory

Example: Monster CFT in (flat space) chiral gravity

Witten '07

Li, Song & Strominger '08

Bagchi, Detournay & DG '12

$$Z(q) = J(q) = \frac{1}{q} + (1 + 196883)q + \mathcal{O}(q^2)$$

Note: $\ln 196883 \approx 12.2 = 4\pi + \text{quantum corrections}$

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Examples: too many!



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Goal of this talk:

Apply algorithm above to flat space holography in 3D higher spin theories

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- ▶ Example where it does not work at all: highest weight conditions!

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Interesting example:

- ▶ unitarity of flat space quantum gravity

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Many open issues in flat space holography!

(Higher spin) gravity as Chern–Simons gauge theory...

...with weird boundary conditions (Achúcarro & Townsend '86; Witten '88; Bañados '96)

CS action (for AdS: $\mathfrak{sl}(2) \oplus \mathfrak{sl}(2)$):

$$S_{\text{CS}} = \frac{k}{4\pi} \int \text{CS}(A) - \frac{k}{4\pi} \int \text{CS}(\bar{A})$$

Variational principle:

$$\delta S_{\text{CS}}|_{\text{EOM}} = \frac{k}{4\pi} \int \text{Tr} (A \wedge \delta A - \bar{A} \wedge \delta \bar{A})$$

Well-defined for boundary conditions (similarly for \bar{A})

$$A_+ = 0 \quad \text{or} \quad A_- = 0 \quad \text{boundary coordinates } x^\pm$$

Example: asymptotically AdS_3 (Cartan-version of Brown–Henneaux)

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$$\begin{aligned} A_\rho &= L_\rho & \bar{A}_\rho &= -L_\rho \\ A_+ &= e^\rho L_1 + e^{-\rho} L(x^+) L_{-1} & \bar{A}_+ &= 0 \\ A_- &= 0 & \bar{A}_- &= -e^\rho L_{-1} - e^{-\rho} \bar{L}(x^-) L_1 \end{aligned}$$

Dreibein: $e/\ell \sim A - \bar{A}$, spin-connection: $\omega \sim A + \bar{A}$

İnönü–Wigner contraction of Virasoro (Barnich & Compère '06)

BMS₃ and GCA₂ (or rather, URCA₂)

- ▶ Take two copies of Virasoro, generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$, central charges c, \bar{c}

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- ▶ Take two copies of Virasoro, generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$, central charges c, \bar{c}
- ▶ Define superrotations L_n and supertranslations M_n

$$L_n := \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n := \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

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- ▶ Take two copies of Virasoro, generators $\mathcal{L}_n, \bar{\mathcal{L}}_n$, central charges c, \bar{c}
- ▶ Define superrotations L_n and supertranslations M_n

$$L_n := \mathcal{L}_n - \bar{\mathcal{L}}_{-n} \quad M_n := \frac{1}{\ell} (\mathcal{L}_n + \bar{\mathcal{L}}_{-n})$$

- ▶ Make ultrarelativistic boost, $\ell \rightarrow \infty$

$$[L_n, L_m] = (n - m) L_{n+m} + c_L \frac{1}{12} \delta_{n+m,0}$$

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$$[M_n, M_m] = 0$$

Global part: contracted to $\text{isl}(2)$ (generators: $L_{\pm 1}, L_0, M_{\pm 1}, M_0$)

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- ▶ Is precisely the (centrally extended) BMS₃ algebra!
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Example TMG (with gravitational CS coupling μ and Newton constant G):

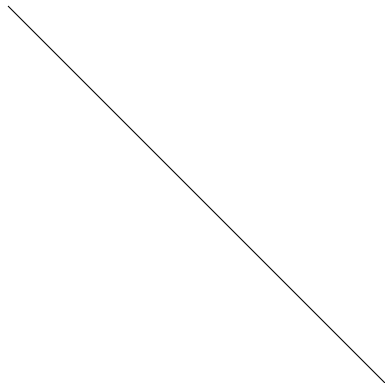
$$c_L = \frac{3}{\mu G} \quad c_M = \frac{3}{G}$$

Consequence of ultrarelativistic boost for AdS boundary

AdS-boundary:



Flat space boundary:



Limit $l \rightarrow \infty$

Null infinity holography!

