# Black hole thermodynamics

Daniel Grumiller

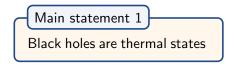
Institute for Theoretical Physics Vienna University of Technology

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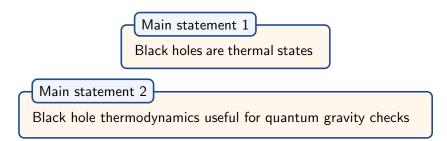
with R. McNees and J. Salzer: 1402.5127

#### Main statements and overview



- four laws of black hole mechanics/thermodynamics
- phase transitions between black holes and vacuum

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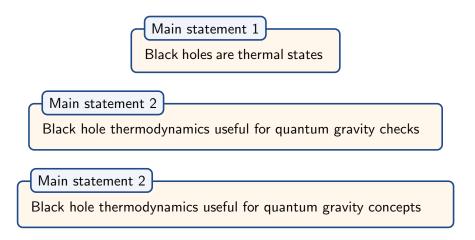


quantum gravity entropy matching with semi-classical prediction

$$S_{BH} = \underbrace{\frac{k_B c^3}{\hbar G_N}}_{=1 \text{ in this talk}} \frac{A_h}{4} + \mathcal{O}(\ln A_h)$$

semi-classical log corrections of entropy

#### Main statements and overview



- information loss, fuzzballs, firewalls, ...
- black hole holography, AdS/CFT, gauge/gravity correspondence, …

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$$S_{EH} \sim \frac{1}{\kappa} \int d^4x \sqrt{|g|} (\Lambda + R) + \text{marginal} + \text{irrelevant}$$

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$$ds^{2} = -(1 - 2M/r) dt^{2} + dr^{2}/(1 - 2M/r) + r^{2} d\Omega_{S^{2}}^{2}$$

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Everything phrased in terms of geometry! Classical gravity = geometry!

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causal limit  $p = \rho$ : speed of sound = speed of light

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Even before Bekenstein-Hawking:

Non-extensive entropy expected from/predicted by GR!

# Thermodynamics

Zeroth law:

T = const. in equilibrium

# T: temperature

Black hole mechanics Zeroth law:

 $\kappa = {\rm const.}\ {\rm f.}\ {\rm stationary}\ {\rm black}\ {\rm holes}$ 

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E: energy

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Third law:  $T \rightarrow 0$  impossible

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M: mass

A: area (of event horizon)

Formal analogy or actual physics?

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Hawking: indeed!

$$S_{BH} = \frac{1}{4} A_h$$

using semi-classical gravity

Hawking effect from QFT on fixed (curved) background: particle production with thermal spectrum at Hawking temperature

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• near horizon approximation:  $r = 2M + x^2/(8M)$ 

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- periodicity in Euclidean time = inverse temperature
- Result: Hawking temperature!

$$T_H = \frac{1}{8\pi M}$$

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- Exploit relationship between  $\mathcal{Z}$  and Euclidean partition function

$$\mathcal{Z} \sim e^{-\beta \Omega}$$

- Ω: thermodynamic potential for appropriate ensemble
- $\beta$ : periodicity in Euclidean time

# Free energy from Euclidean path integral Main idea

Consider Euclidean path integral (Gibbons, Hawking, 1977)

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Requires periodicity in Euclidean time and accessibility of semi-classical approximation

Consider small perturbation around classical solution

$$I_E[g_{cl} + \delta g, X_{cl} + \delta X] = I_E[g_{cl}, X_{cl}] + \delta I_E[g_{cl}, X_{cl}; \delta g, \delta X]$$
$$+ \frac{1}{2} \delta^2 I_E[g_{cl}, X_{cl}; \delta g, \delta X] + \dots$$

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If nothing goes wrong:

$$\left(\mathcal{Z} \sim \exp\left(-\frac{1}{\hbar} I_E[g_{cl}, X_{cl}]\right) \int \mathscr{D} \delta g \, \mathscr{D} \delta X \exp\left(-\frac{1}{2\hbar} \, \delta^2 I_E\right) \times \dots\right)$$

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$$\left[ \delta I_E \big|_{\text{EOM}} \sim \int_{\partial \mathcal{M}} dx \sqrt{\gamma} \left[ \pi^{ab} \, \delta \gamma_{ab} + \pi_X \, \delta X \right] \neq 0 \right]$$

Free energy from Euclidean path integral What could go Wrong? ...everything!

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3. Frequently violated: Gaussian integral may diverge Holographic renormalization resolves first and second problem!