Contraction on gravity side

AdS metric ($\varphi \sim \varphi + 2\pi$):

$$ds_{\text{AdS}}^2 = d(\ell\rho)^2 - \cosh^2\left(\frac{\ell\rho}{\ell}\right) dt^2 + \ell^2 \sinh^2\left(\frac{\ell\rho}{\ell}\right) d\varphi^2$$

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Limit $\ell \rightarrow \infty$ ($r = \ell\rho$):

$$ds_{\text{Flat}}^2 = dr^2 - dt^2 + r^2 d\varphi^2 = -du^2 - 2du dr + r^2 d\varphi^2$$

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BTZ metric:

$$ds_{\text{BTZ}}^2 = -\frac{\left(\frac{r^2}{\ell^2} - \frac{r_+^2}{\ell^2}\right)(r^2 - r_-^2)}{r^2} dt^2 + \frac{r^2 dr^2}{\left(\frac{r^2}{\ell^2} - \frac{r_+^2}{\ell^2}\right)(r^2 - r_-^2)} + r^2 \left(d\varphi - \frac{r_+}{r^2} r_- dt\right)^2$$

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Limit $\ell \rightarrow \infty$ ($\hat{r}_+ = \frac{r_+}{\ell} = \text{finite}$):

$$ds_{\text{FSC}}^2 = \hat{r}_+^2 \left(1 - \frac{r_-^2}{r^2}\right) dt^2 - \frac{1}{1 - \frac{r_-^2}{r^2}} \frac{dr^2}{\hat{r}_+^2} + r^2 \left(d\varphi - \frac{\hat{r}_+ r_-}{r^2} dt\right)^2$$

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Shifted-boost orbifold studied by Cornalba & Costa more than decade ago

Describes expanding (contracting) Universe in flat space

Cosmological horizon at $r = r_-$, screening CTCs at $r < 0$

Outline

Motivations

Holography basics

Flat space holography

Flat space higher spin holography

Flat space higher spin gravity

Afshar, Bagchi, Fareghbal, DG, Rosseel '13, Gonzalez, Matulich, Pino, Troncoso '13

- ▶ AdS gravity in CS formulation: $\text{spin } 2 \rightarrow \text{spin } 3 \sim \text{sl}(2) \rightarrow \text{sl}(3)$

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- ▶ Flat space: similar!

$$S_{\text{CS}}^{\text{flat}} = \frac{k}{4\pi} \int \text{CS}(\mathcal{A})$$

with $\mathfrak{isl}(3)$ connection ($e^a =$ “zuvielbein”)

$$\mathcal{A} = e^a T_a + \omega^a J_a \quad T_a = (M_n, V_m) \quad J_a = (L_n, U_m)$$

$\mathfrak{isl}(3)$ algebra (spin 3 extension of global part of BMS/GCA algebra)

$$[L_n, L_m] = (n - m)L_{n+m}$$

$$[L_n, M_m] = (n - m)M_{n+m}$$

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$$Q[\varepsilon_M, \varepsilon_L, \varepsilon_V, \varepsilon_U] \sim \oint (\varepsilon_M(\varphi)M(\varphi) + \varepsilon_L(\varphi)L(\varphi) + \varepsilon_V(\varphi)V(\varphi) + \varepsilon_U(\varphi)U(\varphi))$$

Flat space higher spin gravity

Asymptotic symmetry algebra at finite level k Afshar, Bagchi, Fareghbal, DG, Rosseel '13

- ▶ Do either Brown–Henneaux type of analysis or İnönü–Wigner contraction of two copies of quantum W_3 -algebra

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- ▶ Do either Brown–Henneaux type of analysis or İnönü–Wigner contraction of two copies of quantum W_3 -algebra
- ▶ Obtain new type of W -algebra as extension of BMS (“BMW”)

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

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$$[U_n, U_m] = (n - m)(2n^2 + 2m^2 - nm - 8)L_{n+m} + \frac{192}{c_M} (n - m)\Lambda_{n+m} \\ - \frac{96(c_L + \frac{44}{5})}{c_M^2} (n - m)\Theta_{n+m} + \frac{c_L}{12} n(n^2 - 1)(n^2 - 4) \delta_{n+m, 0}$$