Subtract suitable boundary terms from the action

$$\Gamma = I_E - I_{\rm CT}$$

such that second problem resolved; typically also resolves first problem

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$$\mathcal{Z} \sim \sum_{g_{cl}} \exp\left(-\frac{1}{\hbar} \Gamma[g_{cl}, X_{cl}]\right) \int \mathscr{D} \delta g \, \mathscr{D} \delta X \exp\left(-\frac{1}{2\hbar} \, \delta^2 \Gamma\right) \times \dots$$

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Leading term is finite

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Leading order (set  $\hbar = 1$ ):

$$\mathcal{Z}(T, X) = e^{-\Gamma(T, X)} = e^{-\beta F(T, X)}$$

Here F is the Helmholtz free energy

• Euclidean action: 
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- (Euclidean)  $AdS_3$  ( $t_E \sim t_E + \beta$ ,  $\varphi \sim \varphi + 2\pi$ )

$$\mathrm{d}s^2 = \cosh^2 \rho \, \mathrm{d}t_E^2 + \sinh^2 \rho \, \mathrm{d}\varphi^2 + \mathrm{d}\rho^2$$

yields free energy

$$F_{\text{AdS}} = -\frac{1}{8}$$

Example: Hawking–Page phase transition of AdS<sub>3</sub>/BTZ

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(non-rotating) BTZ BH

$$\mathrm{d}s^2 = -(r^2 - r_+^2)\,\mathrm{d}t^2 + \frac{\mathrm{d}r^2}{r^2 - r_+^2} + r^2\,\mathrm{d}\varphi^2$$

yields free energy  $\left(T=r_+/(2\pi)
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m DTTZ}=-rac{\pi^2T^2}{\pi^2}=-1$ 

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yields free energy ( $T=r_+/(2\pi)$ )

$$F_{\rm BTZ} = -\frac{\pi^2 T^2}{2} = -\frac{1}{8} \frac{T^2}{T_{\rm crit.}^2}$$

- Critical Hawking–Page temperature:  $T_{\rm crit.} = 1/(2\pi)$ 

Works also for flat space and expanding universe (in 2+1) Bagchi, Detournay, Grumiller & Simon PRL '13

Hot flat space

$$(\varphi \sim \varphi + 2\pi)$$

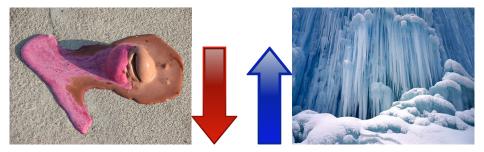
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$$(\varphi \sim \varphi + 2\pi)$$

$$\mathrm{d}s^2 = \mathrm{d}t^2 + \mathrm{d}r^2 + r^2 \,\mathrm{d}\varphi^2$$



$$ds^{2} = d\tau^{2} + \frac{(E\tau)^{2} dx^{2}}{1 + (E\tau)^{2}} + \left(1 + (E\tau)^{2}\right) \left(dy + \frac{(E\tau)^{2}}{1 + (E\tau)^{2}} dx\right)^{2}$$
  
Flat space cosmology  $(y \sim y + 2\pi r_{0})$ 

Daniel Grumiller — Black hole thermodynamics

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Regarding the other two main points: Black hole thermodynamics useful for quantum gravity checks/concepts

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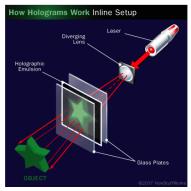
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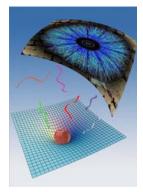
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#### Holography — Main idea aka gauge/gravity duality, aka AdS/CFT correspondence

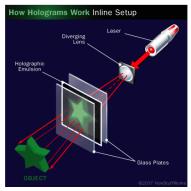


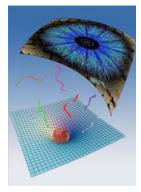


One of the most fruitful ideas in contemporary theoretical physics:

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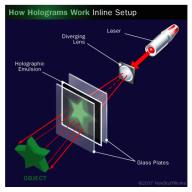


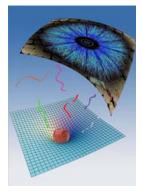


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One of the most fruitful ideas in contemporary theoretical physics:

- The number of dimensions is a matter of perspective
- We can choose to describe the same physical situation using two different formulations in two different dimensions
- The formulation in higher dimensions is a theory with gravity
- The formulation in lower dimensions is a theory without gravity

Why gravity? The holographic principle in black hole physics

Boltzmann/Planck: entropy of photon gas in d spatial dimensions  $S_{
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e.g. 
$$\langle T_{\mu\nu} \rangle_{\text{gauge}} = T^{BY}_{\mu\nu} \qquad \delta(\text{gravity action}) = \int d^d x \sqrt{|h|} T^{BY}_{\mu\nu} \,\delta h^{\mu\nu}$$

...and why were there > 9700 papers on holography in the past 17 years?

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We can expect many new applications in the next decade!

# Thanks for your attention!



## References

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