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$$\Lambda_n = \sum_p : L_p M_{n-p} : - \frac{3}{10} (n + 2)(n + 3)M_n \quad \Theta_n = \sum_p M_p M_{n-p}$$

other commutators as in $\text{isl}(3)$ with $n \in \mathbb{Z}$

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- ▶ Note **quantum shift** and **poles** in central terms!
- ▶ Analysis generalizes to flat space contractions of other W -algebras

Unitarity in flat space

Unitarity leads to further contraction DG, Riegler, Rosseel '14

Facts:

- ▶ Unitarity in GCA requires $c_M = 0$ (see paper for caveats!)

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Doubly contracted algebra has unitary representations:

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Higher spin states decouple and become null states!

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Generic flat space W -algebras DG, Riegler, Rosseel '14

1. NO-GO:

Generically (see paper) you can have only two out of three:

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- ▶ Flat space
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Example:

Flat space chiral gravity

Bagchi, Detournay, DG, 1208.1658

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Example:

Minimal model holography

Gaberdiel, Gopakumar, 1011.2986, 1207.6697

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Example:

Flat space higher spin gravity (Galilean W_3 algebra)

Afshar, Bagchi, Fareghbal, DG and Rosseel, 1307.4768

Gonzalez, Matulich, Pino and Troncoso, 1307.5651

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2. YES-GO:

There is (at least) one counter-example, namely a Vasiliev-type of theory, where you can have all three properties!

Unitary, non-trivial flat space higher spin algebra exists!
Vasiliev-type flat space chiral higher spin gravity?

Unitarity in flat space

Flat space W_∞ -algebra compatible with unitarity DG, Riegler, Rosseel '14

- ▶ We do not know if flat space **chiral** higher spin gravity exists...

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- ▶ We do not know if flat space **chiral** higher spin gravity exists...
- ▶ ...but its existence is at least not ruled out by the no-go result!

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- ▶ We do not know if flat space **chiral** higher spin gravity exists...
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- ▶ If it exists, this must be its asymptotic symmetry algebra:

$$[\mathcal{V}_m^i, \mathcal{V}_n^j] = \sum_{r=0}^{\lfloor \frac{i+j}{2} \rfloor} g_{2r}^{ij}(m, n) \mathcal{V}_{m+n}^{i+j-2r} + c_{\mathcal{V}}^i(m) \delta^{ij} \delta_{m+n,0}$$

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where

$$c_{\mathcal{V}}^i(m) = \#(i, m) \times c \quad \text{and} \quad c = -\bar{c}$$

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$$[\mathcal{V}_m^i, \mathcal{V}_n^j] = \sum_{r=0}^{\lfloor \frac{i+j}{2} \rfloor} g_{2r}^{ij}(m, n) \mathcal{V}_{m+n}^{i+j-2r} + c_{\mathcal{V}}^i(m) \delta^{ij} \delta_{m+n,0}$$

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where

$$c_{\mathcal{V}}^i(m) = \#(i, m) \times c \quad \text{and} \quad c = -\bar{c}$$

- ▶ Vacuum descendants $\mathcal{W}_m^i|0\rangle$ are null states for all i and m !

Unitarity in flat space

Flat space W_∞ -algebra compatible with unitarity DG, Riegler, Rosseel '14

- ▶ We do not know if flat space **chiral** higher spin gravity exists...
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- ▶ Vacuum descendants $\mathcal{W}_m^i |0\rangle$ are null states for all i and m !
- ▶ AdS parent theory: no trace anomaly, but **gravitational anomaly** (Like in conformal Chern–Simons gravity \rightarrow Vasiliev type analogue?)

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Works nicely in Chern–Simons formulation!

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Line-element with spin-2 and spin-3 chemical potentials:

$$g_{\mu\nu} dx^\mu dx^\nu = \left(r^2 (\mu_L^2 - 4\mu_U'' \mu_U + 3\mu_U'^2 + 4\mathcal{M}\mu_U^2) + r g_{uu}^{(r)} + g_{uu}^{(0)} + g_{uu}^{(0')} \right) du^2 + \left(r^2 \mu_L - r\mu_M' + \mathcal{N}(1 + \mu_M) + 8\mathcal{Z}\mu_V \right) 2 du d\varphi - (1 + \mu_M) 2 dr du + r^2 d\varphi^2$$

$$g_{uu}^{(0)} = \mathcal{M}(1 + \mu_M)^2 + 2(1 + \mu_M)(\mathcal{N}\mu_L + 12\mathcal{V}\mu_V + 16\mathcal{Z}\mu_U) + 16\mathcal{Z}\mu_L\mu_V + \frac{4}{3}(\mathcal{M}^2\mu_V^2 + 4\mathcal{M}\mathcal{N}\mu_U\mu_V + \mathcal{N}^2\mu_U^2)$$

Spin-3 field with same chemical potentials:

$$\Phi_{\mu\nu\lambda} dx^\mu dx^\nu dx^\lambda = \Phi_{uuu} du^3 + \Phi_{ruu} dr du^2 + \Phi_{uu\varphi} du^2 d\varphi - (2\mu_U r^2 - r\mu_V' + 2\mathcal{N}\mu_V) dr du d\varphi + \mu_V dr^2 du - (\mu_U' r^3 - \frac{1}{3}r^2(\mu_V'' - \mathcal{M}\mu_V + 4\mathcal{N}\mu_U) + r\mathcal{N}\mu_V' - \mathcal{N}^2\mu_V) du d\varphi^2$$

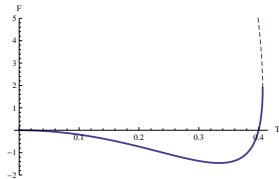
$$\begin{aligned} \Phi_{uuu} = & r^2 [2(1 + \mu_M)\mu_U(\mathcal{M}\mu_L - 4\mathcal{V}\mu_U) - \frac{1}{3}\mu_L^2(\mathcal{M}\mu_V - 4\mathcal{N}\mu_U) + 16\mu_L\mu_U(\mathcal{V}\mu_V + \mathcal{Z}\mu_U) - \frac{4}{3}\mathcal{M}\mu_U^2(\mathcal{M}\mu_V \\ & + 2\mathcal{N}\mu_U)] + 2\mathcal{V}(1 + \mu_M)^3 + \frac{2}{3}(1 + \mu_M)^2(6\mathcal{Z}\mu_L + \mathcal{M}^2\mu_V + 2\mathcal{M}\mathcal{N}\mu_U) + 16\mu_L\mu_V^2(\mathcal{N}\mathcal{V} - \frac{1}{3}\mathcal{M}\mathcal{Z}) \\ & + \frac{2}{3}(1 + \mu_M)((\mathcal{N}\mu_L + 16\mathcal{Z}\mu_U)(2\mathcal{M}\mu_V + \mathcal{N}\mu_U) + 12\mathcal{M}\mathcal{V}\mu_V^2) + \frac{64}{3}\mathcal{Z}\mu_U\mu_V(\mathcal{N}\mu_L + 12\mathcal{V}\mu_V + 12\mathcal{Z}\mu_U) \\ & + \mathcal{N}^2\mu_L^2\mu_V + 64\mathcal{V}^2\mu_V^3 - \frac{8}{27}(\mathcal{M}^3\mu_V^3 - \mathcal{N}^3\mu_U^3) - \frac{4}{9}\mathcal{M}\mathcal{N}\mu_U\mu_V(4\mathcal{M}\mu_V + 5\mathcal{N}\mu_U) + \sum_{n=0}^3 r^n \Phi_{uuu}^{(r^n)} \end{aligned}$$

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Interesting novel phase transitions of zeroth/first order:



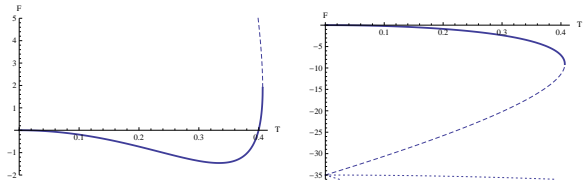
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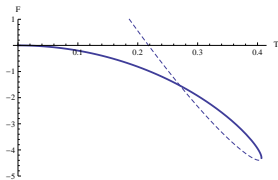
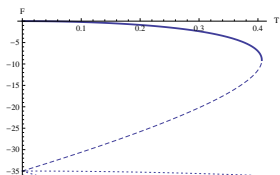
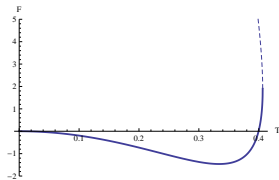
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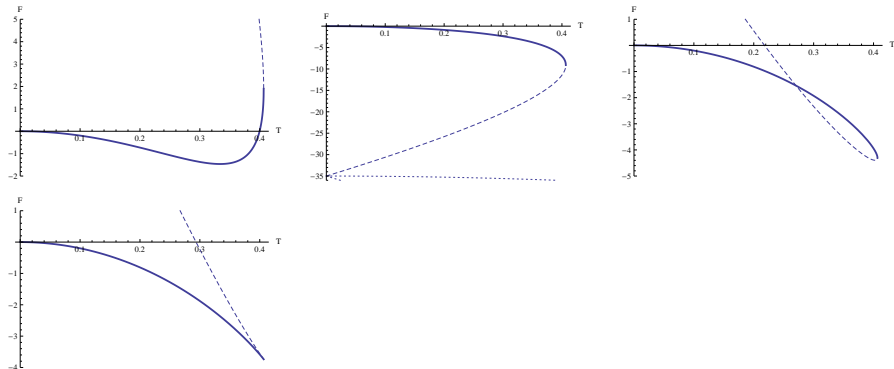
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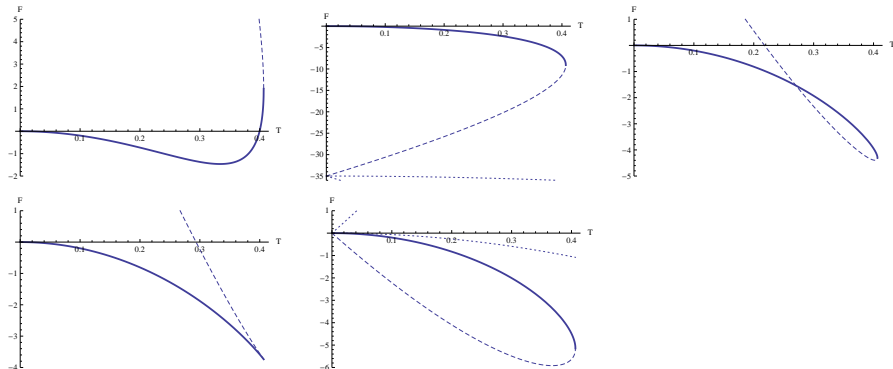
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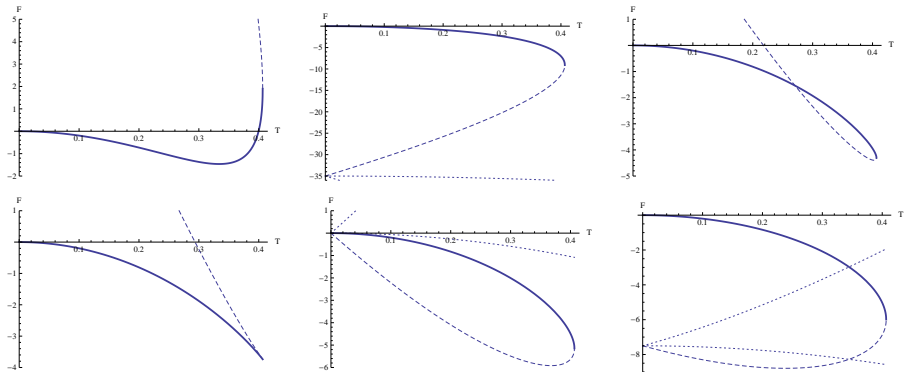
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Flat space higher spin holography provides a new playground
Contributes to long-term goal: find how general is holography

Thanks for your attention!

